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Global Koppelman Formulas on (Singular) Projective Varieties

Mats Andersson¹ 

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Abstract

Let $i: X \rightarrow \mathbb{P}^N$ be a projective manifold of dimension n embedded in projective space \mathbb{P}^N , and let L be the pullback to X of the line bundle $\mathcal{O}_{\mathbb{P}^N}(1)$. We construct global explicit Koppelman formulas on X for smooth $(0, *)$ -forms with values in L^s for any s . The same construction works for singular, even non-reduced, X of pure dimension, if the sheaves of smooth forms are replaced by suitable sheaves \mathcal{A}_X^* of $(0, *)$ -currents with mild singularities at X_{sing} . In particular, if $s \geq \text{reg } X - 1$, where $\text{reg } X$ is the Castelnuovo–Mumford regularity, we get an explicit representation of the well-known vanishing of $H^{0,q}(X, L^{s-q})$, $q \geq 1$. Also some other applications are indicated.

Keywords Projective variety · Integral formula · dbar-equation

Mathematics Subject Classification 32A26 · 32C30 · 32J99

1 Introduction

During the last decade, global Koppelman formulas for $\bar{\partial}$ on various special projective varieties have been constructed, see, e.g., [17–19, 22–24]. The aim of this paper is to present a quite general explicit construction of intrinsic¹ Koppelman formulas on any projective, possibly non-reduced, subvariety $i: X \rightarrow \mathbb{P}^N$ of pure dimension n .

Let us first assume that X is smooth; even in this case, such global formulas are previously known only in case X is (locally) a complete intersection in \mathbb{P}^N . Let $L \rightarrow X$ be the restriction of the ample line bundle $\mathcal{O}_{\mathbb{P}^N}(1)$ to X , and let $\mathcal{E}_X^{0,q}(L^r)$ denote the

¹ Here “intrinsic” means that the operators in the formula only depend on intrinsic forms on X .

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✉ Mats Andersson
matsa@math.chalmers.se

¹ Department of Mathematics, Chalmers University of Technology and the University of Gothenburg, 412 96 Gothenburg, Sweden

sheaf of smooth $(0, q)$ -forms on X with values in L^r . We introduce integral operators $\mathcal{K}: \mathcal{E}^{0, q+1}(X, L^s) \rightarrow \mathcal{E}^{0, q}(X_{reg}, L^s)$ and $\mathcal{P}: \mathcal{E}^{0, q}(X, L^s) \rightarrow \mathcal{E}^{0, q}(X, L^s)$ such that the Koppelman formula

$$\phi(z) = \bar{\partial}\mathcal{K}\phi + \mathcal{K}(\bar{\partial}\phi) + \mathcal{P}\phi \quad (1.1)$$

holds on X . In some situations, see below, we can choose the operators so that $\mathcal{P}\phi = 0$; if $\bar{\partial}\phi = 0$, then $\psi = \mathcal{K}\phi$ is a smooth solution to $\bar{\partial}\psi = \phi$ on X . In certain cases, we can choose \mathcal{P} such that $\mathcal{P}: \mathcal{E}^{0, 0}(X, L^s) \rightarrow \mathcal{O}(\mathbb{P}^N, \mathcal{O}(s))$. Then $\mathcal{P}\phi$ is a holomorphic extension to \mathbb{P}^N of a holomorphic section ϕ of L^r on X . We get no new existence results; the novelty is that we have explicit formulas for the global solutions and for the holomorphic extensions. The operators \mathcal{K} are given by kernels $k(\zeta, z)$ that are defined on $X \times X$ and integrable in ζ for any $z \in X$. Simply speaking the operators locally behave like standard integral operators for $\bar{\partial}$ in \mathbb{C}^n ; in particular, they extend to L^p -spaces, etc., and all classical local norm estimates hold.

Let us now turn our attention to the case when $i: X \rightarrow \mathbb{P}^N$ is a subvariety of pure dimension n . In case X is reduced, there are well-known definitions of smooth forms and currents on X , cf. Sect. 3. In [9] are introduced reasonable definitions of sheaves $\mathcal{E}_X^{0, *}$ of smooth $(0, *)$ -forms and suitable sheaves of currents even in the non-reduced case, see Sect. 6 below. In [6] and [9] are introduced, by means of local Koppelman formulas, fine sheaves \mathcal{A}_X^k (in fact, modules over $\mathcal{E}_X^{0, *}$) of $(0, q)$ -currents on X , or any pure-dimensional analytic space, with the following properties: There are sheaf inclusions $\mathcal{E}_X^{0, q} \subset \mathcal{A}_X^q$ with equality on X_{reg} , whereas \mathcal{A}_X^q have “mild” singularities at X_{sing} , and

$$0 \rightarrow \mathcal{O} \rightarrow \mathcal{A}_X^0 \xrightarrow{\bar{\partial}} \mathcal{A}_X^1 \xrightarrow{\bar{\partial}} \dots \quad (1.2)$$

is a (fine) resolution of the structure sheaf \mathcal{O}_X of holomorphic functions on X . By the abstract de Rham theorem, we therefore have canonical isomorphisms

$$H^k(X, L^s) \simeq \frac{\text{Ker}(\mathcal{A}^k(X, L^s) \xrightarrow{\bar{\partial}} \mathcal{A}^{k+1}(X, L^s))}{\text{Im}(\mathcal{A}^{k-1}(X, L^s) \xrightarrow{\bar{\partial}} \mathcal{A}^k(X, L^s))}, \quad k \geq 1. \quad (1.3)$$

In this paper, we construct integral operators

$$\mathcal{K}: \mathcal{A}^{q+1}(X, L^s) \rightarrow \mathcal{A}^q(X, L^s), \quad \mathcal{P}: \mathcal{A}^q(X, L^s) \rightarrow \mathcal{E}^{0, q}(X, L^s)$$

such that again the Koppelman formula (1.1) holds on X . In the reduced case, the operators \mathcal{K} are given by kernels $k(\zeta, z)$ that are defined on $X_{reg} \times X$, locally integrable in ζ for $z \in X_{reg}$, \mathcal{P} are given by kernels $p(\zeta, z)$ that are smooth on $X_{reg} \times X$, and the integrals

$$\mathcal{K}\phi(z) = \int_X k(\zeta, z) \wedge \phi(\zeta), \quad z \in X_{reg}, \quad \mathcal{P}\phi(z) = \int_X p(\zeta, z) \wedge \phi(\zeta), \quad z \in X,$$

exist as principal values at X_{sing} . For the non-reduced case, see Sect. 6. As in the smooth case, in good situations $\mathcal{P}\phi$ vanishes so we get explicit solutions to $\bar{\partial}\psi = \phi$, and extensions of holomorphic sections from X to \mathbb{P}^N .

Remark 1 It was proved already in [21] that if X is a local reduced complete intersection and ϕ is a smooth $\bar{\partial}$ -closed form, then there is locally a smooth solution to $\bar{\partial}\psi = \phi$ on X_{reg} . The case with a general X was proved only in [5]. The analogous result for a local non-reduced space is a special case of the main result in [9]. It is known that in general there is no solution that is smooth across X_{sing} , see, e.g., [6, Example 1.1].

Let J_X be the homogeneous ideal in the graded ring $S = \mathbb{C}[z_0, \dots, z_N]$ associated with X . Let $S(-r)$ be the module S but with the grading shifted by r . There is a free graded resolution

$$0 \rightarrow M_{N_0} \xrightarrow{a_{N_0}} \dots \xrightarrow{a_3} M_2 \xrightarrow{a_2} M_1 \xrightarrow{a_1} M_0 \quad (1.4)$$

of the homogeneous module $M_J := S/J_X$; i.e.,

$$M_0 = S, \quad M_k = S(-d_k^1) \oplus \dots \oplus S(-d_k^{r_k}),$$

$a_k = (a_k^{ij})$ are matrices of homogeneous forms with

$$\deg a_k^{ij} = d_k^j - d_{k-1}^i \quad (1.5)$$

(1.4) is exact, and the cokernel of the right-most mapping is precisely S/J_X . Since 0 is not an associated prime ideal of J_X it follows from [11, Corollary 20.14] that one can choose (1.4) such that $N_0 \leq N$. Our integral formulas are explicitly constructed out of a resolution (1.4).

Recall that the (Castelnuovo–Mumford) *regularity* of X is defined as the regularity of the ideal J_X which turns out to be 1 plus the regularity of the module S/J_X , so that

$$\operatorname{reg} X = \max_{k,i} (d_k^i - k) + 1,$$

if (1.4) is a minimal free resolution of S/J_X , cf. [12, Ch. 4]. It is well known, see, e.g., [12, Proposition 4.16], that

$$H^q(X, L^{s-q}) = 0, \quad s \geq \operatorname{reg} X - 1, \quad q \geq 1, \quad (1.6)$$

and that the natural mapping

$$\mathcal{O}(\mathbb{P}^N, \mathcal{O}(s)) \rightarrow \mathcal{O}(X, L^s) \quad (1.7)$$

is surjective for $s \geq \operatorname{reg} X - 1$.

In [4, Example 3.4] is described an extension operator that provides an explicit proof of surjectivity of (1.7), in case X is reduced. The non-reduced case is obtained in precisely the same way following the ideas in [9]. By appropriate choices of operators \mathcal{K} we can give an explicit proof of the vanishing of (1.6), provided that X is irreducible. That is, we have

Theorem 1.1 *Let X be a, possibly non-reduced, irreducible, subvariety of \mathbb{P}^N of pure dimension n and assume that $s \geq \operatorname{reg} X - 1$. For each $q \geq 1$ there is an integral operator $\mathcal{K}: \mathcal{A}^q(X, L^{s-q}) \rightarrow \mathcal{A}^{q-1}(X, L^{s-q})$ such that $\bar{\partial}\mathcal{K}\phi = \phi$ if $\phi \in \mathcal{A}^q(X, L^{s-q})$ and $\bar{\partial}\phi = 0$.*

In fact, for fixed q it is enough that $s \geq \max_{\ell \leq N-q} d_\ell^i - (N - q)$; this follows from the proof below. When X is not irreducible a slightly less sharp version of the theorem still holds, see Proposition 7.2.

Koppelman formulas on \mathbb{P}^n were found by Götmark, [17], and on more general symmetric spaces in [18]. In [20], explicit formulas for the $\bar{\partial}$ -equation are used on a smooth Riemann surface embedded in \mathbb{P}^2 , for L^1 -estimates. Similar formulas were also introduced in [19], cf. Sect. 8.2 below. Koppelman formulas for global, even non-reduced, complete intersections are constructed in the recent papers [23,24], cf. Sect. 8 below.

