



Electromagnetically Induced Acoustic Transparency with a Superconducting Circuit

Downloaded from: <https://research.chalmers.se>, 2020-09-19 10:30 UTC

Citation for the original published paper (version of record):

Andersson, G., Ekström, M., Delsing, P. (2020)

Electromagnetically Induced Acoustic Transparency with a Superconducting Circuit

Physical Review Letters, 124(24)

<http://dx.doi.org/10.1103/PhysRevLett.124.240402>

N.B. When citing this work, cite the original published paper.

Electromagnetically Induced Acoustic Transparency with a Superconducting Circuit

Gustav Andersson[✉],* Maria K. Ekström, and Per Delsing

*Department of Microtechnology and Nanoscience MC2, Chalmers University of Technology,
Kemivägen 9 SE-41296 Göteborg, Sweden*



(Received 3 December 2019; accepted 28 May 2020; published 16 June 2020)

We report the observation of electromagnetically induced transparency (EIT) of a mechanical field, where a superconducting artificial atom is coupled to a 1D-transmission line for surface acoustic waves. An electromagnetic microwave drive is used as the control field, rendering the superconducting transmon qubit transparent to the acoustic probe beam. The strong frequency dependence of the acoustic coupling enables EIT in a ladder configuration due to the suppressed relaxation of the upper level. Our results show that superconducting circuits can be engineered to interact with acoustic fields in parameter regimes not readily accessible to purely electromagnetic systems.

DOI: [10.1103/PhysRevLett.124.240402](https://doi.org/10.1103/PhysRevLett.124.240402)

Electromagnetically induced transparency (EIT) is a quantum interference effect where an electromagnetic field controls the response of a three-level medium to a probe field [1]. In contrast to the related phenomenon of Autler-Townes splitting, EIT arises due to the interference of excitation pathways in the coherent interaction of the atom(s) with the radiation field, and its signature is the appearance of a narrow transparent window inside the atomic absorption spectrum. EIT has been observed in atomic three-level systems with either Λ or ladder-type configuration [2,3], and an analogue form of induced transparency has been demonstrated in optomechanical devices where light beams interact with a mechanical resonator through radiation pressure coupling [4,5]. Because of the difficulty in engineering the requisite metastable states, EIT in a circuit quantum electrodynamics architecture [6,7] was demonstrated only recently [8,9], using the combined states of an artificial atom in the form of a superconducting qubit and a three-dimensional microwave cavity. For superconducting circuits coupled to open transmission lines, what was thought to be observations of EIT [10,11] have been shown to in fact be the Autler-Townes splitting [12].

Since its discovery, a range of potential applications for EIT in nonlinear optics and quantum information have been suggested, including quantum memories, routing and cross-phase modulation [13–15]. The modulation of absorption and emission cross sections due to EIT could also have applications in the context of heat engines [16].

Strong coupling of superconducting qubits to surface acoustic waves (SAW) has been demonstrated in both waveguide and cavity settings [17,18], and used to generate nonclassical phonon states [19] as well as for quantum state transfer [20]. Here we exploit the strong frequency dependence of the piezoelectric coupling between a superconducting transmon circuit and SAW to control the relative decoherence rates of the transmon states, effectively enhancing the relative lifetime of the second excited state sufficiently to allow genuine EIT to occur. The delay line setup is commonly used to probe the properties of physical systems with SAW [21,22] and provide an acoustic drive field, including demonstrations of coherent interference effects in optomechanics [23]. Using a SAW delay line, we show electromagnetically induced acoustic transparency in both reflection and transmission of the probe. To our knowledge this constitutes the first observation of EIT for a propagating mechanical mode. For strong dressing fields the system crosses over to the Autler-Townes regime, where routing of SAW phonons has been shown by fast switching of the control tone [24].

