



Rahman's Biorthogonal Rational Functions and Superconformal Indices

Downloaded from: <https://research.chalmers.se>, 2024-04-20 03:47 UTC

Citation for the original published paper (version of record):

Rosengren, H. (2018). Rahman's Biorthogonal Rational Functions and Superconformal Indices. *Constructive Approximation*, 47(3): 529-552. <http://dx.doi.org/10.1007/s00365-017-9393-3>

N.B. When citing this work, cite the original published paper.



Rahman's Biorthogonal Rational Functions and Superconformal Indices

Hjalmar Rosengren¹

Received: 19 December 2016 / Accepted: 9 August 2017 / Published online: 25 September 2017
© The Author(s) 2017. This article is an open access publication

Abstract We study biorthogonal functions related to basic hypergeometric integrals with coupled continuous and discrete components. Such integrals appear as superconformal indices for three-dimensional quantum field theories and also in the context of solvable lattice models. We obtain explicit biorthogonal systems given by products of two of Rahman's biorthogonal rational ${}_{10}W_9$ -functions or their degenerate cases. We also give new bilateral extensions of the Jackson and q -Saalschütz summation formulas and new continuous and discrete biorthogonality measures for Rahman's functions.

Keywords Biorthogonal rational function · Basic hypergeometric series · Basic hypergeometric integral · Superconformal index

Mathematics Subject Classification 33D15 · 33D45 · 33D60 · 81T60

Dedicated to the memory of Mizan Rahman.

Communicated by Erik Koelink.

Supported by the Swedish Science Research Council (Vetenskapsrådet).

✉ Hjalmar Rosengren
hjalmar@chalmers.se
<http://www.math.chalmers.se/~hjalmar>

¹ Department of Mathematical Sciences, Chalmers University of Technology and University of Gothenburg, 412 96 Göteborg, Sweden

1 Introduction

During its long history, the notion of *classical orthogonal polynomials* has gradually been expanded to include more general systems. The most restrictive definition includes only Jacobi, Laguerre, and Hermite polynomials (with special cases such as Chebyshev and Legendre polynomials). By Sonine's theorem, these are precisely the orthogonal polynomials whose derivatives are again orthogonal. After the work of Hahn and others in the early twentieth century, it appeared natural to include also systems related to difference rather than differential equations. This development culminated in the Askey scheme of hypergeometric and basic hypergeometric orthogonal polynomials [18], consisting of Askey–Wilson and q -Racah polynomials together with numerous degenerate cases.

Jacobi polynomials are orthogonal with respect to the measure in Euler's beta integral

$$\int_{-1}^1 (1-x)^{a-1} (1+x)^{b-1} dx = 2^{a+b-1} \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)}.$$

More generally, the Askey scheme is related to a scheme of integral evaluations and summations, many of which are also called beta integrals. From this perspective, the Askey scheme appears to be incomplete. For instance, the q -Racah polynomials correspond to Rogers' ${}_6W_5$ -summation [10, Eq. (II.21)], which is a degenerate case of Jackson's ${}_8W_7$ -summation [10, Eq. (II.22)]. Thus, the top level in the Askey scheme corresponds to the next-to-top level in the scheme of beta integrals.

Mizan Rahman and Jim Wilson realized that to find the missing level in the Askey scheme, it is necessary to extend its scope from orthogonal polynomials to *biorthogonal rational functions*. To our knowledge, the earliest occurrence of hypergeometric biorthogonal rational functions is in Wilson's thesis [38], which contains discrete biorthogonality relations for rational functions of hypergeometric type ${}_9F_8$. This system generalizes Racah polynomials. Rahman [20] independently discovered some degenerate cases of Wilson's system. At the end of the paper, he remarks that Wilson had also found continuous biorthogonality relations; this piece of research seems never to have been published. Rahman [21] gave more general continuous biorthogonality relations for rational functions of type ${}_{10}W_9$, generalizing Askey–Wilson polynomials. The corresponding discrete system, generalizing q -Racah polynomials, was found by Wilson [39]. The complete extension of the Askey scheme to biorthogonal rational functions has been worked out (in the q -case) by van de Bult and Rains [2].

The work of Rahman and Wilson indicates that ${}_{10}W_9$ is the most general “classical” hypergeometric function. It must have been a great surprise to the special functions community when even more general functions turned up in the work of Date et al. [3] on solvable lattice models. Their elliptic $6j$ -symbols generalized ${}_{10}W_9$ -sums in a completely unexpected direction, providing the first known example of *elliptic hypergeometric functions*; see [27] for an introduction.

It took some time before the study of elliptic hypergeometric functions gained momentum. In one of the first papers on the subject, Spiridonov and Zhedanov [36]

constructed a system of biorthogonal rational functions generalizing Wilson's discrete $_{10}W_9$ -functions.

A first step towards elliptic extensions of Askey–Wilson polynomials was Spiridonov's discovery of the elliptic beta integral [31]

$$\oint \frac{\prod_{j=1}^6 \Gamma(t_j z; p, q) \Gamma(t_j / z; p, q)}{\Gamma(z^2; p, q) \Gamma(z^{-2}; p, q)} \frac{dz}{2\pi i z} = \frac{2}{(p; p)_\infty (q; q)_\infty} \prod_{1 \leq i < j \leq 6} \Gamma(t_i t_j; p, q). \quad (1.1)$$

Here, the integration is over a contour separating geometric sequences of poles going to infinity from sequences going to zero. The parameters satisfy the balancing condition $t_1 \cdots t_6 = pq$, the elliptic gamma function is defined by [29]

$$\Gamma(z; p, q) = \prod_{j,k=0}^{\infty} \frac{1 - p^{j+1} q^{k+1} / z}{1 - p^j q^k z}$$

and $(p; p)_\infty = \prod_{j=0}^{\infty} (1 - p^{j+1})$.

In [32], Spiridonov introduced systems biorthogonal with respect to the measure in (1.1). As a direct extension of Rahman's work, he found explicit functions Q_k and R_k satisfying a biorthogonality relation $\mu(Q_k R_l) = C_k \delta_{kl}$, where

$$\mu(f) = \frac{(p; p)_\infty (q; q)_\infty}{2 \prod_{1 \leq i < j \leq 6} \Gamma(t_i t_j; p, q)} \oint f(z) \frac{\prod_{j=1}^6 \Gamma(t_j z; p, q) \Gamma(t_j / z; p, q)}{\Gamma(z^2; p, q) \Gamma(z^{-2}; p, q)} \frac{dz}{2\pi i z}.$$

These functions are elliptic [in a multiplicative coordinate, namely, $Q_k(pz) = Q_k(z)$ and $R_k(pz) = R_k(z)$] and invariant under $z \mapsto z^{-1}$. It follows that they are rational after a change of variables. More generally, Spiridonov found that if \tilde{Q}_k and \tilde{R}_k denote the functions obtained from Q_k and R_k after interchanging p and q , then

$$\mu \left(Q_{k_1} \tilde{Q}_{k_2} R_{l_1} \tilde{R}_{l_2} \right) = C_{k_1} \tilde{C}_{k_2} \delta_{k_1 l_1} \delta_{k_2 l_2}. \quad (1.2)$$

This two-index biorthogonality for a one-variable integral is quite unusual. It is not a biorthogonality of rational functions as the functions involved are rational in two different variables.

An intriguing consequence of (1.2) is a property that we will call the *decoupling phenomenon*. We can write (1.2) as

$$\mu(fg) = \mu(f)\mu(g), \quad (1.3)$$

where $f = Q_{k_1} R_{l_1}$ and $g = \tilde{Q}_{k_2} \tilde{R}_{l_2}$. Taking linear combinations of these basis elements, it follows that (1.3) holds for any functions f and g satisfying appropriate

conditions. In particular, f and g should be elliptic with distinct periods [$f(pz) = f(z)$ and $g(qz) = g(z)$] and invariant under $z \mapsto z^{-1}$. There are also restrictions on the location of their poles. In the present work, we will investigate other instances of the decoupling phenomenon.

Dolan and Osborn [4] discovered that (1.1) (and more general integral identities of Rains [24]) can be interpreted in terms of supersymmetric quantum field theories. It is believed that if two theories are dual in the sense of Seiberg [30], then their so-called superconformal indices [17, 25] agree. For a particular pair of dual theories, this gives (1.1). Exploiting this idea has led to long lists of elliptic hypergeometric identities, most of which still lack rigorous proof, see, e.g., [34, 35].

The quantum field theories related to elliptic hypergeometric integrals live on four-dimensional space-times. In a similar way, one can obtain basic hypergeometric integrals from three-dimensional theories [6, 7, 11, 12, 14, 19, 41]. Interestingly, the resulting integrals are not of a type considered in the classical literature but involve a mixture of continuous and discrete integration (this can also happen for four-dimensional theories [16, 33], but then with a finite rather than infinite discrete component). As an example, a top level integral of this type is [7]

$$\begin{aligned} & \sum_{x=-\infty}^{\infty} \oint \frac{(1 - q^x z^2)(1 - q^x z^{-2})}{q^x z^{6x}} \\ & \prod_{j=1}^6 \frac{(q^{1+x/2}/b_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{N_j+x/2} b_j z, q^{N_j-x/2} b_j/z)_{\infty}} \frac{dz}{2\pi i z} \\ & = \frac{2}{\prod_{j=1}^6 q^{\binom{N_j}{2}} b_j^{N_j}} \prod_{1 \leq i < j \leq 6} \frac{(q/b_i b_j)_{\infty}}{(b_i b_j q^{N_i+N_j})_{\infty}}, \end{aligned} \quad (1.4)$$

valid for generic parameters b_j and integer parameters N_j subject to

$$b_1 \cdots b_6 = q \quad (1.5)$$

and $N_1 + \cdots + N_6 = 0$. The identity (1.4) and some related results can also be interpreted as star-triangle relations for solvable lattice models [5, 8, 15, 16, 40].