As already mentioned, the main novelty in this paper is Koppelman formulas for an arbitrary embedded projective variety X . We think that these formulas will be of interest even when X is smooth. We prove Theorem 1.1 as an illustration of the utility and indicate some other applications in Sect. 7.1. We hope that our Koppelman formulas will be useful for other purposes as well.

In Sect. 2 we describe, based on [3,4,17], how one can obtain weighted integral formulas on \mathbb{P}^N . We need some elements from residue theory that we have collected in Sect. 3. In Sects. 4 and 5, we then describe the construction of our Koppelman formulas on a pure-dimensional subvariety. In order to keep the technicalities on a reasonable level, we restrict here to the case where X is reduced. The reader who is mainly interested in the smooth case can just think that X_{sing} is empty, in that way avoiding a lot of technicalities. In Sect. 6, we discuss the non-reduced case.

Remark 2 The local study of the $\bar{\partial}$ -equation on non-smooth spaces by L^2 -methods was initiated by Pardon and Stern, [31–33], and has been developed by a number of authors since then, notably Fornaess, Gavosto, Ovrelid, Ruppenthal, Vassiliadou, see., e.g., [14–16,28–30,34,35].

The equation has also been studied by local and semi-global integral formulas in, e.g., [5,6,21,25,26,37]. \square

2 Integral Representation on \mathbb{P}^N

We first describe how one can generate weighted Koppelman formulas on \mathbb{P}^N for sections of a holomorphic vector bundle $F \rightarrow \mathbb{P}^N$. This is an adaption of an idea from [1] to \mathbb{P}^N , following [3,4]; see also [17].

Let $\pi: \mathbb{C}_z^{N+1} \setminus \{0\} \rightarrow \mathbb{P}_z^N$ be the natural projection and let $\mathcal{U} \subset \mathbb{P}^N$ be an open set. Recall that a form ξ in $\pi^{-1}\mathcal{U} \subset \mathbb{C}_z^{N+1} \setminus \{0\}$ is projective, i.e., the pullback of a form ξ in \mathcal{U} , if and only if ξ is homogeneous and $\delta_z \xi = \delta_{\bar{z}} \xi = 0$, where δ_z and $\delta_{\bar{z}}$ are interior multiplication by $\sum_0^N z_j (\partial/\partial z_j)$ and its conjugate, respectively. We will identify forms in \mathcal{U} by projective forms in $\pi^{-1}\mathcal{U}$.

Let $\mathcal{O}_z(k)$ denote the pullback of $\mathcal{O}(k) \rightarrow \mathbb{P}_z^N$ to $\mathbb{P}_z^N \times \mathbb{P}_{\bar{z}}^N$ under the projection $\mathbb{P}_z^N \times \mathbb{P}_{\bar{z}}^N \rightarrow \mathbb{P}_z^N$ and define $\mathcal{O}_{\bar{z}}(k)$ analogously.

Throughout this paper we will only consider forms and currents that only contain holomorphic differentials with respect to ζ , whereas anti-holomorphic differentials with respect to both z and ζ may occur.

Notice that

$$\eta = 2\pi i \sum_0^N z_i \frac{\partial}{\partial \zeta_i},$$

is a section of $\mathcal{O}_z(1) \otimes \mathcal{O}_\zeta(-1) \otimes T_{1,0}(\mathbb{P}_\zeta^N)$ on $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N$. Contraction δ_η by η defines a mapping

$$\delta_\eta : \mathcal{C}^{\ell+1,q}(\mathcal{O}_\zeta(k) \otimes \mathcal{O}_z(j)) \rightarrow \mathcal{C}^{\ell,q}(\mathcal{O}_\zeta(k-1) \otimes \mathcal{O}_z(j+1)),$$

where $\mathcal{C}^{\ell,q}(\mathcal{O}_\zeta(k) \otimes \mathcal{O}_z(j))$ denotes the sheaf of currents of bidegree (ℓ, q) that take values in $\mathcal{O}_\zeta(k) \otimes \mathcal{O}_z(j)$. Notice that δ_η only affects holomorphic differentials with respect to ζ . Given a vector bundle $L \rightarrow \mathbb{P}_\zeta^n \times \mathbb{P}_z^n$, let

$$\mathcal{L}^v(L) = \bigoplus_j \mathcal{C}^{j,j+v}(\mathcal{O}_\zeta(j) \otimes \mathcal{O}_z(-j) \otimes L).$$

If

$$\nabla_\eta = \delta_\eta - \bar{\partial},$$

then $\nabla_\eta : \mathcal{L}^v(L) \rightarrow \mathcal{L}^{v+1}(L)$ and $\nabla_\eta^2 = 0$. Furthermore, if L' is a line bundle and ϕ, ψ are sections of $\mathcal{L}^v(L)$ and $\mathcal{L}^{v'}(L')$, respectively, then $\phi \wedge \psi$ is a section of $\mathcal{L}^{v+v'}(L \otimes L')$, and

$$\nabla_\eta(\phi \wedge \psi) = \nabla_\eta \phi \wedge \psi + (-1)^{\deg \phi} \phi \wedge \nabla_\eta \psi.$$

Notice that

$$b = \frac{|\zeta|^2 \bar{z} \cdot d\zeta - (\bar{z} \cdot \zeta) \bar{\zeta} \cdot d\zeta}{|\zeta|^2 |z|^2 - |\bar{\zeta} \cdot z|^2}$$

is a $(1, 0)$ -form on $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N \setminus \Delta$ with values in $\mathcal{O}_\zeta(1) \otimes \mathcal{O}_z(-1)$ such that $\delta_\eta b = 1$; here Δ is the diagonal in $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N$.

Lemma 2.1 *The form*

$$B = \frac{b}{\nabla_\eta b} = \frac{b}{1 - \bar{\partial} b} = b + b \wedge \bar{\partial} b + \dots + b \wedge (\bar{\partial} b)^{N-1}$$

is an integrable section of \mathcal{L}^{-1} and

$$\nabla_\eta B = 1 - [\Delta]_{d\zeta}, \quad (2.1)$$

where the last term is the component of the current of integration $[\Delta]$ that has full degree N in $d\zeta$.

Proof Let us consider the affinization where $\zeta_0 \neq 0$. In the affine coordinates $\zeta'_j = \zeta_j/\zeta_0$, $j = 1, \dots, N$, and the frame z_0/ζ_0 for $\mathcal{O}_z(1) \otimes \mathcal{O}_\zeta(-1)$, we have

$$\eta = 2\pi i \sum_1^N (\zeta'_j - z'_j) \frac{\partial}{\partial \zeta'_j}.$$

It is readily checked that $|b| \leq C|\zeta' - z'|$ and $|\delta_\eta b| \geq C|\zeta' - z'|^2$, and therefore (2.1) follows, cf. [1, Example 4]. \square

Given a vector bundle $F \rightarrow \mathbb{P}^N$, let F_z denote the pullback of F to $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N$ under the natural projection $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N \rightarrow \mathbb{P}_z^N$ and define F_ζ analogously. A weight with respect to F is a smooth section g of $\mathcal{L}^0(\text{Hom}(F_\zeta, F_z))$ such that $\nabla_\eta g = 0$ and $g_0 = I_F$ on the diagonal in $\mathbb{P}_\zeta^N \times \mathbb{P}_z^N$, where g_0 denotes the term in g with bidegree $(0, 0)$. In general, we let lower index on a form denote degree with respect to holomorphic differentials of ζ . Notice that if g is a weight with respect to F , then from (2.1) we get

$$\nabla_\eta(g \wedge B) = (g - [\Delta]_{d\zeta})I_F. \quad (2.2)$$

Identifying terms of full degree in $d\zeta$ thus

$$\bar{\partial}(g \wedge B)_N = ([\Delta]_{d\zeta} - g_N)I_F.$$

By Stokes' theorem we get the following Koppelman formula, cf. [17] and [18].

Proposition 2.2 *Let g be a weight with respect to $F \rightarrow \mathbb{P}^N$. Then for $\phi \in \mathcal{E}^{0,q}(\mathbb{P}^N, F \otimes \mathcal{O}(-N))$ we have*

$$\phi(z) = \bar{\partial}_z \int_\zeta (g \wedge B)_N \wedge \phi + \int_\zeta (g \wedge B)_N \wedge \bar{\partial} \phi + \int_\zeta g_N \wedge \phi. \quad (2.3)$$

Example 1 It is easy to check that

$$\alpha = \alpha_{0,0} + \alpha_{1,1} = \frac{z \cdot \bar{\zeta}}{|\zeta|^2} - \bar{\partial} \frac{\bar{\zeta} \cdot d\zeta}{2\pi i |\zeta|^2}$$

is a projective form in $\mathcal{L}^0(\mathcal{O}_\zeta(-1) \oplus \mathcal{O}_z(1))$ such that

$$\nabla_\eta \alpha = 0. \quad (2.4)$$

Since α_0 is equal to $I_{\mathcal{O}(1)}$ on the diagonal, α is a weight with respect to $F = \mathcal{O}(1)$. For each natural number ρ therefore $g = \alpha^\rho$ is a weight with respect to $F = \mathcal{O}(\rho)$. For $\ell \geq -N$, we thus have the Koppelman formula

$$\phi(z) = \bar{\partial}_z \int_\zeta (\alpha^{\ell+N} \wedge B)_N \wedge \phi + \int_\zeta (\alpha^{\ell+N} \wedge B)_N \wedge \bar{\partial} \phi + \int_\zeta (\alpha^{N+\ell})_N \wedge \phi. \quad (2.5)$$

Since α is holomorphic in z and has no differentials $d\bar{z}$, the last term in (2.5) vanishes if $q \geq 1$, and for degree reasons also if $q = 0$ and $\ell \leq -1$.

We thus have an explicit proof of the well-known vanishing $H^{0,q}(\mathbb{P}^N, \mathcal{O}(\ell)) = 0$ for $\ell \geq -N$ if $1 \leq q \leq N$, and for $q = 0$ if $\ell \leq -1$.

Remark 3 Using the transposed integral operators in Example 1, we get an explicit proof of the vanishing $H^{N,q}(\mathbb{P}^N, \mathcal{O}(\ell)) = 0$ for $\ell \leq N$, $0 \leq 1 \leq N - 1$. Since $K_{\mathbb{P}^N} = \mathcal{O}(-N - 1)$ we therefore get that $H^{0,q}(\mathbb{P}^N, \mathcal{O}(\ell)) = 0$ for $\ell \leq -1$, $0 \leq q \leq N - 1$.

