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Unraveling the impact of different thermal quenching routes on the luminescence efficiency of the $Y_3Al_5O_{12}:Ce^{3+}$ phosphor for white light emitting diodes†

Yuan-Chih Lin, Marco Bettinelli, Suchinder K. Sharma, Britta Redlich, Adolfo Spighi and Maths Karlsson

Cerium doped yttrium aluminium garnet, $Y_{3-z}CeAl_5O_{12}$, is the prototype material for solid-state white lighting and it still is an important white LED phosphor. However, fundamental understanding of the thermal quenching of luminescence, which leads to a pronounced reduction of the emission intensity under high-power light-emitting diode operation, remains to be obtained. Here we show, through a multitechnique approach based on photoluminescence, thermoluminescence and mode-selective vibrational excitation experiments that thermal quenching of luminescence in $Y_{3-z}CeAl_5O_{12}$ is caused by a combined effect of thermal ionization, thermally activated concentration quenching, and thermally activated 5d→4f crossover relaxation via electron–phonon coupling, and establish the general trends upon variation of the Ce$^{3+}$ concentration and temperature. Thermal quenching below 600 K is primarily the result of concentration quenching and crossover relaxation, which can be suppressed by keeping the Ce$^{3+}$ dopant concentration far below 0.7 mol%, whereas for temperatures above 600 K thermal ionization is the dominating quenching process. This new insight into the interplay between different thermal quenching processes provides design principles for optimizing the light emittance and colour stability of new phosphor materials used in white lighting devices characterized by certain operating temperatures.

1 Introduction

Inorganic phosphors are of considerable interest for application in phosphor converted white light emitting diodes (pc-WLEDs) and over the last decades a large number of compounds have been synthesized. Among the most widely used phosphors is yttrium aluminium garnet ($Y_3Al_5O_{12}$, YAG), which when substituted with a few percent of the activator ion Ce$^{3+}$ to replace electronic 4f–5d transitions is thought to result from one or several migration among Ce$^{3+}$ ions to luminescence killer centers.6 Generally, thermal quenching in Ce$^{3+}$-doped phosphors exhibiting electronic 4f–5d transitions is thought to result from one or several of the three processes: (1) thermal ionization of the Ce$^{3+}$ 5d electron into the conduction band (CB) of the YAG host crystal followed by charge trapping at defects acting as luminescence killer centres [Fig. 1(a)];7,9 (2) thermally activated nonradiative energy migration among Ce$^{3+}$ ions to luminescence killer centers (known as concentration quenching) [Fig. 1(b)];10–12 and (3) thermally activated 5d→4f crossover relaxation via electron–phonon coupling [Fig. 1(c)].9,13–15
The first type of process (thermal ionization) is predominantly dependent on the energy difference between the emitting level of the Ce$^{3+}$ ion and the CB minimum, which is here called ionization energy. Recent bandgap engineering studies on YAG:Ce$^{3+}$ and cation-substituted variants, combined with studies of vacuum referred binding energy (VRBE) diagrams and thermoluminescence (TL) excitation experiments, show that the quenching is stronger when the ionization energy is smaller.\textsuperscript{17–21} The second type of process (nonradiative energy migration among Ce$^{3+}$ ions to killer centers (the overlap between the excitation (Exc.) and emission (Emi.) spectra is indicated to the left, and the extent of overlap determines the probability of resonance transfer of excitation energy between Ce$^{3+}$ ions), and (c) thermally activated crossover from the 5d$^1$ excited state to the 4f ground state via electron–phonon coupling. 5d$^1$ denotes the lowest-energy-lying level of the 5d configuration of Ce$^{3+}$. The figure is modified from ref. 16 (Copyright 2018 American Chemical Society).