The main purpose of the present work is to investigate the “classical orthogonal polynomials” corresponding to the integral (1.4) and another integral from [7] [see (4.2) below]. It turns out that there is again a decoupling phenomenon such as (1.3), but with f a rational function of $zq^{-x/2}$ and g a rational function of $zq^{x/2}$ (subject to certain restrictions). We can then obtain two-index biorthogonal functions, which in the case of (1.4) are products of Rahman’s ${}_{10}W_9$ -functions. In the case of (4.2), there are biorthogonalities involving degenerate cases of Rahman’s functions due to Al-Salam and Ismail [1] and van de Bult and Rains [2]. An intriguing question is whether the decoupling phenomenon and the two-index biorthogonality have any relevance for physics, in the context of either quantum field theory or solvable lattice models.

The plan of our paper is as follows. In Sect. 2, we give two basic hypergeometric summation formulas, which provide bilateral extensions of the Jackson and

q -Saalschütz summation. Although these follow easily from known results, they seem not to have been stated explicitly before. In Sect. 3, we discuss Rahman's biorthogonal rational functions. We obtain some new results, in particular, new continuous and discrete biorthogonality measures. Finally, in Sect. 4, we study the decoupling phenomenon and related two-index biorthogonal systems for integrals arising as superconformal indices.

Throughout, q will be a fixed complex number with $0 < |q| < 1$. We also fix a choice of square root $q^{1/2}$. We follow the standard notation of [10], but suppress q from the notation. Thus, we write q -shifted factorials as

$$(a)_k = \begin{cases} (1-a)(1-aq)\cdots(1-aq^{k-1}), & k \in \mathbb{Z}_{\geq 0}, \\ ((1-aq^{-1})(1-aq^{-2})\cdots(1-aq^k))^{-1}, & k \in \mathbb{Z}_{< 0}, \end{cases}$$

and

$$(a)_\infty = \prod_{j=0}^{\infty} (1-aq^j).$$

We will use the theta function

$$\theta(x) = (x)_\infty (q/x)_\infty,$$

which satisfies the identity

$$\theta(xq^k) = (-1)^k q^{-\binom{k}{2}} x^{-k} \theta(x), \quad k \in \mathbb{Z}. \quad (1.6)$$

Repeated arguments of shifted factorials and theta functions stand for products, for instance,

$$(ab^\pm, c)_k = (ab)_k (a/b)_k (c)_k.$$

Finally, we recall the standard basic hypergeometric series

$$\begin{aligned} {}_{r+1}\phi_r \left(\begin{matrix} a_1, \dots, a_{r+1} \\ b_1, \dots, b_r \end{matrix}; z \right) &= \sum_{k=0}^{\infty} \frac{(a_1, \dots, a_{r+1})_k}{(q, b_1, \dots, b_r)_k} z^k, \\ {}_r\psi_r \left(\begin{matrix} a_1, \dots, a_r \\ b_1, \dots, b_r \end{matrix}; z \right) &= \sum_{k=-\infty}^{\infty} \frac{(a_1, \dots, a_r)_k}{(b_1, \dots, b_r)_k} z^k, \\ {}_{r+1}W_r(a; b_1, \dots, b_{r-2}; z) &= \sum_{k=0}^{\infty} \frac{1-aq^{2k}}{1-a} \frac{(a, b_1, \dots, b_{r-2})_k}{(q, aq/b_1, \dots, aq/b_{r-2})_k} z^k. \end{aligned}$$

2 Bilateral Summations

2.1 A Bilateral Jackson Summation

Although it is a straightforward consequence of results that have been known since the 1950s, the following identity seems to fill a gap in the literature.

Proposition 2.1 *If $\lambda, \mu, b_1, \dots, b_6$ are generic parameters subject to (1.5), then*

$$\left\{ \theta(\mu b_1, \dots, \mu b_6) - \frac{q}{\mu^2} \theta(\mu/b_1, \dots, \mu/b_6) \right\} \\ \times \frac{(1 - \lambda^2) \prod_{j=1}^6 (q\lambda^\pm/b_j)_\infty}{\theta(\mu/\lambda)} {}_8\psi_8 \left(\begin{matrix} \lambda q, -\lambda q, \lambda b_1, \dots, \lambda b_6 \\ \lambda, -\lambda, \lambda q/b_1, \dots, \lambda q/b_6 \end{matrix}; q \right) + \text{idem}(\lambda; \mu) \\ = (q)_\infty \theta(\lambda^2, \mu^2, \lambda\mu) \prod_{1 \leq i < j \leq 6} (q/b_i b_j)_\infty. \quad (2.1)$$

Here, we are using the notation

$$f(\lambda, \mu) + \text{idem}(\lambda; \mu) = f(\lambda, \mu) + f(\mu, \lambda).$$

If we let $\lambda = b_6$ and $\mu = b_1$ in (2.1), then both ${}_8\psi_8$ -series reduce to ${}_8W_7$ -series. The resulting identity can be written

$${}_8W_7 \left(b_6^2; b_1 b_6, \dots, b_5 b_6; q \right) - \frac{b_1}{b_6} \frac{(q b_6^2, q b_1/b_6)_\infty}{(q b_1^2, q b_6/b_1)_\infty} \prod_{j=2}^5 \frac{(q b_1/b_j, b_6 b_j)_\infty}{(q b_6/b_j, b_1 b_j)_\infty} \\ \times {}_8W_7 \left(b_1^2; b_1 b_2, \dots, b_1 b_6; q \right) = \frac{(q b_6^2, b_1/b_6)_\infty \prod_{2 \leq i < j \leq 5} (q/b_i b_j)_\infty}{\prod_{j=2}^5 (q b_6/b_j, b_1 b_j)_\infty}, \quad (2.2)$$

which is the nonterminating Jackson summation [10, Eq. (II.25)]. Consequently, we call (2.1) the *bilateral Jackson summation*.

To prove Proposition 2.1, we start from Jackson's transformation [13] (due to Margaret Jackson, not Frank Hilton Jackson of the Jackson summation)

$$\frac{\prod_{j=1}^6 (q a^\pm/b_j)_\infty}{(q a^{\pm 2})_\infty} {}_8\psi_8 \left(\begin{matrix} a q, -a q, a b_1, \dots, a b_6 \\ a, -a, a q/b_1, \dots, a q/b_6 \end{matrix}; \frac{q^2}{b_1 \cdots b_6} \right) \\ = \frac{\theta(\mu a^\pm) \prod_{j=1}^6 (q \lambda^\pm/b_j)_\infty}{\theta(\mu \lambda^\pm) (q \lambda^{\pm 2})_\infty} \\ {}_8\psi_8 \left(\begin{matrix} \lambda q, -\lambda q, \lambda b_1, \dots, \lambda b_6 \\ \lambda, -\lambda, \lambda q/b_1, \dots, \lambda q/b_6 \end{matrix}; \frac{q^2}{b_1 \cdots b_6} \right) + \text{idem}(\lambda; \mu),$$

which is a special case of Slater's transformations for ${}_{2r}\psi_{2r}$ -series [10, Eq. (5.5.2)]. We will only need the case $a = b_6$; that is,

$$\begin{aligned} & \frac{(q)_\infty \prod_{j=1}^5 (qb_6^\pm/b_j)_\infty}{(qb_6^2)_\infty} {}_8W_7(b_6^2; b_1b_6, \dots, b_5b_6; q^2/b_1 \cdots b_6) \\ &= \frac{\theta(\mu b_6^\pm) \prod_{j=1}^6 (q\lambda^\pm/b_j)_\infty}{\theta(\mu\lambda^\pm)(q\lambda^{\pm 2})_\infty} \\ & {}_8\psi_8 \left(\begin{matrix} \lambda q, -\lambda q, \lambda b_1, \dots, \lambda b_6 \\ \lambda, -\lambda, \lambda q/b_1, \dots, \lambda q/b_6 \end{matrix}; \frac{q^2}{b_1 \cdots b_6} \right) \\ & + \text{idem}(\lambda; \mu). \end{aligned} \quad (2.3)$$

Assuming (1.5), we may apply (2.3) to both ${}_8W_7$ series in (2.2), obtaining after simplification

$$\begin{aligned} & \left\{ \theta(b_1b_2, \dots, b_1b_5, \mu b_6^\pm) - \frac{b_1}{b_6} \theta(b_2b_6, \dots, b_5b_6, \mu b_1^\pm) \right\} \\ & \times \frac{\prod_{j=1}^6 (q\lambda^\pm/b_j)_\infty}{\theta(\mu\lambda^\pm)(q\lambda^{\pm 2})_\infty} {}_8\psi_8 \left(\begin{matrix} \lambda q, -\lambda q, \lambda b_1, \dots, \lambda b_6 \\ \lambda, -\lambda, \lambda q/b_1, \dots, \lambda q/b_6 \end{matrix}; q \right) + \text{idem}(\lambda; \mu) \\ & = (q)_\infty \theta(b_1/b_6) \prod_{1 \leq i < j \leq 6} (q/b_i b_j)_\infty. \end{aligned}$$

The proof of (2.1) is now reduced to the theta function identity

$$\begin{aligned} & \left\{ \theta(b_1b_2, \dots, b_1b_5, \mu b_6^\pm) - \frac{b_1}{b_6} \theta(b_2b_6, \dots, b_5b_6, \mu b_1^\pm) \right\} \\ & = \frac{\theta(b_1/b_6)}{\theta(\mu^2)} \left\{ \theta(\mu b_1, \dots, \mu b_6) - \frac{q}{\mu^2} \theta(\mu/b_1, \dots, \mu/b_6) \right\}, \end{aligned} \quad (2.4)$$

which is a special case of [37, Ex. 20.53.3], see also [26, Lemma 6.4].

2.2 A Bilateral q -Saalschütz Summation

The analogue of Proposition 2.1 at the level of the Saalschütz summation combines three rather than two bilateral sums. We find it convenient to write the identity as a determinant evaluation.