The ladder-type three-level system is formed by the ground and two first excited states of a superconducting transmon circuit [25]. The transmon, which is fabricated on a GaAs substrate, couples piezoelectrically to the propagating SAW field via an interdigitated transducer (IDT). The IDT spans $N_p = 25$ periods and has a double-finger structure to suppress internal mechanical reflections [26]. A SQUID loop allows for tuning of the transition frequencies. Whereas a quantum emitter in an electromagnetic transmission line couples to all modes of the propagating field, the periodic structure of the IDT restricts the coupling of the SAW-qubit interaction to a bandwidth of approximately $\Delta_{\text{IDT}}/2\pi = 0.9f_{\text{IDT}}/N_p$ [27,28]. With an IDT center frequency f_{IDT} of 2.26 GHz, we obtain $\Delta_{\text{IDT}}/2\pi \approx 81$ MHz. The transition frequency ω_{10} is tuned into this band using

Published by the American Physical Society under the terms of the Creative Commons Attribution 4.0 International license. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Funded by Bibsam.

an external magnetic flux. The IDT finger structure provides a capacitance to the transmon circuit giving rise to a charging energy of $E_C/h = 129$ MHz. As the transmon anharmonicity, given by E_C , is sufficient to ensure $\Delta_{\text{IDT}} < \omega_{10} - \omega_{12}$, maximizing the SAW coupling of the first transition by setting $\omega_{10}/2\pi = f_{\text{IDT}}$ implies $\Gamma_{10} \gg \Gamma_{21}$ where Γ_{ij} denotes the spontaneous emission rate from the state i to state j . The filtering provided by the IDT thus strongly suppresses the coupling of higher transmon levels to SAW.

A weak SAW probe beam is launched towards the transmon using a double-finger IDT with 150 periods, located a distance $300 \mu\text{m}$ away. The acoustic signal reflected back to the launcher IDT is measured using a microwave circulator while a control field is applied via a capacitively coupled electrical gate, as shown schematically in Fig. 1.

The reflection coefficient of the probe beam is given by [30]

$$r = i \frac{\Gamma_{10}}{\Omega_p} \langle \sigma_- \rangle, \quad (1)$$

where Ω_p is the Rabi frequency given by the probe amplitude and σ_- denotes the lowering operator between the $|0\rangle$ and $|1\rangle$ transmon states. For a weak probe under the application of a control field this yields in the steady state [10]

$$r = - \frac{\Gamma_{10}}{2(\gamma_{10} - i\Delta_p) + \frac{\Omega_c^2}{2(\gamma_{20} - i\Delta_p - i\Delta_c)}}, \quad (2)$$

where $\Delta_p = \omega_p - \omega_{10}$ is the probe detuning. The control field has an effective Rabi frequency Ω_c and a detuning $\Delta_c = \omega_c - \omega_{21}$. The decoherence rates γ_{ij} correspond to decay rates of the off diagonal density matrix elements and determine the conditions for realizing EIT. A common procedure in optical EIT measurements is to sweep the probe detuning Δ_p while applying a resonant control field ($\Delta_c = 0$) [1]. The limited bandwidth of the probe IDT makes this approach impractical in our case. Instead, we adopt a reversed scheme where the frequency of the acoustic probe is fixed on resonance while sweeping the frequency and amplitude of the control field. As the control field interaction is purely electromagnetic it can be applied