Let us now consider a weight with respect to $\mathcal{O}(-1)$. Let

$$\tau = \frac{d\bar{z} \cdot d\zeta}{2\pi i |z|^2} = \frac{\sum_0^N d\bar{z}_j \wedge d\zeta_j}{2\pi i |z|^2}.$$

Then

$$\beta = 2\pi i \left[\delta_\zeta \frac{\delta_{\bar{z}} \tau}{1 - \tau} \right] = 2\pi i \left[\frac{\delta_\zeta \delta_{\bar{z}} \tau}{1 - \tau} + \frac{\delta_\zeta \tau \wedge \delta_{\bar{z}} \tau}{(1 - \tau)^2} \right]$$

is a projective form. More explicitly,

$$\beta = \frac{\bar{z} \cdot \zeta}{|z|^2} + \frac{\bar{z} \cdot \zeta}{|z|^2} \frac{d\bar{z} \cdot d\zeta}{2\pi i |z|^2} + \frac{\bar{z} \cdot d\zeta \wedge \zeta \cdot d\bar{z}}{2\pi i |z|^4} + \cdots, \quad (2.6)$$

so each term has the same degree in $d\zeta$ as in $d\bar{z}$, and is holomorphic in ζ .

Proposition 2.3 *The form β is a weight with respect to $\mathcal{O}(-1)$.*

Proof Clearly, cf. (2.6), $\beta_{0,0} = I_{\mathcal{O}(-1)}$ on the diagonal. We claim that

$$\nabla_\eta \frac{\delta_{\bar{z}} \tau}{1 - \tau} = 1. \quad (2.7)$$

Since δ_ζ anti-commutes with ∇_η and $\delta_\zeta 1 = 0$ the proposition follows. To see (2.7), notice that

$$\delta_\eta \tau = -\frac{z \cdot d\bar{z}}{|z|^2}, \quad \bar{\partial} \tau = \delta_\eta \tau \wedge \tau, \quad \delta_\eta \delta_{\bar{z}} \tau = 1,$$

so that

$$\bar{\partial} \delta_{\bar{z}} \tau = \tau + \delta_\eta \tau \wedge \delta_{\bar{z}} \tau = \tau + \gamma.$$

Therefore,

$$\nabla_\eta \delta_{\bar{z}} \tau = 1 - \tau - \gamma, \quad (2.8)$$

and

$$\nabla_\eta \tau = \delta_\eta \tau \wedge (1 - \tau), \quad (2.9)$$

so that

$$\nabla_{\eta} \tau \wedge \delta_{\bar{z}} \tau = (1 - \tau) \wedge \gamma. \quad (2.10)$$

It follows that

$$\nabla_{\eta} \frac{\delta_{\bar{z}} \tau}{1 - \tau} = \frac{1 - \tau - \gamma}{1 - \tau} + \frac{\nabla_{\eta} \tau \wedge \delta_{\bar{z}} \tau}{(1 - \tau)^2} = 1,$$

in view of (2.8), (2.9), and (2.10). \square

Thus for $(0, q)$ -forms ϕ with values in $\mathcal{O}(\ell)$, $\ell \leq -N$, we get from Proposition 2.2 the Koppelman formula

$$\phi(z) = \bar{\partial} \int_{\zeta} (B \wedge \beta^{-N-\ell})_N \wedge \phi + \int_{\zeta} (B \wedge \beta^{-N-\ell})_N \wedge \bar{\partial} \phi + \int_{\zeta} (\beta^{-N-\ell})_N \wedge \phi.$$

For degree reasons the last term vanishes if $0 \leq q \leq N - 1$, so we get back the well-known vanishing $H^{0,q}(\mathbb{P}^N, \mathcal{O}(\ell)) = 0$ for $\ell \leq -N$. In case $q = N$, the obstruction term vanishes when $\ell = -N$, and when $\ell \leq -N - 1$ it vanishes if and only if

$$\int \psi \wedge \phi = 0 \quad (2.11)$$

for each holomorphic $(N, 0)$ -form with values in $\mathcal{O}(-\ell)$. That is, $\bar{\partial} v = \phi$ is solvable if and only if (2.11) holds. Of course this is precisely what we get by considering the transposed operators with the weight α , cf. Remark 3. However, in the non-smooth case, we have no obvious canonical bundle so we cannot consider transposed operators in the same simple way; therefore this weight β will play a role.

For future reference we prove

Proposition 2.4 *The forms*

$$\gamma_j = \delta_{\zeta} \left[\frac{\delta_{\bar{z}} \tau}{1 - \tau} \wedge d\zeta_j \right]$$

are projective and

$$\nabla_{\eta} \gamma_j = \beta z_j - \zeta_j. \quad (2.12)$$

Proof Clearly $\delta_{\zeta} \gamma_j = 0 = \delta_{\bar{z}} \gamma_j$ and thus γ_j is a projective form. By (2.7) we have that

$$\nabla_{\eta} \left[\frac{\delta_{\bar{z}} \tau}{1 - \tau} \wedge d\zeta_j \right] = d\zeta_j - 2\pi i \frac{\delta_{\bar{z}} \tau}{1 - \tau} z_j.$$

Since ∇_{η} and δ_{ζ} anti-commute, the proposition follows. \square

3 Some Preliminaries

Let X be any reduced analytic space of pure dimension n . By definition there is, locally, some embedding $i: X \rightarrow \Omega \subset \mathbb{C}^N$. Let $\mathcal{I}_X \subset \mathcal{O}_\Omega$ be the ideal sheaf of holomorphic functions in Ω that vanish on X . Then the sheaf of holomorphic functions on X , the structure sheaf \mathcal{O}_X , is represented as $\mathcal{O}_X = \mathcal{O}_\Omega / \mathcal{I}_X$. If Φ is a smooth form in Ω we say that Φ is in $\text{Ker } i^*$ if $i^*\Phi$ vanishes on X_{reg} . We define the sheaf

$$\mathcal{E}_X^{p,*} = \mathcal{E}_\Omega^{p,*} / \text{Ker } i^*.$$

of smooth forms on X , and have a natural mapping $i^*: \mathcal{E}_\Omega^{p,*} \rightarrow \mathcal{E}_X^{p,*}$. One can prove that $\mathcal{E}_X^{p,*}$ so defined is independent of the choice of embedding and is thus an intrinsic sheaf on X . We define the sheaf $\mathcal{C}_X^{p,*}$ of currents as the dual of $\mathcal{E}_X^{n-p,n-*}$. More concretely this means that currents τ in $\mathcal{C}_X^{p,*}$ are identified with currents $i_*\tau$ in $\mathcal{C}_\Omega^{p+N-n,*+N-n}$ such that $i_*\tau$ vanish on $\text{Ker } i^*$ so that $\tau \cdot i^*\Phi = i_*\tau \cdot \Phi$ for test forms Φ . Clearly $\bar{\partial}$ is defined on smooth forms and extends to currents by duality. Also the wedge product $\phi \wedge \tau$ is well defined as long as at least one of the factors is smooth. Thus the currents form a module over the smooth forms.

We say that a current in \mathbb{C}_s^M of the form

$$\frac{\gamma}{s_1^{a_1} \cdots s_r^{a_r}} \bar{\partial} \frac{1}{s_{r+1}^{a_{r+1}}} \wedge \cdots \bar{\partial} \frac{1}{s_{r'}^{a_{r'}}},$$

where γ is a test form, is *elementary*. A current τ on X is *pseudomeromorphic* if locally it is a finite sum of direct images under holomorphic mappings of elementary currents; see, e.g., [10] for a precise definition and basic properties. The pseudomeromorphic currents form a sheaf \mathcal{PM}_X that is closed under multiplication by $\mathcal{E}_X^{p,*}$ and the action of $\bar{\partial}$. Given a pseudomeromorphic current τ in an open set \mathcal{U} and a subvariety $V \subset \mathcal{U}$, the natural restriction of τ to $\mathcal{U} \setminus V$ has a canonical extension to a pseudomeromorphic current $\mathbf{1}_{\mathcal{U} \setminus V} \tau$ such that

$$\mathbf{1}_V \tau := \tau - \mathbf{1}_{\mathcal{U} \setminus V} \tau$$

has support on V . If ξ is a smooth form, then

$$\xi \wedge \mathbf{1}_V \tau = \mathbf{1}_V \xi \wedge \tau. \quad (3.1)$$

Let χ be a smooth function on $[0, \infty)$ that is 0 in a neighborhood of 0 and 1 in a neighborhood of ∞ and let h be a tuple of holomorphic functions, or a section of some holomorphic Hermitian vector bundle such that the zero set of h is precisely V . Then

$$\mathbf{1}_{\mathcal{U} \setminus V} \tau = \lim_{\delta \rightarrow 0} \chi(|h|/\delta) \tau. \quad (3.2)$$

We say that a current a in X is *almost semi-meromorphic*, $a \in \text{ASM}(X)$, if there is a smooth modification $\pi: X' \rightarrow X$, a generically non-vanishing holomorphic section σ of a line bundle $L \rightarrow X'$ and a smooth L -valued form γ such that

$$a = \pi_*(\gamma/\sigma).$$

Let $ZSS(a)$ be the smallest analytic subset of X such that a is smooth in $X \setminus ZSS(a)$. It follows that $ZSS(a)$ has positive codimension. Clearly an almost semi-meromorphic a is pseudomeromorphic.

Proposition 3.1 [10, Theorem 4.8] *Given any pseudomeromorphic τ and $a \in ASM(X)$ the current $a \wedge \tau$ a priori defined in $X \setminus ZSS(a)$ has a unique pseudomeromorphic extension to a pseudomeromorphic current in X , also denoted $a \wedge \tau$, such that $\mathbf{1}_{ZSS(a)} a \wedge \tau = 0$.*

Pseudomeromorphic currents have some important geometric properties, see, e.g., [10]:

Proposition 3.2 *Assume that the pseudomeromorphic current τ has support on a germ of an analytic variety V .*

- (i) *If the holomorphic function h vanishes on V , then $\bar{h}\tau = 0$ and $d\bar{h} \wedge \tau = 0$.*
- (ii) *If τ has bidegree $(*, p)$ and V has codimension $\geq p + 1$, then $\tau = 0$.*

We refer to (ii) as the *dimension principle*.