Proposition 2.2 *For generic parameters subject to*

$$a_1a_2a_3b_1b_2b_3 = q, \quad (2.5)$$

let

$$\begin{aligned} f_1(\lambda) &= \lambda \theta(b_1/\lambda, b_2/\lambda, b_3/\lambda), \\ f_2(\lambda) &= \lambda^{-1} \theta(a_1\lambda, a_2\lambda, a_3\lambda), \\ f_3(\lambda) &= \prod_{j=1}^3 (q/a_j\lambda, q\lambda/b_j)_\infty \cdot {}_3\psi_3 \left(\begin{matrix} a_1\lambda, a_2\lambda, a_3\lambda \\ q\lambda/b_1, q\lambda/b_2, q\lambda/b_3 \end{matrix}; q \right). \end{aligned}$$

Then,

$$\det_{1 \leq i, j \leq 3} (f_j(\lambda_i)) = \frac{(q)_\infty \theta(t\lambda_1\lambda_2\lambda_3)}{\lambda_1\lambda_2\lambda_3} \prod_{i,j=1}^3 (q/a_i b_j)_\infty \prod_{1 \leq i < j \leq 3} \lambda_j \theta(\lambda_i/\lambda_j), \quad (2.6)$$

where

$$t = a_1 a_2 a_3 = q/b_1 b_2 b_3. \quad (2.7)$$

One may prove Proposition 2.2 in a similar way as Proposition 2.1, but we prefer to illustrate a different method (which can also be adapted to give an alternative proof of Proposition 2.1). We first observe that each of the three functions f_j satisfies $f_j(q\lambda) = -f_j(\lambda)/qt\lambda^3$. For f_1 and f_2 , this follows from the quasi-periodicity $\theta(qx) = -\theta(x)/x$, and for f_3 , it is easy to verify by shifting the summation index. It then follows from [28, Prop. 3.4] that (2.6) holds up to a factor independent of the variables λ_j . Thus, it is enough to verify (2.6) for fixed values of these variables. We choose $\lambda_2 = b_2, \lambda_3 = b_3$, so that $f_1(\lambda_2) = f_1(\lambda_3) = 0$. We may then cancel the factors involving λ_1 and are left with the identity

$$\begin{aligned} & b_3(a_1 b_2, a_2 b_2, a_3 b_2, b_3 q/b_1, b_3 q/b_2)_\infty {}_3\phi_2 \left(\begin{matrix} a_1 b_3, a_2 b_3, a_3 b_3 \\ b_3 q/b_1, b_3 q/b_2 \end{matrix}; q \right) \\ & - \text{idem}(b_2; b_3) = b_3 \theta(b_2/b_3) \prod_{j=1}^3 (q/b_1 a_j)_\infty, \end{aligned}$$

which is the nonterminating q -Saalschütz summation [10, Eq. (II.24)]. This proves (2.6) and explains why we call it the bilateral q -Saalschütz summation.

3 Biorthogonal Rational Functions

3.1 The Rahman Functional

Let b_1, \dots, b_6 be complex numbers subject to the balancing condition (1.5). We assume that they are generic in the sense that

$$b_i/b_j \notin q^{\mathbb{Z}}, \quad 1 \leq i < j \leq 6, \quad b_i b_j \notin q^{\mathbb{Z}}, \quad 1 \leq i \leq j \leq 6. \quad (3.1)$$

Let V denote the vector space of rational functions in the variable $(z + z^{-1})/2$ that are regular except for possible single poles at

$$z^{\pm} \in b_j q^{\mathbb{Z}_{<0}}, \quad 1 \leq j \leq 6. \quad (3.2)$$

We also require that the elements of V are regular at infinity; that is, their numerator does not have larger degree than their denominator.

Lemma 3.1 *The space V is spanned by the rational functions*

$$(b_1 z^{\pm})_{k_1} \cdots (b_6 z^{\pm})_{k_6}, \quad (3.3)$$

where k_j are integers such that

$$k_1 + \cdots + k_6 = 0. \quad (3.4)$$

Proof Given $f \in V$, let n_j denote the largest positive integer such that f has a pole at $z = b_j q^{-n_j}$; if there is no such pole, we let $n_j = 0$. Then,

$$f\left(\frac{z + z^{-1}}{2}\right) = (b_1 z^{\pm})_{-n_1} \cdots (b_6 z^{\pm})_{-n_6} p\left(\frac{z + z^{-1}}{2}\right),$$

with p a polynomial of degree at most $N = n_1 + \cdots + n_6$. If we can expand

$$p\left(\frac{z + z^{-1}}{2}\right) = \sum_{\substack{l_1, \dots, l_6 \geq 0, \\ l_1 + \dots + l_6 = N}} C_l \prod_{j=1}^6 (b_j q^{-n_j} z^{\pm})_{l_j}, \quad (3.5)$$

then

$$f\left(\frac{z + z^{-1}}{2}\right) = \sum_{\substack{k_1, \dots, k_6 \geq 0, \\ k_1 + \dots + k_6 = 0}} C_{k+n} \prod_{j=1}^6 (b_j z^{\pm})_{k_j},$$

and we are done.

We claim that there is an expansion of the form (3.5) with $l_3 = \cdots = l_6 = 0$. Indeed, assuming the condition (3.1), the polynomials

$$(b_1 q^{-n_1} z^{\pm})_l (b_2 q^{-n_2} z^{\pm})_{N-l}, \quad l = 0, \dots, N,$$

form a basis for the space of polynomials of degree at most N , see, e.g., [26, Lemma 3.1]. \square

We now recall Rahman's integral evaluation

$$\oint \frac{(z^{\pm 2})_{\infty} \theta(\lambda z^{\pm})}{(b_1 z^{\pm}, \dots, b_6 z^{\pm})_{\infty}} \frac{dz}{2\pi iz} = \frac{2 \{ \theta(\lambda b_1, \dots, \lambda b_6) - q \lambda^{-2} \theta(\lambda/b_1, \dots, \lambda/b_6) \}}{(q)_{\infty} \theta(\lambda^2) \prod_{1 \leq i < j \leq 6} (b_i b_j)_{\infty}}, \quad (3.6)$$

where (1.5) is assumed. The integration is over a positively oriented contour separating the poles at $b_j q^{\mathbb{Z}_{\geq 0}}$ from their reciprocals. The identity (3.6) is a special case of an integral from [21] (see [10, Ex. 6.7]), which was explicitly stated in [22].

Let us replace b_j in (3.6) with $b_j q^{k_j}$, where k_j are integers subject to (3.4). Using (1.6) and other elementary facts, the resulting identity can be expressed as

$$\oint \frac{(z^{\pm 2})_{\infty} \theta(\lambda z^{\pm})}{(b_1 z^{\pm}, \dots, b_6 z^{\pm})_{\infty}} (b_1 z^{\pm})_{k_1} \cdots (b_6 z^{\pm})_{k_6} \frac{dz}{2\pi i z} \\ = \frac{2 \{ \theta(\lambda b_1, \dots, \lambda b_6) - q \lambda^{-2} \theta(\lambda/b_1, \dots, \lambda/b_6) \}}{(q)_{\infty} \theta(\lambda^2) \prod_{1 \leq i < j \leq 6} (b_i b_j)_{\infty}} \\ \frac{\prod_{1 \leq i < j \leq 6} (b_i b_j)_{k_i + k_j}}{\prod_{j=1}^6 q^{\binom{k_j}{2}} b_j^{k_j}}.$$

By Lemma 3.1, this proves the following result.

Proposition 3.2 *There exists a functional \mathbf{J} on the space V such that*

$$\mathbf{J}((b_1 z^{\pm})_{k_1} \cdots (b_6 z^{\pm})_{k_6}) = \frac{\prod_{1 \leq i < j \leq 6} (b_i b_j)_{k_i + k_j}}{\prod_{j=1}^6 q^{\binom{k_j}{2}} b_j^{k_j}}. \quad (3.7)$$

For generic values of λ , it is given by

$$\mathbf{J}(f) = \frac{(q)_{\infty} \theta(\lambda^2) \prod_{1 \leq i < j \leq 6} (b_i b_j)_{\infty}}{2 \{ \theta(\lambda b_1, \dots, \lambda b_6) - q \lambda^{-2} \theta(\lambda/b_1, \dots, \lambda/b_6) \}} \\ \times \oint f \left(\frac{z + z^{-1}}{2} \right) \frac{(z^{\pm 2})_{\infty} \theta(\lambda z^{\pm})}{(b_1 z^{\pm}, \dots, b_6 z^{\pm})_{\infty}} \frac{dz}{2\pi i z}, \quad (3.8)$$

where the integration is over a positively oriented contour encircling all poles of the integrand of the form $z \in b_j q^{\mathbb{Z}}$ and no poles of the form $z \in b_j^{-1} q^{\mathbb{Z}}$.

Note that the choice of contour in (3.8) depends on f . If n_j are defined as in the proof of Lemma 3.1, then the points $z = b_j q^{k - n_j}$, $k \in \mathbb{Z}_{\geq 0}$, should be inside the contour of integration and their reciprocals outside. For fixed f , the existence of such a contour follows from (3.1), but no contour works for all f .