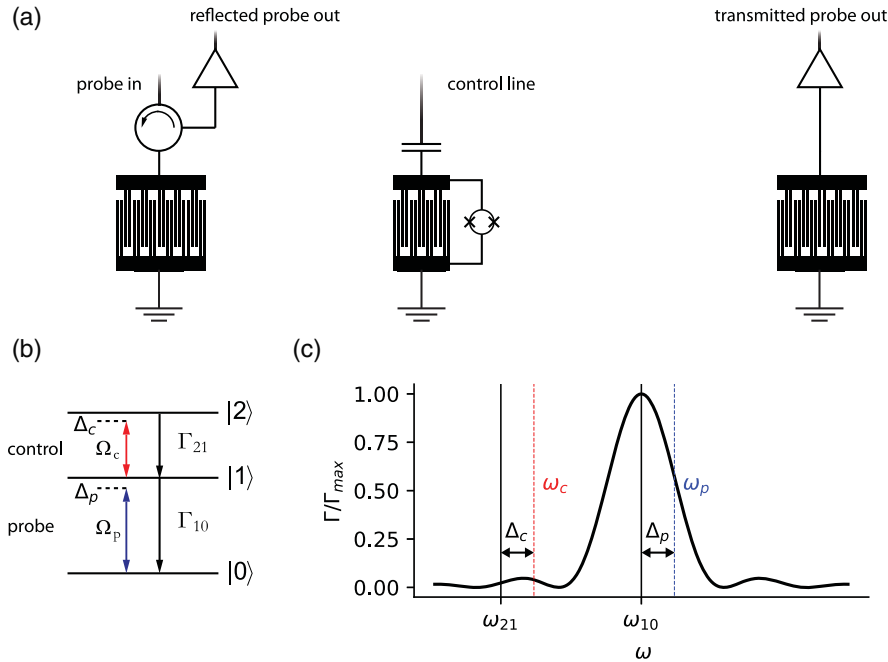


FIG. 1. (a) Schematic of the device layout. An IDT with 150 finger pairs is used to launch the SAW probe beam towards the transmon and detect the reflected signal. The transmon circuit interacts with the acoustic probe via its 25 IDT finger pairs, and the control field is applied to a gate electrode capacitively coupled to the transmon. A further IDT, also with 150 finger pairs, is used to receive and transduce the transmitted SAW signal. The structure is fabricated from aluminum on a piezoelectric GaAs substrate. (b) Level scheme of the ladder system of the artificial atom. The acoustic probe field has an amplitude Ω_p and detuning Δ_p from the $|0\rangle \leftrightarrow |1\rangle$ transition. A purely electromagnetic control field is applied close to the $|1\rangle \leftrightarrow |2\rangle$ transition with detuning Δ_c . The direct $|0\rangle \leftrightarrow |2\rangle$ transition is suppressed by parity conservation and spontaneous emission occurs to the next-lower level. (c) Schematic illustrating the spontaneous emission rate of the artificial atom into the SAW transmission line, which depends on frequency as the square of a sinc function [29]. We operate the device in a regime where the $|0\rangle \leftrightarrow |1\rangle$ transition at ω_{01} is maximally coupled, and the $|1\rangle \leftrightarrow |2\rangle$ transition is completely outside the main lobe of the coupling function.

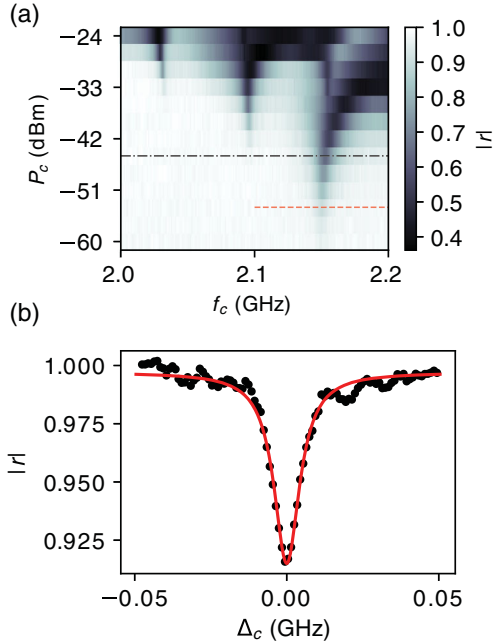


FIG. 2. (a) Normalized reflection coefficient measured as a function of (room temperature) control power P_c and frequency of the control field. The probe frequency is 2.2684 GHz. The black dash-dotted line indicates the crossover from the low power EIT to the Autler-Townes effect. With increasing power the level scheme of Fig. 1 breaks down as multiphoton transitions are strongly driven. The red dashed line indicates the line cut shown with a Lorentzian fit (red line) in (b). The dip in reflection is due to the EIT.