4 A Structure Form Associated to X

Let $i: X \rightarrow \mathbb{P}^N$ be a reduced subvariety of pure dimension n , and let (1.4) be a free graded resolution of the S -module $M_X = S/J_X$. In particular, then $a_1 = (a_1^{11}, \dots, a_1^{1r_1})$ is a tuple of homogeneous forms that define the homogeneous ideal J_X in the graded ring $S = \mathbb{C}[z_0, \dots, z_N]$. Let E_k^j be disjoint trivial line bundles over \mathbb{P}^N with basis elements $e_{k,j}$ and let

$$E_k = (E_k^1 \otimes \mathcal{O}(-d_k^1)) \oplus \dots \oplus (E_k^{r_k} \otimes \mathcal{O}(-d_k^{r_k})), \quad E_0 \simeq \mathbb{C}.$$

Then

$$0 \rightarrow E_{N_0} \xrightarrow{a_{N_0}} \dots \xrightarrow{a_3} E_2 \xrightarrow{a_2} E_1 \xrightarrow{a_1} E_0 \rightarrow 0 \quad (4.1)$$

is a complex of vector bundles over \mathbb{P}^N that is pointwise exact outside Z , and the corresponding complex of locally free sheaves

$$0 \rightarrow \mathcal{O}(E_{N_0}) \xrightarrow{a_{N_0}} \dots \xrightarrow{a_3} \mathcal{O}(E_2) \xrightarrow{a_2} \mathcal{O}(E_1) \xrightarrow{a_1} \mathcal{O}(E_0) \quad (4.2)$$

over \mathbb{P}^N is a resolution of the sheaf $\mathcal{O}(E_0)/\mathcal{I}_X$, where $\mathcal{I}_X \subset \mathcal{O}_{\mathbb{P}^N}$ is the ideal sheaf associated with X . See, e.g., [7, Sect. 6]. We equip E_k with the natural Hermitian metric

$$|\xi(z)|_{E_k}^2 = \sum_{j=1}^{r_k} |\xi_j(z)|^2 |z|^{2d_k^j}$$

if $\xi = (\xi_1, \dots, \xi_{r_k})$, so that (4.1) becomes a Hermitian complex. In [7] were introduced pseudomeromorphic currents

$$U = U_1 + \dots + U_{N_0}, \quad R = R_1 + \dots + R_{N_0}$$

on \mathbb{P}^N associated to (4.1) with the following properties: The currents U_k are almost semi-meromorphic $(0, k-1)$ -currents, smooth outside X , that take values in $\text{Hom}(E_0, E_k) \simeq E_k$, and R_k are $(0, k)$ -currents with support on X , taking values in $\text{Hom}(E_0, E_k) \simeq E_k$. Moreover, we have the relations

$$a_1 U_1 = I_{E_0}, \quad a_{k+1} U_{k+1} - \bar{\partial} U_k = -R_k, \quad k \geq 1, \quad (4.3)$$

which can be compactly written as

$$\nabla_a U = I_{E_0} - R \quad (4.4)$$

if

$$\nabla_a = a - \bar{\partial} = a_1 + a_2 + \dots + a_{N_0} - \bar{\partial}.$$

If Φ is a section of $\mathcal{O} = \mathcal{O}(E_0)$, then the current $R\Phi$ vanishes if and only if Φ is in \mathcal{J}_X , see [7, Theorem 1.1].

Let X_k be the analytic subset of \mathbb{P}^N where a_k does not have optimal rank. Then

$$\dots X_{k+1} \subset X_k \dots X_{N-n+1} \subset X_{\text{sing}} \subset X = X_{N-n} = \dots = X_1. \quad (4.5)$$

Since \mathcal{J}_X has pure dimension

$$\text{codim } X_k \geq k+1, \quad k \geq N-n+1, \quad (4.6)$$

see [11, Corollary 20.14].

By the dimension principle $R_k = 0$ for $k < N-n$. Moreover, there are almost semi-meromorphic $\text{Hom}(E_k, E_{k+1})$ -valued $(0, 1)$ -currents α_{k+1} , smooth outside X_{k+1} , such that

$$R_{k+1} = \alpha_{k+1} R_k$$

there. By (4.6) and the dimension principle it follows that this equality must hold across X_{k+1} if the right-hand side is interpreted in the sense of Proposition 3.1. By a simple induction argument, using (4.6) and the dimension principle, it follows that

$$\mathbf{1}_{X_{\text{sing}}} R = 0. \quad (4.7)$$

Lemma 4.1 *If Φ is a smooth $(0, *)$ -form, then $i^* \Phi = 0$ on X_{reg} if and only if $\Phi R = 0$.*

It follows that $R\phi$ is well-defined for ϕ in $\mathcal{E}_X^{0,*}$.

Proof Locally at a point $x \in X_{reg}$ we can choose coordinates (z, w) such that $X = \{w = 0\}$. By a Taylor expansion of Φ in w , using that $w_j R = \bar{w}_j R = d\bar{w}_j \wedge R = 0$, cf. Proposition 3.2 (i), we find that $R\Phi = 0$ if and only if $i^*\Phi = 0$. If $\Phi R = 0$ on X_{reg} it follows from (3.1) and (4.7) that $\Phi R = 0$ identically. \square

Let $\Omega = \delta_z(dz_0 \wedge \cdots \wedge dz_N)$ be the unique, up to a multiplicative constant, non-vanishing global $(N, 0)$ -form with values in $\mathcal{O}(N+1)$. From [6, Proposition 3.3] we have

Proposition 4.2 *There is a unique almost semi-meromorphic current $\omega = \omega_0 + \cdots + \omega_n$ on X that is smooth on X_{reg} , ω_ℓ have bidegree (n, ℓ) and take values in $E^\ell := i^*E_{N-n+\ell}$, and*

$$i_*\omega = \Omega \wedge R, \quad i_*\omega_\ell = \Omega \wedge R_{N-n+\ell}. \quad (4.8)$$

We say that ω is a *structure form* on X .

For any smooth form ξ on \mathbb{P}^N there is a unique form $\vartheta(\xi)$ such that

$$\vartheta(\xi) \wedge \Omega = \xi_N, \quad (4.9)$$

where ξ_N denotes the components of ξ of bidegree $(N, *)$. From (4.8) and (4.9) we have that

$$\xi_N \wedge R = \vartheta(\xi) \wedge \Omega \wedge R = i_*(\vartheta(\xi) \wedge \omega), \quad (4.10)$$

where we in the last term, for simplicity, write $\vartheta(\xi) \wedge \omega$ rather than $i^*\vartheta(\xi) \wedge \omega$.

Lemma 4.3 *Let $\chi(t)$ be a smooth function as in (3.2). If h is a holomorphic section of a Hermitian vector bundle that does not vanish identically on any irreducible component of X and $\chi_\delta = \chi(|h|/\delta)$, then*

$$\chi_\delta \omega \rightarrow \omega, \quad \bar{\partial} \chi_\delta \wedge \omega \rightarrow 0. \quad (4.11)$$

Proof Let W be the zero set of h . Notice that $i_*\mathbf{1}_{X_{sing}}\omega = \mathbf{1}_{X_{sing}}i_*\omega = \mathbf{1}_{X_{sing}}\Omega \wedge R = \Omega \wedge \mathbf{1}_{X_{sing}}R = 0$ in view of Lemma 4.1. Thus $\mathbf{1}_{X_{sing}}\omega = 0$, and hence $\mathbf{1}_W\mathbf{1}_{X_{sing}}\omega = 0$. Since ω is smooth on X_{reg} we have that $\mathbf{1}_W\mathbf{1}_{X_{reg}}\omega = 0$. By simple computational rules, see, e.g., [10], we conclude that $\mathbf{1}_W\omega = 0$. In view of (3.2) thus the first part of (4.11) follows. Notice that (4.4) implies that $(a - \bar{\partial})R = 0$, and by (4.8) thus $(a - \bar{\partial})\omega = 0$. Applying $(a - \bar{\partial})$ to the first limit in (4.11) now the second one follows. \square

5 Koppelman Formulas on a Projective Variety

Let $U^\lambda = |a_1|^{2\lambda}U$ and $R^\lambda = 1 - |a_1|^{2\lambda} + \bar{\partial}|a_1|^{2\lambda} \wedge U$. Then R^λ and U^λ are as smooth as we may wish if $\operatorname{Re} \lambda$ is sufficiently large. In particular, R^λ and U^λ are well-defined currents. Moreover, they admit analytic continuations to $\operatorname{Re} \lambda > -\epsilon$, and the values at $\lambda = 0$ are precisely R and U , respectively, see [7].

Let ρ be an integer. Following [3, Definition 1] we say that $H = (H_k^\ell)$ is a Hefer morphism for the complex $E_\bullet \otimes \mathcal{O}(\rho)$, a , cf. (4.1), if H_k^ℓ are smooth sections of

$$\mathcal{L}^{-k+\ell}(\mathrm{Hom}(E_{\zeta,k} \otimes \mathcal{O}_\zeta(\rho), E_{z,\ell} \otimes \mathcal{O}_z(\rho))),$$

$H_k^\ell = 0$ for $k < \ell$, the term $(H_\ell^\ell)_0$ of bidegree $(0, 0)$ is the identity I_{E_ℓ} on Δ , and

$$\nabla_\eta H_k^\ell = H_{k-1}^\ell a_k - a_{\ell+1}(z) H_k^{\ell+1}, \quad (5.1)$$

where a_k stands for $a_k(\zeta)$. From [4, Lemma 2.5] we have

Lemma 5.1 *Assume that H is a Hefer morphism for the complex $E_\bullet \otimes \mathcal{O}(\rho)$, a . For $\mathrm{Re} \lambda \gg 0$, the form*

$$g^\lambda = a_1(z) H^1 U^\lambda + H^0 R^\lambda \quad (5.2)$$

is a weight with respect to $\mathcal{O}(\rho)$.

$$\text{Here } H^1 U = \sum_j H_j^1 U_j \text{ and } H^0 R = \sum_j H_j^0 R_j.$$

Proposition 5.2 *Let F be holomorphic vector bundle over \mathbb{P}^N and let g be a weight with respect to F . Moreover, assume that H is a Hefer morphism for $E_\bullet \otimes L^\rho$. For $\phi \in \mathcal{E}^{0,q}(X, F \otimes L^{\rho-N})$ we have the Koppelman formula*

$$\begin{aligned} \phi(z) = & \bar{\partial}_z \int_\zeta (H R \wedge g \wedge B)_N \wedge \phi + \int_\zeta (H R \wedge g \wedge B)_N \wedge \bar{\partial} \phi \\ & + \int_\zeta (H R \wedge g)_N \wedge \phi, \quad z \in X_{\mathrm{reg}}. \end{aligned} \quad (5.3)$$

By blowing up $\mathbb{P}^N \times \mathbb{P}^N$ along the diagonal one can verify that B is almost semi-meromorphic. In view of Proposition 3.1 thus $(H R \wedge g \wedge B)_N \wedge \phi$ is a well-defined pseudomeromorphic current on $\mathbb{P}^N \times \mathbb{P}^N$. Formally (5.3) means that

$$\phi = \bar{\partial} \pi_*((H R \wedge g \wedge B)_N \wedge \phi) + \pi_*((H R \wedge g \wedge B)_N \wedge \bar{\partial} \phi) + \pi_*((H R \wedge g)_N \wedge \phi),$$

where π is the projection $(\zeta, z) \mapsto z$.