We will refer to \mathbf{J} as *Rahman's functional*. Rahman preferred to work with (3.8) in the special case $\lambda = b_6$. This yields the more compact but less symmetric expression

$$\mathbf{J}(f) = \frac{(q)_{\infty} \prod_{1 \leq i < j \leq 5} (b_i b_j)_{\infty}}{2 \prod_{j=1}^5 (b_1 \cdots b_5 / b_j)_{\infty}} \oint f \left(\frac{z + z^{-1}}{2} \right) \\ \frac{(z^{\pm 2}, b_1 \cdots b_5 z^{\pm})_{\infty}}{(b_1 z^{\pm}, \dots, b_5 z^{\pm})_{\infty}} \frac{dz}{2\pi i z}. \quad (3.9)$$

We will also obtain expressions for \mathbf{J} in terms of discrete measures. To this end, replace b_j in (2.1) by $b_j q^{k_j}$, where, as before, we assume (3.4). The resulting identity can be written

$$\begin{aligned} & \frac{(1 - \lambda^2) \left\{ \theta(\mu b_1, \dots, \mu b_6) - q \mu^{-2} \theta(\mu/b_1, \dots, \mu/b_6) \right\} \prod_{j=1}^6 (q \lambda^{\pm}/b_j)_{\infty}}{(q)_{\infty} \theta(\lambda^2, \mu^2, \mu \lambda^{\pm}) \prod_{1 \leq i < j \leq 6} (q/b_i b_j)_{\infty}} \\ & \times \sum_{x=-\infty}^{\infty} \frac{1 - \lambda^2 q^{2x}}{1 - \lambda^2} q^x \prod_{j=1}^6 \frac{(\lambda b_j)_x}{(q \lambda/b_j)_x} (\lambda b_j q^x, b_j q^{-x}/\lambda)_{k_j} \\ & + \text{idem}(\lambda; \mu) = \frac{\prod_{1 \leq i < j \leq 6} (b_i b_j)_{k_i + k_j}}{\prod_{j=1}^6 q^{\binom{k_j}{2}} b_j^{k_j}}. \end{aligned}$$

This proves the following result.

Proposition 3.3 *For generic values of λ and μ ,*

$$\begin{aligned} \mathbf{J}(f) &= \frac{(1 - \lambda^2) \left\{ \theta(\mu b_1, \dots, \mu b_6) - \frac{q}{\mu^2} \theta(\mu/b_1, \dots, \mu/b_6) \right\} \prod_{j=1}^6 (q \lambda^{\pm}/b_j)_{\infty}}{(q)_{\infty} \theta(\lambda^2, \mu^2, \lambda \mu, \mu/\lambda) \prod_{1 \leq i < j \leq 6} (q/b_i b_j)_{\infty}} \\ & \times \sum_{x=-\infty}^{\infty} \frac{1 - \lambda^2 q^{2x}}{1 - \lambda^2} q^x \prod_{j=1}^6 \frac{(\lambda b_j)_x}{(q \lambda/b_j)_x} f\left(\frac{\lambda q^x + \lambda^{-1} q^{-x}}{2}\right) \\ & + \text{idem}(\lambda; \mu). \end{aligned} \quad (3.10)$$

In contrast to (3.8), where the contour of integration depends on f , the identity (3.10) holds uniformly on the whole space V .

Assume for simplicity that $f((z + z^{-1})/2)$ has no poles at $z^{\pm} \in b_5 q^{\mathbb{Z}_{<0}}$ or $b_6 q^{\mathbb{Z}_{<0}}$. Then, we may let $\lambda = b_5$ and $\mu = b_6$ in (3.10) and conclude that

$$\begin{aligned} \mathbf{J}(f) &= \frac{\prod_{j=1}^4 (q b_5/b_j, b_j b_6)_{\infty}}{(q b_5^2, b_6/b_5)_{\infty} \prod_{1 \leq i < j \leq 4} (q/b_i b_j)_{\infty}} \\ & \times \sum_{x=0}^{\infty} \frac{1 - b_5^2 q^{2x}}{1 - b_5^2} q^x \prod_{j=1}^6 \frac{(b_5 b_j)_x}{(q b_5/b_j)_x} f\left(\frac{b_5 q^x + b_5^{-1} q^{-x}}{2}\right) \\ & + \text{idem}(b_5; b_6). \end{aligned} \quad (3.11)$$

For general functions in V , one obtains in the same way an expression involving also residues of f at points corresponding to $x < 0$. This is discussed in [2, Prop. 6.4], but the general case of (3.10) seems to be new.

Van de Bult and Rains [2] also found integral formulas for \mathbf{J} that are nonsymmetric in the sense that the measure is not invariant under $z \mapsto 1/z$. We will now give a more general result of this type.

Proposition 3.4 *The functional \mathbf{J} can be expressed as*

$$\mathbf{J}(f) = \frac{1}{C} \oint f\left(\frac{z+z^{-1}}{2}\right) \frac{z-z^{-1}}{z^2} \prod_{j=1}^6 \frac{\theta(\lambda_j z)}{(b_j z^{\pm})_{\infty}} \frac{dz}{2\pi i z}, \quad (3.12)$$

where the integration is as above, λ_j are generic parameters subject to $\lambda_1 \cdots \lambda_6 = q$, and

$$C = \frac{\prod_{j=1}^4 \theta(\lambda_j \lambda_2) \left\{ \prod_{j=1}^6 \theta(\lambda_1 b_j) - q \lambda_1^{-2} \prod_{j=1}^6 \theta(\lambda_1 / b_j) \right\}}{\lambda_2 \theta(\lambda_1^2, \lambda_1 / \lambda_2)(q)_{\infty} \prod_{1 \leq i < j \leq 6} (b_i b_j)_{\infty}} + \text{idem}(\lambda_1; \lambda_2).$$

Proof Symmetrizing the integral, we may write the right-hand side of (3.12) as

$$\frac{1}{2C} \oint f\left(\frac{z+z^{-1}}{2}\right) \frac{(z-z^{-1}) \left\{ z^{-2} \prod_{j=1}^6 \theta(\lambda_j z) - z^2 \prod_{j=1}^6 \theta(\lambda_j / z) \right\}}{\prod_{j=1}^6 (b_j z^{\pm})_{\infty}} \frac{dz}{2\pi i z}.$$

By (2.4), with μ replaced with z and b_j with an appropriate permutation of λ_j , the factor in brackets can be expressed as

$$\frac{z^{-1} \theta(z^2, \lambda_2 \lambda_3, \dots, \lambda_2 \lambda_6, \lambda_1 z^{\pm})}{\lambda_2 \theta(\lambda_1 / \lambda_2)} + \text{idem}(\lambda_1; \lambda_2).$$

As

$$z^{-1}(z-z^{-1})\theta(z^2) = (z^{\pm 2})_{\infty},$$

we may then apply (3.8) to express the integral in terms of \mathbf{J} . \square

The nonsymmetric integral formulas given in [2] correspond to the special case of (3.12) when $\lambda_j = b_j$ for $j = 1, \dots, 4$. In that case,

$$C = \frac{\theta(b_1 b_2, b_1 b_3, b_1 b_4, b_2 b_3, b_2 b_4)}{b_2 \theta(b_1 / b_2)(q)_{\infty} \prod_{1 \leq i < j \leq 4} (b_i b_j)_{\infty}} \\ \times \{ \theta(b_2 \lambda_5, b_2 \lambda_6, b_1 b_5, b_1 b_6) - \theta(b_1 \lambda_5, b_1 \lambda_2, b_2 b_5, b_2 b_6) \}.$$

Since $b_5 b_6 = \lambda_5 \lambda_6$, we may apply Weierstrass' identity (see, e.g., [27, §1.4]) to conclude that the factor in brackets equals

$$-b_2 \lambda_6 \theta(b_1 / b_2, b_3 b_4, \lambda_5 / b_5, \lambda_5 / b_6).$$

In the corresponding specialization of (3.12), we find it convenient to replace z with $1/z$ and write $\lambda = q/\lambda_5$. We then obtain

$$\begin{aligned} \mathbf{J}(f) = & \frac{b_5 b_6 \lambda (q, b_5 b_6)_\infty \prod_{j=1}^4 (b_j b_5, b_j b_6)_\infty}{\theta(b_5 \lambda, b_6 \lambda) \prod_{1 \leq i < j \leq 4} (q/b_i b_j)_\infty} \\ & \times \oint f\left(\frac{z+z^{-1}}{2}\right) (z-z^{-1}) \frac{\theta(\lambda z, z/b_5 b_6 \lambda)}{(b_5 z^\pm, b_6 z^\pm)_\infty} \prod_{j=1}^4 \frac{(qz/b_j)_\infty}{(b_j z)_\infty} \frac{dz}{2\pi i z}. \end{aligned} \quad (3.13)$$

This is also contained in [2, Prop. 6.3]. As is discussed in [2], (3.11) can be obtained from (3.13) through shrinking the contour of integration to zero, picking up residues at the points $b_5 q^{\mathbb{Z}_{\geq 0}}$ and $b_6 q^{\mathbb{Z}_{\geq 0}}$.

3.2 Rahman's Biorthogonal Functions

Rahman [21, 23] constructed rational functions that are biorthogonal with respect to the functional \mathbf{J} . To explain this, we introduce the subspace V_n^i of V consisting of rational functions having no poles except possibly at $z^\pm = b_i q^{-k-1}$, where $i = 1, \dots, 6$ is fixed and $0 \leq k \leq n$. It is easy to see that $\dim(V_n^i) = n+1$. If $1 \leq i \neq j \leq 6$, we may look for a function $Q_n^{(i,j)} \in V_n^i$ such that $I(Q_n^{(i,j)} f) = 0$ for all $f \in V_{n-1}^j$. As this is an n -dimensional condition on an $(n+1)$ -dimensional space, it is not surprising that for generic parameters the solution is unique up to normalization. We then have the biorthogonality relations $\mathbf{J}(Q_n^{(i,j)} Q_m^{(j,i)}) = 0$ for $m \neq n$.

We may assume that $i = 5, j = 6$. Then, $Q_n^{(5,6)}$ is given by the function (our notation differs from that of Rahman)

$$\begin{aligned} Q_n\left(\frac{z+z^{-1}}{2}; b_1, \dots, b_6; q\right) &= \frac{(b_1 b_2, b_1 b_3, b_1 b_4, 1/b_1 b_6)_n}{(q b_1/b_5)_n} \\ &\times {}_{10}W_9(b_1/b_5; b_1 z, b_1/z, q/b_2 b_5, q/b_3 b_5, q/b_4 b_5, q^n/b_5 b_6, q^{-n}; q). \end{aligned}$$

It follows from Bailey's ${}_{10}W_9$ -transformation [10, Eq. (III.28)] that Q_n is symmetric in the parameters b_1, \dots, b_4 .