across a wide frequency range. Figure 2(a) shows the reflected probe amplitude as a function of applied control power and frequency for a probe beam at 2.2684 GHz. As the control frequency is swept into resonance with the $|1\rangle \leftrightarrow |2\rangle$ transition at $f_{12} = 2.15$ GHz, a dip appears in the reflected amplitude. Together with an independent estimate of γ_{10} , this measurement allows for extracting all parameters relevant to discerning EIT from the Autler-Townes splitting [12]. The rate γ_{10} is obtained from analyzing the line shape obtained while sweeping the transmon frequency ω_{10} around the probe frequency in the absence of a control field ($\Omega_c = 0$), and found to be $\gamma_{10}/2\pi = 21$ MHz.

A resonant probe beam implies setting $\Delta_p = 0$ in Eq. (2), which yields a negative Lorentzian with a half width at half maximum given by

$$\gamma_{\text{EIT}} = \gamma_{20} + \frac{\Omega_c^2}{4\gamma_{10}}. \quad (3)$$

This linewidth is linear in the control power and limited in sharpness by the decoherence γ_{20} of the upper level in the ladder. In Fig. 3 we plot the linewidth of this Lorentzian as a function of the control tone power applied at room temperature as well as linear fit. From the fit we extract a

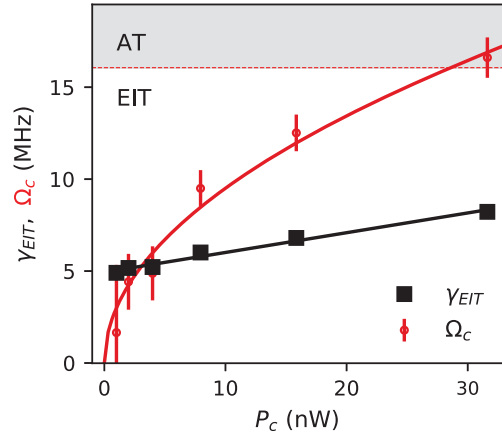


FIG. 3. EIT linewidth γ_{EIT} as a function of the power P_c applied to the control gate at room temperature. The black line shows a linear fit. The extracted fit parameters are used to obtain $\gamma_{20} = 4.94$ MHz and the effective control strength Ω_c as a function of power (solid red line) from Eq. (3). The red dots show the extracted Ω_c for each data point with error bars. The threshold $\Omega_c < \gamma_{10} - \gamma_{20}$ is indicated by the dashed red line.

$\gamma_{20}/2\pi$ of 4.94 ± 0.14 MHz. As the coupling of the $|1\rangle \leftrightarrow |2\rangle$ transition to SAW is suppressed by more than an order of magnitude relative to $|1\rangle$, this decoherence rate is not dominated by SAW emission, but rather other sources of pure dephasing. We stress that the condition $\gamma_{10} > \gamma_{20}$ arises due to the frequency dependence of the acoustic coupling, and is necessary to enable EIT in ladder type systems. While the constant term of the linear fit yields γ_{20} , we use the known γ_{10} and the slope to extract the control field drive strength Ω_c as a function of power applied at room temperature.