Proof Since $g^\lambda \wedge g$ is a weight with respect to $\mathcal{O}(\rho) \otimes F$, and $a_1(z) = 0$ when $z \in X$, from Proposition 2.2 we get (5.3) with R^λ instead of R . In view of (the proof of) [6, Lemma 5.2], see also [9, Lemma 9.5], we can take $\lambda = 0$ and so we get the proposition, keeping in mind that the product $B \wedge \tau$ can be defined as the value of $|\eta|^{2\lambda} B \wedge \tau$ at $\lambda = 0$ in view of [6, (2.2) and (2.3)]. \square

Let

$$\mathcal{K}\phi = \pi_*((H R \wedge g \wedge B)_N \wedge \phi), \quad \mathcal{P}\phi = \pi_*((H R \wedge g)_N \wedge \phi). \quad (5.4)$$

Then we can write (5.3) as

$$\phi = \bar{\partial} \mathcal{K}\phi + \mathcal{K} \bar{\partial} \phi + \mathcal{P}\phi \quad (5.5)$$

for $z \in X_{reg}$.

In view of (4.10) we have that

$$\mathcal{K}\phi(z) = \int_X k(\zeta, z) \wedge \phi(\zeta), \quad \mathcal{P}\phi(z) = \int_X p(\zeta, z) \wedge \phi(\zeta), \quad (5.6)$$

where

$$k = \pm \vartheta(g \wedge B \wedge H) \wedge \omega, \quad p = \pm \vartheta(g \wedge H) \wedge \omega. \quad (5.7)$$

It is apparent from (5.6) that \mathcal{K} and \mathcal{P} are intrinsic integral operators on X . Locally they are precisely of the type in [6], so it follows that $\mathcal{K}\phi$ is smooth on X_{reg} if ϕ is smooth. Moreover, from [6, Theorem 1.4] we get:

Theorem 5.3 *Let F be holomorphic vector bundle over \mathbb{P}^N and let g be a weight with respect to F on X . Moreover, assume that H is a Hefer morphism for $E_\bullet \otimes \mathcal{O}(\rho)$ and \mathcal{K} and \mathcal{P} are defined by (5.6). Then*

$$\begin{aligned} \mathcal{K}: \mathcal{A}^{k+1}(X, E \otimes L^{\rho-N}) &\rightarrow \mathcal{A}^k(X, E \otimes L^{\rho-N}), \\ \mathcal{P}: \mathcal{A}^k(X, E \otimes L^{\rho-N}) &\rightarrow \mathcal{E}^{0,k}(\mathbb{P}^N, E \otimes \mathcal{O}(\rho - N)) \end{aligned}$$

and the global Koppelman formula (5.5) holds on X for $\phi \in \mathcal{A}^q(X, F \otimes L^{\rho-N})$.

6 The Non-reduced Case

Now assume that $i: X \rightarrow \mathbb{P}^N$ has pure dimension n but is non-reduced. Then we still have an ideal sheaf $\mathcal{I}_X \subset \mathcal{O}_{\mathbb{P}^N}$ that has pure dimension n but \mathcal{I}_X is no longer radical, i.e., there are nilpotent elements. Still the structure sheaf of X has the representation $\mathcal{O}_X = \mathcal{O}_{\mathbb{P}^N} / \mathcal{I}_X$. The underlying reduced space X_{red} is associated with the radical ideal $\sqrt{\mathcal{I}_X} = \mathcal{I}_{X_{red}}$.

Let X_{reg} be the subset of X where X_{red} is smooth and in addition \mathcal{I}_X is Cohen–Macaulay. In a neighborhood \mathcal{U} of a point x in X_{reg} , which is an open dense subset of X , we can choose local coordinates (z, w) such that $X_{red} \cap \mathcal{U} = \{w = 0\}$. It turns out, see, e.g., [9], that there are monomials $1, w^{\alpha_1}, \dots, w^{v-1}$ such that each ϕ in \mathcal{O}_X has a unique representation

$$\phi = \hat{\phi}_0(z) \otimes 1 + \dots + \hat{\phi}_{v-1}(z) \otimes w^{v-1}. \quad (6.1)$$

Thus \mathcal{O}_X has the structure of a free $\mathcal{O}_{X_{red}}$ -module in \mathcal{U} .

We say that Φ in $\mathcal{E}_{\mathbb{P}^N}^{0,*}$ is in $\text{Ker } i^*$ if in a neighborhood of each point in X_{reg} , Φ is in the subsheaf of $\mathcal{E}_{\mathbb{P}^N}^{0,*}$ generated by $\mathcal{I}_X, \bar{\mathcal{I}}_{X_{red}}$, and $d\bar{\mathcal{I}}_{X_{red}}$. As in the reduced case we define $\mathcal{E}_X^{0,*} = \mathcal{E}_{\mathbb{P}^N}^{0,*} / \text{Ker } i^*$, and again it is independent of the choice of local embedding of X . It turns out that at each point in X_{reg} and coordinates (z, w) as above we have a unique representation (6.1) of ϕ in $\mathcal{E}_X^{0,*}$ where $\hat{\phi}_j$ are in $\mathcal{E}_{X_{red}}^{0,*}$.

We define the sheaf of $(n, *)$ -currents on X as the dual of $\mathcal{E}_X^{0, n-*}$, so that such a current τ is represented by a $(N, N - n + *)$ -current $i_*\tau$ in \mathbb{P}^N that is annihilated by $\text{Ker } i^*$.

Basically all facts in Sect. 4 now hold verbatim, except for that one has to replace X by X_{red} in (4.5) and slightly modify the proof of Lemma 4.1. It follows from Lemma 4.1 that there is a current ω such that (4.8) holds. However, in the non-reduced case we give no meaning to that ω is almost semi-meromorphic and smooth on X_{reg} . The first part of (4.11) just means that $\mathbf{1}_W R = 0$ and this follows from [9, Corollary 6.3]. The second part of (4.11) follows from the first part precisely as before.

Following [9] we can also make the construction of Koppelman formulas in Sect. 5 and define sheaves \mathcal{A}^* on X so that Theorem 5.3 holds.

Remark 4 Recall that a holomorphic differential operator L is Noetherian with respect to the ideal \mathcal{J} if $L\phi$ vanishes on Z if $\phi \in \mathcal{J}$. As in the local case, cf. [6, Remark 6.6], there is a tuple \mathcal{L} of global Noetherian operators on \mathbb{P}^N with almost semi-meromorphic coefficients so that

$$\omega.\xi = \int_Z \mathcal{L}\xi,$$

cf. [8, Theorem 4.1 and Proposition 5.1].

7 Global Solutions

To begin with we consider a Hefer morphism, introduced in [3], for the complex $E_\bullet \otimes \mathcal{O}(\rho)$, a for large ρ . Let E' denote the complex of trivial bundles over \mathbb{C}^{N+1} that we get from E , and let A denote the corresponding mappings (which then formally are just the original matrices a). Let δ_{w-z} denote interior multiplication by

$$2\pi i \sum_0^N (w_j - z_j) \frac{\partial}{\partial w_j}.$$

in $\mathbb{C}_w^{N+1} \times \mathbb{C}_z^{N+1}$.

Proposition 7.1 *There exist $(k - \ell, 0)$ -form-valued mappings*

$$h_k^\ell = \sum_{ij} (h_k^\ell)_{ij} e_{\ell i} \otimes e_{kj}^* : \mathbb{C}_w^{N+1} \times \mathbb{C}_z^{N+1} \rightarrow \text{Hom}(E'_k, E'_\ell),$$

such that $h_k^\ell = 0$ for $k < \ell$, $h_\ell^\ell = I_{E'_\ell}$,

$$\delta_{z-w} h_k^\ell = h_{k-1}^\ell A_k(w) - A_{\ell+1}(z) h_k^{\ell+1}, \quad (7.1)$$

and the coefficients in the form $(h_k^\ell)_{ij}$ are homogeneous polynomials of degree $d_k^j - d_\ell^i - (k - \ell)$.

Notice that

$$\gamma_j = d\zeta_j - \frac{\bar{\zeta} \cdot d\zeta}{|\zeta|^2} \zeta_j \quad (7.2)$$

is a projective form and that

$$\nabla_\eta \gamma_j = 2\pi i (z_j - \alpha \zeta_j). \quad (7.3)$$

Given h_k^ℓ in Proposition 7.1 we let $\tau^* h_k^\ell$ be the projective form we obtain by replacing w by $\alpha \zeta$ and dw_j by γ_j . We then have

$$\nabla_\eta \tau^* h = \tau^* (\delta_{w-z} h), \quad (7.4)$$

in light of (7.3) and (2.4). It is proved in [3] that if

$$\kappa_0 = \kappa_0(X) = \max d_k^i,$$

then

$$H_k^\ell = \sum_{ij} (\tau^* h_k^\ell)_{ij} \wedge \alpha^{\kappa_0 - d_k^j} e_{\ell,i} \otimes e_{k,j}^*,$$

is a Hefer morphism for (4.1) with $\rho = \kappa_0$. Clearly, H is holomorphic in z .

Recall that $g = \alpha^v$ is a holomorphic weight with respect to $F = \mathcal{O}(v)$ for $v \geq 0$. From Proposition 5.2 we thus obtain explicit solutions to the $\bar{\partial}$ -equation in L^ℓ for $\ell \geq \kappa_0(X) - N$.

Proposition 7.2 *Assume that X is a possibly singular projective subvariety of \mathbb{P}^N of pure dimension n , and $s \geq \kappa_0 - N$. If ϕ is a $\bar{\partial}$ -closed section in $\mathcal{A}^q(X, L^s)$, $q \geq 1$, then*

$$\mathcal{K}\phi(z) = \int (H_{\kappa_0} R \wedge \alpha^{s - \kappa_0 + N} \wedge B)_N \wedge \phi$$

is a solution in $\mathcal{A}^{q-1}(X, L^s)$ to $\bar{\partial}u = \phi$.