More precisely, Rahman proved that if

$$Q_n = Q_n\left(\frac{z+z^{-1}}{2}; b_1, b_2, b_3, b_4, b_5, b_6\right), \quad (3.14a)$$

$$R_n = Q_n\left(\frac{z+z^{-1}}{2}; b_1, b_2, b_3, b_4, b_6, b_5\right), \quad (3.14b)$$

then

$$\mathbf{J}(Q_m R_n) = \delta_{mn} \frac{1 - q^{-1} b_1 b_2 b_3 b_4}{1 - q^{2n-1} b_1 b_2 b_3 b_4} \frac{(q)_n \prod_{1 \leq i < j \leq 4} (b_i b_j)_n}{q^n (q^{-1} b_1 b_2 b_3 b_4)_n}. \quad (3.15)$$

Rahman formulated this result with \mathbf{J} given by (3.9) (and modifications, where the contour is deformed to the unit circle and then moved to an interval by a change of variables). Using instead (3.8), (3.10), or (3.12) gives new explicit forms of (3.15). Note, in particular, that (3.10) gives biorthogonality relations with respect to a fixed measure for all m and n , whereas in (3.8), (3.9), and (3.12), the integrals must be deformed (or discrete terms added) as m and n increase.

3.3 The Al-Salam–Ismail Functional

There are many interesting limit cases of Rahman's biorthogonal functions [2]. From the perspective of superconformal indices (see Sect. 4), a particularly interesting case is when three of the parameters b_j tend to zero and the remaining three to infinity, while the variable z is scaled so that either zb_j or z/b_j is fixed for each j . To describe this limit, we make the substitutions

$$(b_1, b_2, b_3, b_4, b_5, b_6) \mapsto (a_1/c, a_2/c, a_3/c, b_1c, b_2c, b_3c),$$

so that the balancing condition (1.5) becomes (2.5). We also replace z by cz . The possible poles (3.2) are then located at

$$z \in a_j^{-1}q^{\mathbb{Z}_{>0}}, \quad z \in b_jq^{\mathbb{Z}_{<0}}, \quad 1 \leq j \leq 3, \quad (3.16)$$

together with additional poles tending to zero or infinity as c tends to zero.

More concretely, let $f((z+z^{-1})/2)$ denote the spanning function (3.3). We replace (k_1, \dots, k_6) by $(k_1, k_2, k_3, l_1, l_2, l_3)$ and write

$$T = k_1 + k_2 + k_3 = -l_1 - l_2 - l_3. \quad (3.17)$$

Then,

$$\begin{aligned} \lim_{c \rightarrow 0} c^{2T} f\left(\frac{(cz) + (cz)^{-1}}{2}\right) &= \lim_{c \rightarrow 0} \prod_{j=1}^3 c^{2k_j} (a_j z, a_j/c^2 z)_{k_j} (b_j/z, b_j c^2 z)_{l_j} \\ &= (-1)^T \prod_{j=1}^3 q^{\binom{k_j}{2}} a_j^{k_j} \cdot g(z), \end{aligned}$$

where

$$g(z) = z^{-T} \prod_{j=1}^3 (a_j z)_{k_j} (b_j/z)_{l_j}. \quad (3.18)$$

It is easy to see that, as k_j, l_j , and T vary over integers subject to (3.17), the functions g span the space of rational functions in z that are regular everywhere (including at

infinity) except for possible single poles at (3.16). We will denote this space by W . We also compute

$$\begin{aligned} \lim_{c \rightarrow 0} c^{2T} \mathbf{J}(f) &= \lim_{c \rightarrow 0} \frac{c^{2T} \prod_{1 \leq i < j \leq 3} (a_i a_j / c^2)_{k_i + k_j} (b_i b_j c^2)_{l_i + l_j} \prod_{i,j=1}^3 (a_i b_j)_{k_i + l_j}}{\prod_{j=1}^3 q^{\binom{k_j}{2} + \binom{l_j}{2}} (a_j / c)^{k_j} (b_j c)^{l_j}} \\ &= t^T q^{\binom{T}{2}} \frac{\prod_{i,j=1}^3 (a_i b_j)_{k_i + l_j}}{\prod_{j=1}^3 q^{\binom{l_j}{2}} b_j^{l_j}}, \end{aligned}$$

where t is as in (2.7). Thus, \mathbf{J} degenerates to a functional \mathbf{K} on W , which acts on the spanning functions (3.18) as

$$\mathbf{K}(g) = \frac{(-t)^T q^{\binom{T}{2}} \prod_{i,j=1}^3 (a_i b_j)_{k_i + l_j}}{\prod_{j=1}^3 q^{\binom{k_j}{2} + \binom{l_j}{2}} a_j^{k_j} b_j^{l_j}}. \quad (3.19)$$

As we explain in Sect. 3.4, Al-Salam and Ismail constructed explicit biorthogonal systems for the functional \mathbf{K} [1]. For this reason, we call it the Al-Salam–Ismail functional.

The reader may find it puzzling that \mathbf{J} and its degenerate case \mathbf{K} seem to depend on the same number of parameters. The explanation is that we can simultaneously scale z , $1/a_j$, and b_j by the same constant without changing \mathbf{K} . This effectively reduces the number of parameters by one. We also mention the symmetry

$$\tilde{\mathbf{K}}(\tilde{g}) = \mathbf{K}(g), \quad g \in W, \quad (3.20)$$

where $\tilde{g}(z) = g(1/z)$ and $\tilde{\mathbf{K}}$ is obtained from \mathbf{K} after interchanging the parameters a_j and b_j .

One may obtain an integral expression for \mathbf{K} from Gasper's identity [9], [10, Eq. (4.11.3)]

$$\oint \frac{(qz/b_3)_\infty \theta(\lambda z, qz/\lambda b_1 b_2)}{(a_1 z, a_2 z, a_3 z, b_1/z, b_2/z)_\infty} \frac{dz}{2\pi i z} = \frac{\theta(b_1 \lambda, b_2 \lambda)}{(q)_\infty} \prod_{j=1}^3 \frac{(q/b_3 a_j)_\infty}{(b_1 a_j, b_2 a_j)_\infty}, \quad (3.21)$$

which holds for generic parameters subject to (2.5). More symmetric integral formulas arise from the following one-parameter extension of (3.21).

Proposition 3.5 *For generic parameters subject to the relations*

$$\lambda_1 \lambda_2 \lambda_3 = q a_1 a_2 a_3 = q^2 / b_1 b_2 b_3, \quad (3.22)$$

$$\oint \prod_{j=1}^3 \frac{\theta(\lambda_j z)}{(a_j z, b_j/z)_\infty} \frac{dz}{2\pi i z} = \frac{1}{(q)_\infty \theta(\lambda_1/\lambda_2) \prod_{i,j=1}^3 (a_i b_j)_\infty} \times \left\{ \prod_{j=1}^3 \theta(\lambda_1/a_j, \lambda_2 b_j) - \frac{\lambda_1}{\lambda_2} \prod_{j=1}^3 \theta(\lambda_1 b_j, \lambda_2/a_j) \right\}. \quad (3.23)$$

To prove Proposition 3.5, we will use (3.21) and the following fact. The proof is a standard elliptic function argument that we include for completeness.

Lemma 3.6 *Let f be an analytic function on $\mathbb{C} \setminus \{0\}$ that satisfies*

$$f(qx) = f(x)/q^3 x^6, \quad f(1/x) = -f(x). \quad (3.24)$$

Moreover, let c_1 and c_2 be numbers such that $c_1^2, c_2^2, c_1 c_2, c_1/c_2 \notin q^{\mathbb{Z}}$. Then, f is uniquely determined by the values $f(c_1)$ and $f(c_2)$.

Proof Let

$$g(x) = x^{-1} \theta(x^2) \left(\frac{\theta(c_2 x^\pm)}{c_1^{-1} \theta(c_1^2, c_2 c_1^\pm)} f(c_1) + \frac{\theta(c_1 x^\pm)}{c_2^{-1} \theta(c_2^2, c_1 c_2^\pm)} f(c_2) \right).$$

We claim that $f(x) = g(x)$ for all x . This is clearly true for $x = c_1$. Since g satisfies (3.24), it is more generally true for $x^\pm \in c_1 q^{\mathbb{Z}}$. We also observe that (3.24) implies that f vanishes for $x^2 \in q^{\mathbb{Z}}$. Hence, the function

$$h(x) = \frac{f(x) - g(x)}{x^{-1} \theta(x^2, c_1 x^\pm)}$$

is analytic for $x \neq 0$ and satisfies $h(qx) = h(x)$. By Liouville's theorem, h is constant. Since $h(c_2) = 0$, we arrive at the desired conclusion. \square

Proof of Proposition 3.5 We substitute $\lambda_1 = \lambda x$ and $\lambda_2 = \lambda/x$ in (3.23) and multiply the identity by $x^{-1} \theta(x^2)$. Consider both sides as functions of x . It is easy to see that they satisfy (3.24) (as the contour of integration can be chosen independently of x , it is enough to consider the integrand). By Lemma 3.6, it then suffices to verify (3.23) for two generic values of λ_1 . Choosing $\lambda_1 = q/b_3$, (3.23) is reduced to (3.21). By symmetry, we may take $\lambda_1 = q/b_2$ as the other value. \square

Let us now replace a_j and b_j in (3.23) by $a_j q^{k_j}$ and $b_j q^{l_j}$, where k_j and l_j are integers subject to (3.17). In order to respect (3.22), we also replace λ_3 by $\lambda_3 q^T$. The resulting identity can be expressed as

$$\begin{aligned} \mathbf{K}(g) &= \frac{(q)_\infty \theta(\lambda_1/\lambda_2) \prod_{i,j=1}^3 (a_i b_j)_\infty}{\prod_{j=1}^3 \theta(\lambda_1/a_j, \lambda_2 b_j) - \frac{\lambda_1}{\lambda_2} \prod_{j=1}^3 \theta(\lambda_1 b_j, \lambda_2/a_j)} \\ &\quad \times \oint g(z) \prod_{j=1}^3 \frac{\theta(\lambda_j z)}{(a_j z, b_j/z)_\infty} \frac{dz}{2\pi i z}, \end{aligned} \quad (3.25)$$

where g is as in (3.18). As these functions span W , we may draw the following conclusion.