The presence of a transparent window in the scattering off a three-level system does not necessarily imply EIT and substantial theoretical analysis has been developed to determine whether EIT or Autler-Townes splitting occurs under given experimental conditions [12]. EIT relies on the destructive interference of two excitation paths. In the ladder case the direct $|0\rangle \rightarrow |1\rangle$ transition interferes with the path exciting to the upper level and back, $|0\rangle \rightarrow |1\rangle \rightarrow |2\rangle \rightarrow |1\rangle$. Spontaneous emission or dephasing of the state $|2\rangle$ suppresses this interference. Using the criterion from Ref. [12], the quantitative distinction arises from analyzing the poles of Eq. (2). If under resonant control ($\Delta_c = 0$), the poles of Eq. (2) are purely imaginary, the reflection coefficient as a function of probe frequency can be expressed as the difference of two Lorentzians centered at the same frequency. This is the condition for EIT and in our system equivalent to $\Omega_c < \gamma_{10} - \gamma_{20}$, which implies EIT can only be observed if $\gamma_{10} > \gamma_{20}$. The threshold for the drive amplitude that separates the EIT and AT regimes is shown as the red dashed line in Fig. 3. If the drive strength is increased beyond this limit the reflection coefficient is given by the sum of two Lorentzians separated by the drive

strength Ω_c . This is the Autler-Townes splitting. As shown in Fig. 3, the lower control powers are insufficient for Autler-Townes splitting to appear and the transparency features observed are due to the EIT. The crossover to the Autler-Townes regime appears at $\Omega_{c,t} = \gamma_{10} - \gamma_{20} = 2\pi \times 16.1$ MHz, corresponding to a control power of -45 dBm. Figure 2 shows the dip in reflection arising from EIT at a drive amplitude of $\Omega_c/2\pi = 6.1$ MHz, well below $\Omega_{c,t}$.

We further measure the acoustic EIT in transmission, where the transmitted SAW signal is transduced by an IDT at a distance $100 \mu\text{m}$ from the transmon. In this measurement we use a different scheme where the probe frequency is fixed, and an external magnetic flux is used to tune the qubit frequency $\omega_{10}/2\pi$, thereby sweeping Δ_p . As the control frequency also remains fixed, this scheme will simultaneously sweep both Δ_c and Δ_p . To first order, the anharmonicity is not affected by the tuning, yielding $\Delta_c = \Delta_p + \delta$, where δ is the residual control detuning at probe resonance $\Delta_p = 0$. The case $\delta = 0$ corresponds to perfect alignment of control and probe frequencies with the anharmonicity. With the definition $t = 1 + r$ and Δ_p as the independent variable we get the transmission coefficient

$$t = 1 - \frac{\Gamma_{10}}{2(\gamma_{10} - i\Delta_p) + \frac{\Omega_c^2}{2(\gamma_{20} - 2i\Delta_p - i\delta)}}. \quad (4)$$

Figure 4 shows the normalized transmission amplitude for different drive strengths. We extract an upper level decoherence rate of $\gamma_{20}/2\pi = 4.5 \pm 0.6$ MHz. While this

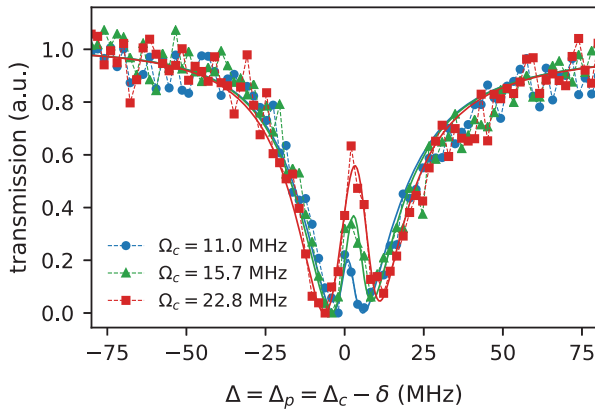


FIG. 4. Normalized transmission coefficient in the EIT (blue), crossover (green), and Autler-Townes (red) regimes. Solid lines are fits to data and yield $\gamma_{20}/2\pi = 4.5$ MHz. The probe tone is fixed at 2.2644 GHz, and the probe and pump detunings are swept by tuning the transmon frequency via an external magnetic field. Because of the fixed anharmonicity of the transmon, the ω_{10} and ω_{21} transition frequencies are swept uniformly. The slight asymmetry in the line shape is due to a residual control detuning $\delta \approx 4$ MHz at $\Delta_p = 0$ as well as the interference of electromagnetic cross talk between the launcher and receiver IDTs.

estimate is less precise than the result obtained in the reflection measurement, they are consistent insofar as the error margin in reflection (4.94 ± 0.14 MHz) falls completely within this range.