Notice that $\kappa_0 - N \leq \max(d_k^i - k) \leq \text{reg } X - 1$. Thus the proposition gives a weaker form of Theorem 1.1.

Remark 5 The associated operator \mathcal{P} is precisely the operator in [4, Example 3.4] that realizes the surjectivity of (1.7). That is, if $\phi \in \mathcal{O}(X, L^s)$, then

$$\phi(z) = \int (H R \wedge \alpha^{s - \kappa_0 + N})_N \wedge \phi$$

is a global section of $\mathcal{O}(s) \rightarrow \mathbb{P}^N$ that coincides with ϕ on X .

Example 2 Assume that X is a complete intersection, i.e., J is generated by homogeneous forms a_1^1, \dots, a_1^p , of degrees d^1, \dots, d^p , where $p = N - n$. Then the Koszul complex generated by a_1^j provides a minimal free resolution, and it is then easy to see that $\kappa_0 = d^1 + \dots + d^p$, cf. Sect. 8 below. Moreover, by the adjunction formula

$$K_X = K_{\mathbb{P}^N} \otimes \mathcal{O}(d^1 + \dots + d^p) = L^{\kappa_0 - N - 1}.$$

Here K_X is the Grothendieck dualizing sheaf, which in this case is a line bundle that is generated by $\omega = \omega_0$. When X is smooth K_X is just the usual canonical bundle. If we define (n, q) -forms as $(0, q)$ -forms with values in K_X , then Proposition 7.2 gives an explicit realization of the vanishing

$$H^{n,q}(X, L^\ell) = 0, \quad 1 \leq q \leq n, \quad \ell \geq 1. \quad (7.5)$$

If X is smooth this follows precisely from Kodaira's theorem, since L is strictly positive on X .

For the proof of Theorem 1.1 we must make a more careful analysis of the kernels. Let us introduce the notation

$$\kappa_q = \kappa_q(X) = \max_{\ell \leq N-q} d_\ell^i. \quad (7.6)$$

Notice that

$$\kappa_q - N \leq \max_{\ell \leq N-q} (d_\ell^i - (N - \ell) - \ell) \leq \text{reg } X - 1 - q. \quad (7.7)$$

Proof of Theorem 1.1 Notice that the section $h(\zeta, z) = \zeta \cdot \bar{z}/|z|^2$ of $\mathcal{O}_\zeta(1) \otimes \mathcal{O}_z(-1)$ is non-vanishing on Δ . Let $\chi(t)$ be a cutoff function as before and let

$$\chi_\delta := \chi(|h|^2/\delta) = \chi(|\bar{\zeta} \cdot z|^2/|z|^2|\zeta|^2\delta).$$

Here $|h|$ denote the natural norm of the section h , whereas in the last term $||$ denotes norm of points in \mathbb{C}^{N+1} . For small δ , χ_δ is identically 1 in a neighborhood of Δ and thus

$$g^\delta := \chi_\delta - \bar{\partial}\chi_\delta \wedge B \quad (7.8)$$

is a smooth weight (with respect to the trivial line bundle). For fixed z , g^δ vanishes in a neighborhood of the hyperplane $h = 0$, and therefore

$$\alpha^{-r} \wedge g^\delta$$

is a smooth weight with respect to $\mathcal{O}(-r)$ for any r , though not holomorphic in z . Now fix $q \geq 1$ and let $t = s - q$. In particular, for $t \geq \kappa_q - N$ and $\phi \in \mathcal{A}^q(X, L^t)$, $q \geq 1$, with $\bar{\partial}\phi = 0$ we have the formula

$$\phi(z) = \bar{\partial} \int_\zeta (H R \wedge \alpha^{t-\kappa_0+N} \wedge g^\delta \wedge B)_N \wedge \phi$$

$$+ \int_{\zeta} (H R \wedge \alpha^{t-\kappa_0+N} \wedge g^{\delta} \wedge \phi)_N \wedge \phi = \bar{\partial} \mathcal{K}^{\delta} \phi + \mathcal{P}^{\delta} \phi. \quad (7.9)$$

We claim that $\mathcal{K}^{\delta} \phi$ tends to a current $\mathcal{K} \phi$ in $\mathcal{A}^{q-1}(X, L^t)$ and that $\mathcal{P}^{\delta} \phi \rightarrow 0$. Taking this for granted, the theorem follows in view of (7.7).

To settle the claim we first consider the expression for $\mathcal{P}^{\delta} \phi$ in (7.9). Since ϕ has bidegree $(0, q)$ only components $H_{N-\ell} R_{N-\ell}$ with $\ell \geq q$ can occur in the integral. Thus the total power of α is

$$\kappa_0 - d_{N-\ell}^i + t - \kappa_0 + N \geq \kappa_0 - d_{N-\ell}^i + \kappa_q - N - \kappa_0 + N = \kappa_q - d_{N-\ell}^i \geq 0$$

in view of (7.6). Thus

$$\mathcal{P}^{\delta} \phi(z) = \sum_{\ell \geq q} \int_X \omega_{n-\ell} \wedge \phi_{0,q} \wedge \vartheta(\xi_{\ell} \wedge (\chi_{\delta} - \bar{\partial} \chi_{\delta} \wedge B)),$$

where ξ_{ℓ} are smooth and holomorphic in z . Since $q \geq 1$ we need some anti-holomorphic differentials with respect to z and they must come from $\bar{\partial} \chi_{\delta} \wedge B$; hence we can forget about χ_{δ} . Since $\chi_{\delta} = 1$ in a neighborhood of the diagonal, we can consider B as smooth. Thus we have to verify that

$$\omega \wedge \phi \wedge \bar{\partial} \chi_{\delta} \rightarrow 0, \quad \delta \rightarrow 0. \quad (7.10)$$

Since X is irreducible, $h = 0$ has positive codimension on X , and if ϕ is smooth thus (7.10) holds in view of Lemma 4.3. If ϕ is in \mathcal{A}^q , then it is in Dom_X , cf. [6, 9], and then (7.10) follows from (the proofs of) [6, Lemma 4.1] and [9, Lemma 8.4]. In fact, (7.10) can be reformulated as $\mathbf{1}_{h=0} \bar{\partial}(\omega \wedge \phi) = 0$. We conclude that $\mathcal{P}^{\delta} \phi \rightarrow 0$. Notice now that $B \wedge B = 0$ so that

$$\mathcal{K}^{\delta} \phi = \int_{\zeta} \chi_{\delta} (H R \wedge \alpha^{t-\kappa_0+N} \wedge B)_N \wedge \phi.$$

It is proved in [6, 9] that $(H R \wedge \alpha^{t-\kappa_0+N} \wedge B)_N \wedge \phi$ is in the space $\mathcal{W}^{X \times \mathbb{P}^N}$, and this implies that

$$\chi_{\delta} (H R \wedge \alpha^{t-\kappa_0+N} \wedge B)_N \wedge \phi \rightarrow (H R \wedge \alpha^{t-\kappa_0+N} \wedge B)_N \wedge \phi.$$

It follows that

$$\mathcal{K}^{\delta} \phi \rightarrow \mathcal{K} \phi = \int_{\zeta} (H R \wedge \alpha^{t-\kappa_0+N} \wedge B)_N \wedge \phi.$$

□

7.1 Examples with Negative Curvature

We now turn our attention to the case of negative curvature. We define a Hefer morphism from h_k^ℓ in Proposition 7.1 by replacing w by ζ , z by βz , and dw_j by the γ_j from Proposition 2.4. The morphism so obtained is a Hefer morphism for (4.1) (i.e., for $E_\bullet \otimes \mathcal{O}(\rho)$, a with $\rho = 0$). This is verified in the same way as [3, Proposition 4.4].

This time H is not holomorphic in z but in ζ instead. Let δ be the depth of the ring S/J . This is a number, $0 \leq \delta \leq n$, and choosing (1.4) minimal, (4.1) will end up at $k = N - \delta$, which means that $R = R_{N-n} + \cdots + R_{N-\delta}$. The variety X is Cohen–Macaulay precisely when $\delta = n$.

From the Koppelman formula we get solutions to $\bar{\partial}$ (representation of the cohomology in the smooth case) for $(0, q)$ -forms ϕ with values in $\mathcal{O}(\ell)$ for $\ell \leq -N$ and thus solutions as soon as the obstruction term

$$\int (HR \wedge \beta^{-N-\ell})_N \wedge \phi \quad (7.11)$$

vanishes. Notice that HR has degree at most $N - \delta$ in $d\bar{\zeta}$ since β and γ_j only contain holomorphic differentials with respect to ζ . Therefore (7.11) must vanish if $N - \delta + q < N$, i.e., $0 \leq q \leq \delta - 1$.

Theorem 7.3 *Assume that X is a subvariety of \mathbb{P}^N of pure dimension n and $\ell \leq -N$. Then for any $\bar{\partial}$ -closed $(0, q)$ -form $\phi \in \mathcal{A}_q(X, L^\ell)$, $0 < q \leq \delta - 1$,*

$$\psi(z) = \int (HR \wedge \beta^{-N-\ell} \wedge B)_N \wedge \phi$$

is a solution in $\mathcal{A}^{q-1}(X, L^\ell)$ to $\bar{\partial}\psi = \phi$.

We thus have an explicit proof of the vanishing $H^{0,q}(X, \mathcal{O}(\ell)) = 0$ for $0 \leq q \leq \delta - 1$, $\ell \leq -N$.