Proposition 3.7 *The functional \mathbf{K} is given by the integral formula (3.25), where λ_j are generic parameters subject to (3.22).*

The special case corresponding to (3.21) is

$$\begin{aligned} \mathbf{K}(g) &= \frac{(q)_\infty}{\theta(b_1 \lambda, b_2 \lambda)} \prod_{j=1}^3 \frac{(b_1 a_j, b_2 a_j)_\infty}{(q/b_3 a_j)_\infty} \\ &\quad \times \oint g(z) \frac{(qz/b_3)_\infty \theta(\lambda z, qz/\lambda b_1 b_2)}{(a_1 z, a_2 z, a_3 z, b_1/z, b_2/z)_\infty} \frac{dz}{2\pi i z}. \end{aligned} \quad (3.26)$$

By (3.20), we may alternatively write

$$\begin{aligned} \mathbf{K}(g) &= \frac{(q)_\infty}{\theta(\lambda/a_1, \lambda/a_2)} \prod_{j=1}^3 \frac{(a_1 b_j, a_2 b_j)_\infty}{(q/a_3 b_j)_\infty} \\ &\quad \times \oint g(z) \frac{(q/a_3 z)_\infty \theta(\lambda z, qz a_1 a_2/\lambda)}{(a_1 z, a_2 z, b_1/z, b_2/z, b_3/z)_\infty} \frac{dz}{2\pi i z}. \end{aligned} \quad (3.27)$$

We also mention the further specialization $\lambda = a_3$ of (3.26) (or $\lambda = q/b_3$ of (3.27)),

$$\begin{aligned} \mathbf{K}(g) &= \frac{(q)_\infty \prod_{i,j=1}^2 (a_i b_j)_\infty}{(q/a_3 b_3)_\infty \prod_{j=1}^2 (q/a_3 b_j, q/b_3 a_j)_\infty} \\ &\quad \times \oint g(z) \frac{(qz/b_3, q/a_3 z)_\infty \theta(a_1 a_2 b_3 z)}{(a_1 z, a_2 z, b_1/z, b_2/z)_\infty} \frac{dz}{2\pi i z}. \end{aligned} \quad (3.28)$$

To obtain discrete integral formulas for \mathbf{K} , we replace a_j by $a_j q^{k_j}$ and b_j by $b_j q^{l_j}$ in (2.6), where we still assume (3.17). This leads to the following result.

Proposition 3.8 *For generic parameters λ_j , the functional \mathbf{K} can be expressed as*

$$\mathbf{K}(g) = C \sum_{k=1}^3 X_k \sum_{x=-\infty}^{\infty} q^x g(\lambda_k q^x) \prod_{j=1}^3 \frac{(a_j \lambda_k)_x}{(q \lambda_k / b_j)_x}, \quad (3.29)$$

where

$$\begin{aligned}
 C &= \frac{\lambda_1 \lambda_2 \lambda_3}{(q)_\infty \theta(t \lambda_1 \lambda_2 \lambda_3) \prod_{i,j=1}^3 (q/a_i b_j)_\infty \prod_{1 \leq i < j \leq 3} \lambda_j \theta(\lambda_i / \lambda_j)}, \\
 X_1 &= \prod_{j=1}^3 (q/a_j \lambda_1, q \lambda_1 / b_j)_\infty \left\{ \frac{\lambda_2}{\lambda_3} \prod_{j=1}^3 \theta(a_j \lambda_3, b_j / \lambda_2) - \frac{\lambda_3}{\lambda_2} \prod_{j=1}^3 \theta(a_j \lambda_2, b_j / \lambda_3) \right\}, \\
 X_2 &= \prod_{j=1}^3 (q/a_j \lambda_2, q \lambda_2 / b_j)_\infty \left\{ \frac{\lambda_3}{\lambda_1} \prod_{j=1}^3 \theta(a_j \lambda_1, b_j / \lambda_3) - \frac{\lambda_1}{\lambda_3} \prod_{j=1}^3 \theta(a_j \lambda_3, b_j / \lambda_1) \right\}, \\
 X_3 &= \prod_{j=1}^3 (q/a_j \lambda_3, q \lambda_3 / b_j)_\infty \left\{ \frac{\lambda_1}{\lambda_2} \prod_{j=1}^3 \theta(a_j \lambda_2, b_j / \lambda_1) - \frac{\lambda_2}{\lambda_1} \prod_{j=1}^3 \theta(a_j \lambda_1, b_j / \lambda_2) \right\}.
 \end{aligned}$$

Note that, in contrast to Proposition 3.7, we are not assuming any balancing condition for the parameters λ_j .

Analogously to (3.11), assuming that g has no poles at $z \in b_1 q^{\mathbb{Z}_{<0}}$ or $b_2 q^{\mathbb{Z}_{<0}}$, we may choose $\lambda_1 = b_1$ and $\lambda_2 = b_2$ in (3.29). This leads to the expression

$$\begin{aligned}
 \mathbf{K}(g) &= \frac{1}{b_1 \theta(b_2 / b_1) \prod_{j=1}^3 (q/a_j b_3)_\infty} \left(b_1 (q b_1 / b_2, q b_1 / b_3, a_1 b_2, a_2 b_2, a_3 b_2)_\infty \right. \\
 &\quad \times \sum_{x=0}^{\infty} q^x g(b_1 q^x) \prod_{j=1}^3 \frac{(a_j b_1)_x}{(q b_1 / b_j)_x} - \text{idem}(b_1; b_2) \Big), \quad (3.30)
 \end{aligned}$$

which can alternatively be obtained from (3.26) by shrinking the contour to zero, picking up residues at the points $b_1 q^{\mathbb{Z}_{\geq 0}}$ and $b_2 q^{\mathbb{Z}_{\geq 0}}$. If one instead expands the contour in (3.27), or substitutes $\lambda_1 = 1/a_1$, $\lambda_2 = 1/a_2$ in (3.29), one finds a similar identity involving values of g at $a_j q^{\mathbb{Z}_{\leq 0}}$.

3.4 Biorthogonal Rational Functions

Let us now consider the limit of Rahman's biorthogonal functions corresponding to the functional \mathbf{K} . There are two essentially different cases, depending on whether the distinguished parameters b_5 and b_6 in (3.15) tend to the same or distinct limits in $\{0, \infty\}$.

For the first case, we write

$$\begin{aligned}
 &q_n(z; b_1, b_2, b_3, a_1, a_2, a_3; q) \\
 &= \lim_{c \rightarrow 0} Q_n \left(\frac{(cz) + (cz)^{-1}}{2}; b_1 c, b_2 c, b_3 c, a_1 / c, a_2 / c, a_3 / c; q \right) \\
 &= (a_1 b_1, 1/a_3 b_1)_n {}_4\phi_3 \left(\begin{matrix} q^{-n}, q/a_2 b_2, q/a_2 b_3, b_1/z \\ a_1 b_1, q^{1-n} a_3 b_1, q/a_2 z \end{matrix}; q \right).
 \end{aligned}$$

This function is symmetric in the parameters (b_1, b_2, b_3) . Though we label q_n by seven parameters (counting q), it is effectively a five-parameter system since (2.5) is assumed and we may scale z , a_j^{-1} and b_j by a fixed constant. If r_n denotes the function q_n with a_2 and a_3 interchanged, then it follows from (3.15) that

$$\mathbf{K}(q_m r_n) = \delta_{mn} q^{-n}(q, a_1 b_1, a_1 b_2, a_1 b_3)_n.$$

Applying (3.20), we also find that if

$$\begin{aligned} \tilde{q}_n(z) &= q_n(1/z; a_1, a_2, a_3, b_1, b_2, b_3; q) \\ &= (a_1 b_1, 1/a_1 b_3)_n {}_4\phi_3 \left(\begin{matrix} q^{-n}, q/a_2 b_2, q/a_3 b_2, a_1 z \\ a_1 b_1, q^{1-n} a_1 b_3, qz/b_2 \end{matrix}; q \right) \end{aligned}$$

and \tilde{r}_n denotes \tilde{q}_n with b_2 and b_3 interchanged, then

$$\mathbf{K}(\tilde{q}_m \tilde{r}_n) = \delta_{mn} q^{-n}(q, a_1 b_1, a_2 b_1, a_3 b_1)_n.$$

For reasons that will be clear in Sect. 4, we think of this as a separate system, even though it is equivalent to (q_n, r_n) by a change of parameters.

The biorthogonal system (q_n, r_n) appears in the scheme of van de Bult and Rains [2] with the label 3100v2. The biorthogonality relations described there correspond to (3.26) and (3.30), possibly after permuting the parameters and applying the symmetry (3.20). Our expressions (3.25) and (3.29) unify and generalize these relations. Just as for Rahman's functions, in the continuous case the contour of integration depends on m and n , whereas the discrete measures are fixed.