In conclusion, we have demonstrated that EIT can be generated in a mechanical mode of propagating surface acoustic waves by using an electromagnetic control signal. We show consistent data in reflection and transmission. By demonstrating quantum interference of acoustically and electromagnetically driven excitations, this experiment suggests applications in phononic quantum information architectures, where SAW phonons couple disparate quantum systems. We show that the piezoelectric coupling of superconducting qubits to short-wavelength SAW has a frequency dependence offering the possibility to engineer relaxation rates, allowing EIT to be observed in waveguide QED with superconducting circuits. This principle could be further exploited in quantum acoustic experiments [31], for example, to generate population inversion and single-atom sound lasing. EIT is further associated with slow light [32], leading to a reduction in the group velocity as the field propagates through the EIT medium. The amount of slowdown is limited by the linewidth of the EIT window. While in our case this limit corresponds approximately to a factor of 3 ($v_g \approx 1000$ m/s), improved coherence in the second excited state would enable larger reductions of the sound velocity. Larger group delays could also be achieved by using an array of artificial atoms rather than only a single transmon.

We acknowledge fruitful discussions with B. Suri and G. Johansson. This work was supported by the Knut and Alice Wallenberg foundation and by the Swedish Research Council, VR. This project has also received funding from the European Union's Horizon 2020 research and innovation programme under Grant Agreement No. 642688 (SAWtrain).

*gustav.andersson@chalmers.se

- [1] M. Fleischhauer, A. Imamoglu, and J.P. Marangos, *Rev. Mod. Phys.* **77**, 633 (2005).
- [2] J. Gea-Banacloche, Y. Q. Li, S. Z. Jin, and M. Xiao, *Phys. Rev. A* **51**, 576 (1995).
- [3] B. Gouraud, D. Maxein, A. Nicolas, O. Morin, and J. Laurat, *Phys. Rev. Lett.* **114**, 180503 (2015).
- [4] S. Weis, R. Riviere, S. Deleglise, E. Gavartin, O. Arcizet, A. Schliesser, and T. J. Kippenberg, *Science* **330**, 1520 (2010).
- [5] A. H. Safavi-Naeini, T. P. M. Alegre, J. Chan, M. Eichenfield, M. Winger, Q. Lin, J. T. Hill, D. E. Chang, and O. Painter, *Nature (London)* **472**, 69 (2011).
- [6] A. Wallraff, D. I. Schuster, A. Blais, L. Frunzio, R. S. Huang, J. Majer, S. Kumar, S. M. Girvin, and R. J. Schoelkopf, *Nature (London)* **431**, 162 (2004).
- [7] A. Blais, R.-S. Huang, A. Wallraff, S. M. Girvin, and R. J. Schoelkopf, *Phys. Rev. A* **69**, 062320 (2004).