8 Global Complete Intersections

Let us compute the resulting formulas in case $i: X \rightarrow \mathbb{P}^N$ is a global complete intersection as in Example 2. Let f_1, \dots, f_p , $f_j = a_1^j$, be our given homogeneous forms of degrees d^j from Example 2, and recall that $p = N - n$. Assume that E_0 is the trivial line bundle and let

$$E = E^1 \otimes \mathcal{O}(-d^1) \oplus \cdots \oplus E^p \otimes \mathcal{O}(-d^p),$$

where E^j are trivial line bundles. Let e_j be basis elements for E^j and let e_j^* be the dual basis elements. We take

$$E_k = \Lambda^k E = \sum_{|I|=k}^p \mathcal{O}(-(d^{I_1} + \cdots + d^{I_k})) E^{I_1} \otimes \cdots \otimes E^{I_k}$$

and $a_k: E_k \rightarrow E_{k-1}$ as interior multiplication by $f = \sum f_j e_j^*$. Now

$$\sigma = \sum_j \frac{\overline{f_j(z)}}{|z|^{2d_j}} e_j / \|f\|^2$$

is the section of E with minimal norm such that $f \cdot \sigma = 1$ outside Z , if $\|f\| = |f|_{E^*}$. Moreover,

$$U = \|f\|^{2\lambda} \sum_{k=1}^m \sigma \wedge (\bar{\partial}\sigma)^{k-1} \Big|_{\lambda=0}$$

and

$$R = 1 - \|f\|^{2\lambda} + \bar{\partial}\|f\|^{2\lambda} \wedge \sum_{k=1}^m \sigma \wedge (\bar{\partial}\sigma)^{k-1} \Big|_{\lambda=0}, \quad (8.1)$$

cf. Sect. 5 and [2]; here $|_{\lambda=0}$ means evaluation at $\lambda = 0$ after analytic continuation.

Since $\text{codim } Z = p$ the resulting residue current R just consists of the term R_p ; it coincides with the classical Coleff–Herrera product

$$\bar{\partial} \frac{1}{f_p} \wedge \cdots \wedge \bar{\partial} \frac{1}{f_1} \wedge e_1 \wedge \cdots \wedge e_p.$$

We now compute Hefer morphisms for the Koszul complex. Let $\tilde{h}_j(w, z)$ be $(1, 0)$ -forms in $\mathbb{C}^{n+1} \times \mathbb{C}^{n+1}$ of polynomial degrees $d^j - 1$ such that

$$\delta_{w-z} \tilde{h}_j = f_j(w) - f_j(z)$$

and let $h_j = \tau^* \tilde{h}_j$. We only have to care about $k \leq p$ so $\kappa_0 = d^1 + \cdots + d^p$. Then

$$H_k^\ell = \sum_{|I|=\ell} \sum_{|J|=k-\ell} \pm h_{J_1} \wedge \cdots \wedge h_{J_{k-\ell}} \wedge e_I \otimes e_{IJ}^* \wedge \alpha^{\kappa_0 - (d^{J_1} + \cdots + d^{J_{k-\ell}} + d^{I_1} + \cdots + d^{I_k})}$$

is a Hefer morphism. The components of most interest for us are H_k^0 and H_k^1 . Since

$$H_k^0 = \sum_{|J|=k} \pm h_{J_1} \wedge \cdots \wedge h_{J_k} \wedge e_J^* \wedge \alpha^{\kappa_0 - (d^{J_1} + \cdots + d^{J_k})}$$

it can be more compactly written formally as

$$H_k^0 = \alpha^{\kappa_0} \wedge (\delta_h)_k,$$

where δ_h denotes formal interior multiplication with

$$h = \sum \alpha^{-d^j} \wedge h_j \wedge e_j^*$$

and $(\delta_h)_k = (\delta_h)^k/k!$. In the same way

$$H_k^1 = \alpha^{\kappa_0} \wedge N(\delta_h)_{k-1},$$

where

$$N = \sum_j \alpha^{-d^j} e_j \otimes e_j^*.$$

Our description of U , H_k^1 , etc., is just to illustrate what these currents look like in the complete intersection case since they play a role in the proofs above. As we have seen, however, in the final Koppelman formula only the term

$$H_p^0 R_p = h_1 \wedge \cdots \wedge h_p \wedge \bar{\partial} \frac{1}{f_p} \wedge \cdots \wedge \bar{\partial} \frac{1}{f_1}$$

of HR occurs. It follows that the operator \mathcal{K} in Proposition 7.2, with

$$\kappa = s + N - (d^1 + \cdots + d^p) \quad (8.2)$$

has the more explicit form

$$\mathcal{K}\phi(z) = \int (\alpha^\kappa \wedge B \wedge h_1 \wedge \cdots \wedge h_p)_N \wedge \bar{\partial} \frac{1}{f_p} \wedge \cdots \wedge \bar{\partial} \frac{1}{f_1} \wedge \phi \quad (8.3)$$

and the operator \mathcal{P} in Remark 5 is

$$\mathcal{P}\phi(z) = \int (\alpha^\kappa \wedge h_1 \wedge \cdots \wedge h_p)_N \wedge \bar{\partial} \frac{1}{f_p} \wedge \cdots \wedge \bar{\partial} \frac{1}{f_1} \wedge \phi. \quad (8.4)$$

Proposition 8.1 *Assume that the projective space $i: X \rightarrow \mathbb{P}^N$ of codimension p is defined by the homogeneous forms f_j on \mathbb{C}^{N+1} of degree d^j , $j = 1, \dots, p$, and assume that $s \geq d^1 + \cdots + d^p - N$. For $\phi \in \mathcal{E}^{0,k}(X, L^s)$, or $\phi \in \mathcal{A}^k(X, L^s)$, we have the Koppelman formula (5.5) with \mathcal{K} and \mathcal{P} defined by (8.3) and (8.4), respectively. Moreover, \mathbb{P} vanishes if $k \geq 1$.*

Remark 6 In [23] similar Koppelman formulas are obtained on a, not necessarily reduced, global complete intersection X for $(0, *)$ -forms with values in $L^{d^1 + \cdots + d^p - N - 1}$ in the notation from Example 2. The authors construct Koppelman formulas on homogeneous subvarieties of \mathbb{C}^{N+1} , keep track of homogeneities and so obtain Koppelman formulas on X . They use the same definition of $\bar{\partial}$ as we do. However, they only consider solutions to $\bar{\partial}u = \phi$ on X_{reg} when ϕ is smooth on X and satisfies a condition $(*)$ that in general is stronger than $\bar{\partial}\phi = 0$. There is no discussion whether their solution has some meaning as a current across X_{sing} . The condition $(*)$ on ϕ means that (locally) there is a smooth extension Φ to ambient space such that $\bar{\partial}\Phi$ is in $\mathcal{E}\mathcal{J}$. Clearly this implies that $\bar{\partial}\phi = 0$ on X but in general the converse

does not hold. In fact, consider a reduced hypersurface $X = \{a = 0\} \subset \mathbb{C}^{n+1}$ so that $\mathcal{J} = (a)$. Then $(*)$ means that there is an extension Φ such that $\bar{\partial}\Phi = \xi a$ for some smooth form ξ . Then $0 = a\bar{\partial}\xi$ and thus $\bar{\partial}\xi = 0$; hence $\xi = \bar{\partial}\eta$ for some smooth η . Now $\Phi - a\eta$ is $\bar{\partial}$ -closed and therefore there is a smooth solution to $\bar{\partial}\Psi = \Phi - a\eta$. It follows that $\psi = i^*\Psi$ is a smooth solution to $\bar{\partial}\psi = \phi$. However, it is well known that there are smooth ϕ with $\bar{\partial}\phi = 0$ such that $\bar{\partial}\psi = \phi$ has no smooth solution, see, e.g., [6, Example 1.1].

8.1 The Reduced Case

Let us consider a more intrinsic-looking representation of \mathcal{K} and \mathcal{P} as in (5.6) and (5.7). In order to avoid a Noetherian operator, cf. Remark 4, let us in addition assume that X is reduced. Let A_1, \dots, A_p be holomorphic vector fields on \mathbb{C}^{N+1} such that

$$\delta_{A_p} \cdots \delta_{A_1} df_1 \wedge \cdots \wedge df_p = (2\pi i)^p,$$

or equivalently,

$$df_1 \wedge \cdots \wedge df_p \wedge \delta_{A_p} \cdots \delta_{A_1} (d\zeta_0 \wedge \cdots \wedge d\zeta_N) = (2\pi i)^p d\zeta_0 \wedge \cdots \wedge d\zeta_N. \quad (8.5)$$

Notice that since δ_ζ anti-commutes with δ_{A_j} ,

$$\omega' := \delta_{A_p} \cdots \delta_{A_1} \Omega \quad (8.6)$$

is a projective form. Following the proof of [4, Proposition 6.3] we see that ω' is a representative for the structure form on X , that is,

$$\omega = \omega_0 = i^* \omega'. \quad (8.7)$$

Example 3 Let

$$\delta_A = \delta_{A_1} \cdots \delta_{A_p}, \quad \bar{\partial} \frac{1}{f} = \bar{\partial} \frac{1}{f_p} \wedge \cdots \wedge \bar{\partial} \frac{1}{f_1}.$$

Then

$$\delta_A \xi_N = \delta_A (\vartheta(\xi) \wedge \Omega) = \pm \vartheta(\xi) \wedge \omega'. \quad (8.8)$$

In view of (5.7) we thus have that

$$\mathcal{K}\phi = \pm \int_X \delta_A (\alpha^\kappa \wedge B \wedge h)_N \wedge \phi, \quad \mathcal{P}\phi = \pm \int_X \delta_A (\alpha^\kappa \wedge h)_N \wedge \phi$$

for $\phi \in \mathcal{A}^q(X, L^s)$, $s = \kappa + d^1 + \cdots + d^p - N$, $\kappa \geq 0$.

8.2 Explicit Formulas for a Curve in \mathbb{P}^N

Following [19] we will now describe how one can find an especially simple expression for the kernel k when X is a curve. Applying δ_η to

$$(\alpha^\kappa \wedge B \wedge h)_N = \vartheta(\alpha^\kappa \wedge B \wedge h) \wedge \Omega,$$

cf. (4.9), we get

$$\delta_\eta(\alpha^\kappa \wedge B \wedge h)_N = \pm \vartheta(\alpha^\kappa \wedge B \wedge h) \wedge \delta_\eta \Omega. \quad (8.9)$$

We claim that

$$\delta_\eta(\alpha^\kappa \wedge B \wedge h)_N = (\alpha^\kappa \wedge h)_{N-1} \quad (8.10)$$

on $X \times X$. In fact, consider the weight $g = \alpha^\kappa \wedge g^\lambda$, cf. (5.2). From (2.2), keeping z on X and taking $\lambda = 0$ we get

$$\nabla_\eta(\alpha^\kappa \wedge B \wedge h) \wedge R = \alpha^\kappa \wedge h \wedge R$$

outside the diagonal. Identifying terms of degree $N - 1$ in $d\zeta$ we get

$$\delta_\eta(\alpha^\kappa \wedge B \wedge h)_N \wedge R = (\alpha^\kappa \wedge h)_{N-1} \wedge R.$$

By the generalized Poincaré–Lelong formula $R \wedge df = (2\pi i)^p[X]$ we can conclude that (8.10) holds on $X \times X$. Combining (8.9) and (8.10) we get that

$$\pm \vartheta(\alpha^\kappa \wedge B \wedge h) \wedge \delta_\eta \Omega = (\alpha^\kappa \wedge h)_{N-1} \quad (8.11)$$

on $X \times X$.