For the second case, we write

$$\begin{aligned} s_n(z; b_1, b_2, b_3, a_1, a_2, a_3; q) &= \lim_{c \rightarrow 0} (-1)^n q^{-\binom{n}{2}} c^{2n} Q_n \left(\frac{(cz) + (cz)^{-1}}{2}; b_1 c, b_2 c, a_1/c, a_2/c, a_3/c, b_3 c; q \right) \\ &= \frac{(a_1 b_1, a_2 b_1)_n}{(b_1 b_3)^n} {}_4\phi_3 \left(\begin{matrix} q^{-n}, q/a_3 b_2, q^n/a_3 b_3, b_1/z \\ a_1 b_1, a_2 b_1, q/a_3 z \end{matrix}; q \right), \\ t_n(z; b_1, b_2, b_3, a_1, a_2, a_3; q) &= \lim_{c \rightarrow 0} Q_n \left(\frac{(cz) + (cz)^{-1}}{2}; a_1/c, b_1 c, b_2 c, a_2/c, b_3 c, a_3/c; q \right) \\ &= (q^{-1} a_2 b_3)^n (a_1 b_1, a_1 b_2)_n {}_4\phi_3 \left(\begin{matrix} q^{-n}, q/a_2 b_3, q^n/a_3 b_3, a_1 z \\ a_1 b_1, a_1 b_2, qz/b_3 \end{matrix}; q \right). \end{aligned}$$

Again, these are effectively five-parameter systems. Both s_n and t_n are symmetric in the pairs (a_1, a_2) and (b_1, b_2) . We obtain from (3.15) the biorthogonality

$$\mathbf{K}(s_m t_n) = \delta_{mn} (q^{-1} a_1 a_2)^n \frac{1 - q^{-1} a_1 a_2 b_1 b_2}{1 - q^{2n-1} a_1 a_2 b_1 b_2} \frac{(q)_n \prod_{i,j=1}^2 (a_i b_j)_n}{(q^{-1} a_1 a_2 b_1 b_2)_n}. \quad (3.31)$$

In this case, it is natural to use the expression (3.28) for \mathbf{K} , as the poles of s_m and t_n are then situated at zeroes of the biorthogonality measure, so that one may use the same contour of integration for all m and n . With the biorthogonality written in this form, (3.31) is due to Al-Salam and Ismail [1]. In the classification of [2], the system (s_n, t_n) is denoted by 2200 uv (in Appendix A.2 it appears with a typo as 2200 vp). The expressions (3.25) and (3.29) for \mathbf{K} generalize the biorthogonality measures found in [1] and [2].

4 Superconformal Indices

As was discussed in the introduction, when considering integral evaluations related to dualities for three-dimensional supersymmetric quantum field theories, (1.4) appears as a top level beta integral. As before, the integration is over a positively oriented contour separating the geometric sequences of poles tending to 0 from those tending to ∞ ; note that the contour necessarily depends on x . Besides the quantum field theory interpretation, (1.4) also appears as the star-triangle relation for a solvable lattice model introduced in [8].

Let us replace b_j in (1.4) by $b_j q^{k_j}$ and N_j by $N_j + l_j - k_j$, where k_j and l_j are integers with $k_1 + \cdots + k_6 = l_1 + \cdots + l_6 = 0$. After simplification, this results in the identity

$$\begin{aligned} & \sum_{x=-\infty}^{\infty} \oint \frac{(1 - q^x z^2)(1 - q^x z^{-2})}{q^x z^{6x}} \prod_{j=1}^6 \frac{(q^{1+x/2}/b_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{N_j+x/2} b_j z, q^{N_j-x/2} b_j/z)_{\infty}} \\ & \times \prod_{j=1}^6 (b_j (q^{-x/2} z)^{\pm})_{k_j} (b_j q^{N_j} (q^{x/2} z)^{\pm})_{l_j} \frac{dz}{2\pi i z} \\ & = \frac{2}{\prod_{j=1}^6 q^{\binom{k_j}{2} + \binom{l_j}{2} + \binom{N_j}{2} + N_j l_j} b_j^{k_j + l_j + N_j}} \\ & \quad \prod_{1 \leq i < j \leq 6} \frac{(q/b_i b_j)_{\infty} (b_i b_j)_{k_i + k_j} (b_i b_j q^{N_i + N_j})_{l_i + l_j}}{(b_i b_j q^{N_i + N_j})_{\infty}}. \end{aligned}$$

This can be interpreted in terms of the functional (3.7). Let V be the space introduced in Sect. 3.1 and V' the space obtained from V through replacing the parameters b_j with $b_j q^{N_j}$. Denoting the Rahman functional on V' by \mathbf{J}' , we obtain the following result.

Proposition 4.1 For $f \in V$ and $g \in V'$,

$$\begin{aligned} & \sum_{x=-\infty}^{\infty} \oint \frac{(1-q^x z^2)(1-q^x z^{-2})}{q^x z^{6x}} \prod_{j=1}^6 \frac{(q^{1+x/2}/b_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{N_j+x/2} b_j z, q^{N_j-x/2} b_j/z)_{\infty}} \\ & \times f\left(\frac{q^{-x/2} z + q^{x/2} z^{-1}}{2}\right) g\left(\frac{q^{x/2} z + q^{-x/2} z^{-1}}{2}\right) \frac{dz}{2\pi i z} \\ & = \frac{2}{\prod_{j=1}^6 q^{\binom{N_j}{2}} b_j^{N_j}} \prod_{1 \leq i < j \leq 6} \frac{(q/b_i b_j)_{\infty}}{(b_i b_j q^{N_i+N_j})_{\infty}} \mathbf{J}(f) \mathbf{J}'(g). \end{aligned} \quad (4.1)$$

We find it remarkable that the right-hand side of (4.1) factors, even though the variables of f and g on the left are coupled. We refer to this as the *decoupling phenomenon*. Just as for Spiridonov's beta integral discussed in the introduction, it can be used to construct two-index biorthogonal systems. To this end, let Q_n and R_n be as in (3.14). Moreover, if (c_1, \dots, c_6) is any permutation of $(b_1 q^{N_1}, \dots, b_6 q^{N_6})$, let

$$\begin{aligned} Q'_n &= Q_n \left(\frac{z + z^{-1}}{2}; c_1, c_2, c_3, c_4, c_5, c_6 \right), \\ R'_n &= Q_n \left(\frac{z + z^{-1}}{2}; c_1, c_2, c_3, c_4, c_6, c_5 \right). \end{aligned}$$

Combining (3.15) and Proposition 4.1 then gives the following result.

Corollary 4.2 In the notation above,

$$\begin{aligned} & \sum_{x=-\infty}^{\infty} \oint \frac{(1-q^x z^2)(1-q^x z^{-2})}{q^x z^{6x}} \prod_{j=1}^6 \frac{(q^{1+x/2}/b_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{N_j+x/2} b_j z, q^{N_j-x/2} b_j/z)_{\infty}} \\ & \times (Q_{n_1} R_{m_1}) \left(\frac{q^{-x/2} z + q^{x/2} z^{-1}}{2} \right) (Q'_{n_2} R'_{m_2}) \left(\frac{q^{x/2} z + q^{-x/2} z^{-1}}{2} \right) \frac{dz}{2\pi i z} \\ & = \delta_{n_1 m_1} \delta_{n_2 m_2} \frac{2}{\prod_{j=1}^6 q^{\binom{N_j}{2}} b_j^{N_j}} \prod_{1 \leq i < j \leq 6} \frac{(q/b_i b_j)_{\infty}}{(b_i b_j q^{N_i+N_j})_{\infty}} \\ & \times \frac{1 - 1/b_5 b_6}{1 - q^{2n_1}/b_5 b_6} \frac{(q)_{n_1} \prod_{1 \leq i < j \leq 4} (b_i b_j)_{n_1}}{q^{n_1} (1/b_5 b_6)_{n_1}} \frac{1 - 1/c_5 c_6}{1 - q^{2n_2}/c_5 c_6} \frac{(q)_{n_2} \prod_{1 \leq i < j \leq 4} (c_i c_j)_{n_2}}{q^{n_2} (1/c_5 c_6)_{n_2}}. \end{aligned}$$

Up to permutation of the parameters, there are three nonequivalent cases of Corollary 4.2, corresponding to whether none, one, or both of the numbers c_5, c_6 are contained in $\{b_5 q^{N_5}, b_6 q^{N_6}\}$.

At the Saalschütz level, we have the superconformal index identity [7]

$$\sum_{x=-\infty}^{\infty} \oint \prod_{j=1}^3 \frac{(q^{1+x/2}/a_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{M_j+x/2} a_j z, q^{N_j-x/2} b_j/z)_{\infty}} \left(-\frac{q^{1/2}}{tz^3} \right)^x \frac{dz}{2\pi iz} \\ = \frac{1}{\prod_{j=1}^3 q^{\binom{M_j}{2} + \binom{N_j}{2}} a_j^{M_j} b_j^{N_j}} \prod_{i,j=1}^3 \frac{(q/a_i b_j)_{\infty}}{(a_i b_j q^{M_i+N_j})_{\infty}}, \quad (4.2)$$

where the parameters satisfy (2.5), t is given by (2.7), and the integers M_j and N_j satisfy $M_1 + M_2 + M_3 = N_1 + N_2 + N_3 = 0$. In [7] this is stated under the additional assumption $t = q^{1/2}$, but the general case follows immediately after rescaling z and the parameters.

We replace the parameters in (4.2) by $a_j \mapsto a_j q^{k_j}$, $b_j \mapsto b_j q^{l_j}$, $M_j \mapsto M_j + m_j - k_j$, $N_j \mapsto N_j + n_j - l_j$, where the integral shifts satisfy

$$k_1 + k_2 + k_3 = -l_1 - l_2 - l_3 = m_1 + m_2 + m_3 = -n_1 - n_2 - n_3 = T.$$

The resulting identity can be written

$$\sum_{x=-\infty}^{\infty} \oint \prod_{j=1}^3 \frac{(q^{1+x/2}/a_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{M_j+x/2} a_j z, q^{N_j-x/2} b_j/z)_{\infty}} \left(-\frac{q^{1/2}}{tz^3} \right)^x \\ \times z^{-2T} \prod_{j=1}^3 (q^{-x/2} a_j z)_{k_j} (q^{x/2} b_j/z)_{l_j} (q^{M_j+x/2} a_j z)_{m_j} (q^{N_j-x/2} b_j/z)_{n_j} \frac{dz}{2\pi iz} \\ = \frac{q^{2\binom{T}{2}} (t^2 q^{M_1+M_2+M_3})^T}{\prod_{j=1}^3 q^{\binom{k_j}{2} + \binom{l_j}{2} + \binom{m_j}{2} + \binom{n_j}{2}} + \binom{M_j}{2} + \binom{N_j}{2} + m_j M_j + n_j N_j} a_j^{k_j+m_j+M_j} b_j^{l_j+n_j+N_j} \\ \times \prod_{i,j=1}^3 \frac{(q/a_i b_j)_{\infty} (a_i b_j)_{k_i+l_j} (a_i b_j q^{M_i+N_j})_{m_i+n_j}}{(a_i b_j q^{M_i+N_j})_{\infty}}.$$

If we now let \mathbf{K} and W be the functional and space of rational functions introduced in (3.19) and \mathbf{K}' , W' the objects obtained from these by replacing the parameters a_j by $a_j q^{M_j}$ and b_j by $b_j q^{N_j}$, we may draw the following conclusion.