- [8] S. Novikov, T. Sweeney, J. E. Robinson, S. P. Premaratne, B. Suri, F. C. Wellstood, and B. S. Palmer, *Nat. Phys.* **12**, 75 (2016).
- [9] J. Long, H. S. Ku, X. Wu, X. Gu, R. E. Lake, M. Bal, Y.-x. Liu, and D. P. Pappas, *Phys. Rev. Lett.* **120**, 083602 (2018).
- [10] A. A. Abdumalikov, O. Astafiev, A. M. Zagoskin, Y. A. Pashkin, Y. Nakamura, and J. S. Tsai, *Phys. Rev. Lett.* **104**, 193601 (2010).
- [11] I.-C. Hoi, C. M. Wilson, G. Johansson, T. Palomaki, B. Peropadre, and P. Delsing, *Phys. Rev. Lett.* **107**, 073601 (2011).
- [12] P. M. Anisimov, J. P. Dowling, and B. C. Sanders, *Phys. Rev. Lett.* **107**, 163604 (2011).
- [13] L. Ma, O. Slattery, and X. Tang, *J. Opt.* **19**, 043001 (2017).
- [14] K. Xia, F. Jelezko, and J. Twamley, *Phys. Rev. A* **97**, 052315 (2018).
- [15] H. Schmidt and A. Imamoglu, *Opt. Lett.* **21**, 1936 (1996).
- [16] S. E. Harris, *Phys. Rev. A* **94**, 053859 (2016).
- [17] M. V. Gustafsson, T. Aref, A. F. Kockum, M. K. Ekstrom, G. Johansson, and P. Delsing, *Science* **346**, 207 (2014).
- [18] B. A. Moores, L. R. Sletten, J. J. Viennot, and K. W. Lehnert, *Phys. Rev. Lett.* **120**, 227701 (2018).
- [19] K. J. Satzinger, Y. P. Zhong, H. S. Chang, G. A. Peairs, A. Bienfait, M.-H. Chou, A. Y. Cleland, C. R. Conner, E. Dumur, J. Grebel, I. Gutierrez, B. H. November, R. G. Povey, S. J. Whiteley, D. D. Awschalom, D. I. Schuster, and A. N. Cleland, *Nature (London)* **563**, 661 (2018).
- [20] A. Bienfait, K. J. Satzinger, Y. P. Zhong, H.-S. Chang, M.-H. Chou, C. R. Conner, É. Dumur, J. Grebel, G. A. Peairs, R. G. Povey, and A. N. Cleland, *Science* **364**, 368 (2019).
- [21] A. Wixforth, J. P. Kotthaus, and G. Weimann, *Phys. Rev. Lett.* **56**, 2104 (1986).
- [22] M. Weiler, H. Huebl, F. S. Goerg, F. D. Czeschka, R. Gross, and S. T. B. Goennenwein, *Phys. Rev. Lett.* **108**, 176601 (2012).
- [23] K. C. Balram, M. Davanco, J. D. Song, and K. Srinivasan, *Nat. Photonics* **10**, 346 (2016).
- [24] M. K. Ekström, T. Aref, A. Ask, G. Andersson, B. Suri, H. Sanada, G. Johansson, and P. Delsing, *New J. Phys.* **21**, 123013 (2019).
- [25] J. Koch, T. M. Yu, J. M. Gambetta, A. A. Houck, D. I. Schuster, J. Majer, A. Blais, M. H. Devoret, S. M. Girvin, and R. J. Schoelkopf, *Phys. Rev. A* **76**, 042319 (2007).
- [26] T. W. Bristol, W. R. Jones, P. B. Snow, and W. R. Smith, in *Proc. IEEE Ultrasonics Symposium* (1972), pp. 343–345.
- [27] S. Datta, *Surface Acoustic Wave Devices* (Prentice Hall, Englewood Cliffs, 1986).
- [28] T. Aref, P. Delsing, M. K. Ekström, A. F. Kockum, M. V. Gustafsson, G. Johansson, P. J. Leek, E. Magnusson, and R. Manenti, in *Superconducting Devices in Quantum Optics*, edited by R. H. Hadfield and G. Johansson (Springer International Publishing, Cham, 2016), pp. 217–244.
- [29] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.124.240402> for acoustic coupling and EIT linewidth estimation.
- [30] O. Astafiev, A. M. Zagoskin, A. A. Abdumalikov, Y. A. Pashkin, T. Yamamoto, K. Inomata, Y. Nakamura, and J. S. Tsai, *Science* **327**, 840 (2010).
- [31] A. Frisk Kockum, P. Delsing, and G. Johansson, *Phys. Rev. A* **90**, 013837 (2014).
- [32] L. V. Hau, S. E. Harris, Z. Dutton, and C. H. Behroozi, *Nature (London)* **397**, 594 (1999).