Let us now compute the kernel k where X is parametrized by $[\zeta_0, \zeta_1]$. Let δ_{ζ_j} denote interior multiplication by $\partial/\partial\zeta_j$. Notice that

$$\delta_\eta \delta_{\zeta_2} \cdots \delta_{\zeta_N} \Omega = 2\pi i (\zeta_1 z_0 - \zeta_0 z_1).$$

If we apply $\delta_{\zeta_2} \cdots \delta_{\zeta_N}$ to (8.11) we get

$$\pm \vartheta(\alpha^\kappa \wedge B \wedge h) = \frac{1}{2\pi i} \frac{\delta_{\zeta_2} \cdots \delta_{\zeta_N} (\alpha^\kappa \wedge h)_{N-1}}{\zeta_1 z_0 - \zeta_0 z_1}. \quad (8.12)$$

Notice that the denominator on the right-hand side is smooth. Recalling that $k(\zeta, z) = \pm \vartheta(\alpha^\kappa \wedge B \wedge h) \wedge \omega$ on $X \times X$, cf. (5.7), we get from (8.12) and (8.6), cf. (8.7),

Proposition 8.2 *With the notation above we have the explicit formula*

$$k(\zeta, z) = \frac{1}{2\pi i} \frac{\delta_{\zeta_2} \cdots \delta_{\zeta_N} (\alpha^\kappa \wedge h)_{N-1}}{\zeta_1 z_0 - \zeta_0 z_1} \delta_A \Omega \quad (8.13)$$

where X is parametrized by $[\zeta_0, \zeta_1]$.

8.3 Curves in \mathbb{P}^2

Assume now that $N = 2$ and $X = \{f = 0\}$, where f is a d -homogeneous form. Thus $s = \kappa + d - 2$. Let $\tilde{h}(w, z) = h_0 dw_0 + h_1 dw_1 + h_2 dw_2$ be a Hefer form for f ; i.e., h_ℓ are $d - 1$ -homogeneous and

$$2\pi i \sum_{\ell=0}^2 h_\ell(w, z)(w_\ell - z_\ell) = f(w) - f(z). \quad (8.14)$$

Notice that for degree reasons,

$$(\alpha^\kappa \wedge h)_1 = \alpha_0^\kappa (h_0(\alpha_0 \zeta, z) \gamma_0 + h_1(\alpha_0 \zeta, z) \gamma_1 + h_2(\alpha_0 \zeta, z) \gamma_2).$$

In what follows, for simplicity, let us write α rather than α_0 . If $\partial f / \partial \zeta_2$ is generically non-vanishing on X and

$$A = \frac{2\pi i}{\partial f / \partial \zeta_2} \frac{\partial}{\partial \zeta_2},$$

then $\delta_A df = 2\pi i$ (generically) on X . We get

$$k(\zeta, z) = \frac{1}{2\pi i} \frac{\zeta_0 d\zeta_1 - \zeta_1 d\zeta_0}{\zeta_1 z_0 - \zeta_0 z_1} \frac{2\pi i \delta_{\zeta_2}(\alpha^\kappa \wedge h)_1}{\partial f / \partial \zeta_2} \quad (8.15)$$

Notice furthermore that

$$\delta_{\zeta_2}(\alpha^\kappa \wedge h)_1 = \alpha^\kappa \left(h_2(\alpha \zeta, z) - \frac{\bar{\zeta}_2}{|\zeta|^2} \sum_{j=0}^2 h_j(\alpha \zeta, z) \zeta_j \right). \quad (8.16)$$

Proposition 8.3 *With the notation above we have*

$$k(\zeta, z) = \frac{1}{2\pi i} \frac{\zeta_0 d\zeta_1 - \zeta_1 d\zeta_0}{\zeta_1 z_0 - \zeta_0 z_1} \frac{2\pi i \alpha^\kappa}{\partial f / \partial \zeta_2} \left(h_2(\alpha \zeta, z) - \frac{\bar{\zeta}_2}{|\zeta|^2} \sum_{j=0}^2 h_j(\alpha \zeta, z) \zeta_j \right). \quad (8.17)$$

The second term in the brackets actually cancels out the singularity if X is smooth.

Proposition 8.4 *If X is smooth and $[\zeta_0, \zeta_1]$ are local homogeneous coordinates, then*

$$k(\zeta, z) = \frac{1}{2\pi i} \frac{\zeta_0 d\zeta_1 - \zeta_1 d\zeta_0}{\zeta_1 z_0 - \zeta_0 z_1} \frac{\alpha^\kappa 2\pi i h_2(\alpha \zeta, z)}{\partial f / \partial \zeta_2} + \dots, \quad (8.18)$$

where \dots is smooth and holomorphic in z .

Notice that $2\pi i h_2(\alpha \zeta, z) = \partial f / \partial \zeta_2$ on the diagonal.

Proof Differentiating (8.14) with respect to w_j gives

$$2\pi i h_j(w, z) = \frac{\partial f}{\partial w_j}(w) - 2\pi i \sum_{\ell} \frac{\partial h_{\ell}}{\partial w_j}(w_{\ell} - z_{\ell}).$$

Since f is d -homogeneous,

$$\sum_{j=0}^2 \zeta_j \frac{\partial f}{\partial \zeta_j} = d \cdot f(\zeta).$$

We conclude that

$$\sum_0^2 \zeta_j \tilde{h}_j(\alpha \zeta_j, z) = d\alpha^{d-1} \cdot f(\zeta) + \sum_0^2 b_{\ell}(\zeta, \alpha \zeta, z)(\alpha \zeta_{\ell} - z_{\ell}),$$

where $b_{\ell}(\zeta, w, z)$ are holomorphic, $d-2$ -homogeneous in (w, z) and 1-homogeneous in ζ . Since ζ is on X , $f(\zeta) = 0$. We thus get (8.18) where

$$\dots = \frac{1}{2\pi i} \frac{\zeta_0 d\zeta_1 - \zeta_1 d\zeta_0}{\zeta_1 z_0 - \zeta_0 z_1} B$$

and

$$B = \frac{\bar{\zeta}_2}{|\zeta|^2} \alpha^{\kappa} \sum_0^2 b_j(\zeta, \alpha \zeta, z)(\alpha \zeta_j - z_j).$$

Without loss of generality we may assume that $\zeta_0 = z_0 = 1$ and that ζ_1 is a local coordinate on X so that $\zeta' = g(\zeta_1)$. Since $\alpha = 1$ on the diagonal we then have that $B = B'(\zeta_1 - z_1)$ where B' is holomorphic in z . After homogenization we get that $B = B''(z_0 \zeta_1 - z_1 \zeta_0)$, where B'' is holomorphic in z . Thus the proposition follows. \square

Corollary 8.5 *If ζ_1 is a local coordinate, then taking $\zeta_1 = \tau$, $z_1 = t$, and $\zeta_0 = z_0 = 1$, we get*

$$k(\tau, t) = \frac{1}{2\pi i} \frac{d\tau}{\tau - t} \frac{\alpha^{\kappa} 2\pi i h_2(1, \tau, \zeta_2; 1, t, z_2)}{\partial f / \partial \zeta_2} + \dots, \quad (8.19)$$

where \dots is smooth and holomorphic in t .

Example 4 If $f(z) = z_0^3 + z_1^3 + z_2^3$, then X is a smooth surface of genus 1. We can take

$$\tilde{h}(z, w) = \sum_{j=0}^2 (z_j^2 + z_j w_j + w_j^2) dw_j / 2\pi i.$$

Then, where $[z_0, z_1]$ are homogeneous coordinates,

$$k(\zeta, z) = \alpha^\kappa \frac{1}{2\pi i} \frac{\zeta_0 d\zeta_1 - \zeta_1 d\zeta_0}{\zeta_1 z_0 - \zeta_0 z_1} \frac{1}{3\zeta_0^2} \left[z_2^2 + \alpha z_2 \zeta_2 + \alpha^2 \zeta_2^2 - \frac{\bar{\zeta}_2}{|\zeta|^2} \left(z_0^2 \zeta_0 + z_1^2 \zeta_1 + z_2^2 \zeta_2 + \alpha(z_0 \zeta_0^2 + z_1 \zeta_1^2 + z_2 \zeta_2^2) \right) \right].$$

Even if X is not smooth, by the same argument, one can identify the principal term of the kernel k .

Example 5 The curve $X = \{z_1^3 - z_2^2 z_0 = 0\}$ has a cusp singularity at $[1, 0, 0]$ and is smooth elsewhere. It is globally parametrized by

$$\mathbb{P}^1 \rightarrow X \subset \mathbb{P}^2, \quad [t_0, t_1] \mapsto [t_0^3, t_1^2 t_0, t_1^3].$$

We can choose the Hefer form

$$2\pi i \tilde{h}(w, z) = z_2^2 dw_0 - (z_1^2 + z_1 w_1 + w_1^2) dw_1 + (z_2 + w_2) w_0 dw_2.$$

By formula (8.17) we can now express k completely in terms of the parameters $[t_0, t_1]$ and $[\tau_0, \tau_1]$. However, we restrict to considering the principal term. Since we are primarily interested in the singularity, we consider the standard affinization where $z_0 = \zeta_0 = 1$. By the recipe above, then

$$k = \frac{1}{2\pi i} \frac{d\zeta_1}{\zeta_1 - z_1} \alpha^\kappa \frac{z_2 + \alpha \zeta_2}{2\zeta_2} + \dots \quad (8.20)$$

In this case \dots is not smooth but at least the singularity is smaller than in the leading term. Since $\alpha - 1 = \mathcal{O}(\zeta_1 - z_1) + \mathcal{O}(\zeta_2 - z_2)$, we can delete α as well in the leading term in (8.20). We then have

$$k = \frac{1}{2\pi i} \frac{(\zeta_2^2 - z_2^2) d\zeta_1}{(\zeta_1 - z_1)(\zeta_2 - z_2)2\zeta_2} + \dots$$

Letting $t = t_1/t_0$ and $\tau = \tau_1/\tau_0$ we are then left with

$$k = \frac{1}{2\pi i} \frac{(\tau^6 - t^6) d\tau}{(\tau^2 - t^2)(\tau^3 - t^3)\tau^2} + \dots$$

where the leading term is precisely the kernel in the last example in [6, Sect. 8].

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