Proposition 4.3 For $f \in W$ and $g \in W'$,

$$\sum_{x=-\infty}^{\infty} \oint \prod_{j=1}^3 \frac{(q^{1+x/2}/a_j z, q^{1-x/2} z/b_j)_{\infty}}{(q^{M_j+x/2} a_j z, q^{N_j-x/2} b_j/z)_{\infty}} \left(-\frac{q^{1/2}}{tz^3} \right)^x f(q^{-x/2} z) g(q^{x/2} z) \frac{dz}{2\pi iz} \\ = \frac{1}{\prod_{j=1}^3 q^{\binom{M_j}{2} + \binom{N_j}{2}} a_j^{M_j} b_j^{N_j}} \prod_{i,j=1}^3 \frac{(q/a_i b_j)_{\infty}}{(a_i b_j q^{M_i+N_j})_{\infty}} \mathbf{K}(f) \mathbf{K}'(g).$$

We can now use the results of Sect. 3.4 to construct corresponding biorthogonal systems. This can be done in many ways, since we may for each of the spaces W and W' work with the system (q_n, r_n) , $(\tilde{q}_n, \tilde{r}_n)$, or (s_n, t_n) . Moreover, permuting the parameters gives cases when both, one, or none of the two functions chosen for W have poles at the same geometric sequences as the functions chosen for W' . In total, this gives ten essentially nonequivalent possibilities. We refrain from writing these down explicitly.

Acknowledgements This work is dedicated to the memory of Mizan Rahman, a gentle and generous mathematician whose unsurpassed mastery of q -series has been a great inspiration for me. I thank Ilmar Gahramanov for fruitful discussions on superconformal indices and Slava Spiridonov for encouraging me to finish the present paper.

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (<http://creativecommons.org/licenses/by/4.0/>), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made.

References

1. Al-Salam, W.A., Ismail, M.E.H.: A q -beta integral on the unit circle and some biorthogonal rational functions. *Proc. Am. Math. Soc.* **121**, 553–561 (1994)
2. van de Bult, F.J., Rains, E.M.: Limits of elliptic hypergeometric biorthogonal functions. *J. Approx. Theory* **193**, 128–163 (2015)
3. Date, E., Jimbo, M., Kuniba, A., Miwa, T., Okado, M.: Exactly solvable SOS models II. Proof of the star-triangle relation and combinatorial identities. In: Jimbo, M., Miwa, T., Tsuchiya, A. (eds.) *Conformal Field Theory and Solvable Lattice Models*, pp. 17–122. Academic Press, Boston, MA (1988)
4. Dolan, F.A., Osborn, H.: Applications of the superconformal index for protected operators and q -hypergeometric identities to $N = 1$ dual theories. *Nucl. Phys. B* **818**, 137–178 (2009)
5. Gahramanov, I., Kels, A.P.: The star-triangle relation, lens partition function, and hypergeometric sum/integrals. *J. High Energy Phys.* **2017**, 40
6. Gahramanov, I., Rosengren, H.: A new pentagon identity for the tetrahedron index. *J. High. Energy Phys.* **2013**, 128 (2013)
7. Gahramanov, I., Rosengren, H.: Basic hypergeometry of supersymmetric dualities. *J. Nucl. Phys. B* **913**, 747–768 (2016)
8. Gahramanov, I., Spiridonov, V.P.: The star-triangle relation and 3d superconformal indices. *J. High Energy Phys.* **2015**, 40 (2015)
9. Gasper, G.: q -extensions of Barnes', Cauchy's, and Euler's beta integrals. In: Rassias, T.M. (eds.) *Topics in Mathematical Analysis*, pp. 294–314. World Scientific Publishing, Teaneck, NJ (1989)
10. Gasper, G., Rahman, M.: *Basic Hypergeometric Series*. Cambridge University Press, Cambridge (1990)
11. Hwang, C., Kim, H., Park, K.-J., Park, J.: Index computation for 3d Chern–Simons matter theory: test of Seiberg-like duality. *J. High Energy Phys.* **2011**, 37 (2011)
12. Imamura, Y., Yokoyama, S.: Index for three dimensional superconformal field theories with general R-charge assignments. *J. High Energy Phys.* **2011**, 7 (2011)
13. Jackson, M.: On well-poised bilateral hypergeometric series of the type ${}_8\psi_8$. *Q. J. Math.* **1**, 63–68 (1950)
14. Kapustin, A., Willett, B.: Generalized superconformal index for three dimensional field theories. [arXiv:1106.2484](https://arxiv.org/abs/1106.2484)
15. Kels, A.P.: A new solution of the star-triangle relation. *J. Phys. A* **47**, 055203 (2014)
16. Kels, A.P.: New solutions of the star-triangle relation with discrete and continuous spin variables. *J. Phys. A* **48**, 435201 (2015)
17. Kinney, J., Maldacena, J., Minwalla, S., Raju, S.: An index for 4 dimensional super conformal theories. *Commun. Math. Phys.* **275**, 209–254 (2007)

18. Koekoek, R., Swarttouw, R.F.: The Askey-Scheme of Hypergeometric Orthogonal Polynomials and Its q -Analogue, Delft University of Technology. <http://homepage.tudelft.nl/11r49/askey/> (1998)
19. Krattenthaler, C., Spiridonov, V.P., Vartanov, G.S.: Superconformal indices of three-dimensional theories related by mirror symmetry. *J. High Energy Phys.* **2011**, 8 (2011)
20. Rahman, M.: Families of biorthogonal rational functions in a discrete variable. *SIAM J. Math. Anal.* **12**, 355–367 (1981)
21. Rahman, M.: An integral representation of a ${}_{10}\phi_9$ and continuous bi-orthogonal ${}_{10}\phi_9$ rational functions. *Can. J. Math.* **38**, 605–618 (1986)
22. Rahman, M.: Some extensions of Askey–Wilson’s q -beta integral and the corresponding orthogonal systems. *Can. Math. Bull.* **31**, 467–476 (1988)
23. Rahman, M.: Biorthogonality of a system of rational functions with respect to a positive measure on $[-1, 1]$. *SIAM J. Math. Anal.* **22**, 1430–1441 (1991)
24. Rains, E.M.: Transformations of elliptic hypergeometric integrals. *Ann. Math.* **171**, 169–243 (2010)
25. Römelberger, C.: Counting chiral primaries in $N = 1, d = 4$ superconformal field theories. *Nucl. Phys. B* **747**, 329–353 (2006)
26. Rosengren, H.: An elementary approach to $6j$ -symbols (classical, quantum, rational, trigonometric, and elliptic). *Ramanujan J.* **13**, 133–168 (2007)
27. Rosengren, H.: Elliptic hypergeometric functions. [arXiv:1608.06161](https://arxiv.org/abs/1608.06161)
28. Rosengren, H., Schlosser, M.: Elliptic determinant evaluations and the Macdonald identities for affine root systems. *Compos. Math.* **142**, 937–961 (2006)
29. Ruijsenaars, S.N.M.: Complete integrability of relativistic Calogero–Moser systems and elliptic function identities. *Commun. Math. Phys.* **110**, 191–213 (1987)
30. Seiberg, N.: Electric-magnetic duality in supersymmetric non-abelian gauge theories. *Nuclear Phys. B* **435**, 129–146 (1995)
31. Spiridonov, V.P.: On the elliptic beta function. *Russ. Math. Surv.* **56**, 185–186 (2001)
32. Spiridonov, V.P.: Theta hypergeometric integrals. *St. Petersburg. Math. J.* **15**, 929–967 (2004)
33. Spiridonov, V.P.: Rarefied elliptic hypergeometric functions. [arXiv:1609.00715](https://arxiv.org/abs/1609.00715)
34. Spiridonov, V.P., Vartanov, G.S.: Elliptic hypergeometry of supersymmetric dualities. *Commun. Math. Phys.* **304**, 797–874 (2011)
35. Spiridonov, V.P., Vartanov, G.S.: Elliptic hypergeometry of supersymmetric dualities II. Orthogonal groups, knots, and vortices. *Commun. Math. Phys.* **325**, 421–486 (2014)
36. Spiridonov, V., Zhedanov, A.: Spectral transformation chains and some new biorthogonal rational functions. *Commun. Math. Phys.* **210**, 49–83 (2000)
37. Whittaker, E.T., Watson, G.N.: *A Course of Modern Analysis*. Cambridge University Press, Cambridge (1927)
38. Wilson, J.A.: *Hypergeometric Series, Recurrence Relations and Some New Orthogonal Functions*. University of Wisconsin, Madison (1978)
39. Wilson, J.A.: Orthogonal functions from gram determinants. *SIAM J. Math. Anal.* **22**, 1147–1155 (1991)
40. Yamazaki, M.: New integrable models from the gauge/YBE correspondence. *J. Stat. Phys.* **154**, 895–911 (2014)
41. Yokoyama, S.: Index for three dimensional superconformal field theories and its applications. *J. Phys. Conf. Ser.* **343**, 012134 (2012)