The type of process (nonradiative energy migration among Ce$^{3+}$ ions and thus to an increased probability of trapping the migrating excitation energy by defects.\textsuperscript{11} Bachmann et al.\textsuperscript{6} attributed, in a systematic investigation of YAG:Ce$^{3+}$ with 0.033, 0.33, 1.0, and 3.33 mol% Ce$^{3+}$ doping, the lower quenching temperature for 1.0, and 3.33 mol% Ce$^{3+}$ to an effect of thermally activated concentration quenching. However, the possibility of Ce$^{3+}$ concentration dependent cross-over relaxation (as the incorporation of Ce$^{3+}$ may soften the crystal structure and hence increase the phonon population at a given temperature)\textsuperscript{16,21} and/or thermal ionization (as the incorporation of Ce$^{3+}$ may cause changes of the ionization energy) should not be neglected in the elucidation of a full understanding of the thermal quenching behaviour of YAG:Ce$^{3+}$ and other materials alike. In particular, because of the many experimental challenges of probing explicitly energy migration and electron–phonon coupling processes in materials, it is only through the combinatorial use of experimental and theoretical techniques, including novel machine learning methods to screen thermally robust phosphors,\textsuperscript{24} that a mechanistic understanding of thermal quenching is likely to emerge.

Here we present a systematic analysis of the thermal quenching of luminescence in YAG:x\%Ce$^{3+}$, as a function of Ce$^{3+}$ concentration (x = 0.2, 1, 2, and 3) and temperature (T = 80–860 K), based on a multitechnique approach. The techniques used are photoluminescence (PL), TL, and mode-selective vibrational excitation (together with thermal simulations) techniques. Our results show that the thermal quenching is caused by the combined effect of thermal ionization, thermally activated concentration quenching, and thermally activated 5d $\rightarrow$ 4f cross-over relaxation via electron–phonon coupling, with different weights depending on the temperature. This new insight into the interplay between different thermal quenching processes provides design principles for optimizing the light emittance and colour stability of new phosphor materials used in white lighting devices characterized by certain operating temperatures, e.g. > 150 °C under high-power operating conditions.\textsuperscript{25}

### 2 Results and discussion

#### 2.1 Photoluminescence decay time and quantum yield

Fig. 2(a) shows the temperature dependence of the luminescence decay time, $\tau$, for YAG:Ce$^{3+}$ with a Ce$^{3+}$ concentration corresponding to 0.2, 1, 2, and 3 mol%. All decay curves can be adequately described by a single-barrier quenching model (solid lines) according to

$$\tau(T) = \frac{1}{F_r + F_n \exp(-\Delta E/k_B T)}, \quad (1)$$

where $F_r$ is the radiative rate, $F_n$ is the attempt rate of nonradiative processes, $k_B$ is the Boltzmann constant, and $\Delta E$ is the activation energy for the overall quenching behaviour. $\Delta E$ is found to decrease from 1.11 eV for the lowest Ce$^{3+}$ dopant concentration to 0.32 eV for the highest one [Table 1 and Fig. 2(b)]. The thermal quenching temperature, $T_{50\%}$, which is defined here as the temperature at which the luminescence efficiency (estimated from decay time or quantum yield data)
becomes 50% of the low-temperature value, decreases from 778 K for 0.033 mol% to 556 K for 3 mol% [Table 1 and Fig. 2(c)].

Fig. 3 shows the temperature dependence of the internal PL quantum yield (PLQY) for YAG:Ce 3+ with a Ce 3+ concentration of 1, 2, and 3 mol%, upon excitation at 454 nm and 340 nm, respectively. The PLQY is only weakly affected by temperature increase up to about 300–500 K. For higher temperatures, the PLQY decreases drastically with increasing temperature, in a way quite similar to the luminescence decay time [Fig. 2(a)]. For excitation at 454 nm, the $T_{50\%}$ values are generally lower by 10–20% compared to the $T_{50\%}$ values as extracted from the temperature dependence of the PLQY data, as measured for excitation at 454 nm and 340 nm (see Fig. 3), respectively.

## 2.2 Optical spectra and VRBE diagram analysis

Fig. 4(a) shows the diffuse reflectance/transmittance spectra of YAG:xCe 3+ with $x = 0.2, 1, 2,$ and 3 mol%, as measured at room temperature (RT). The two dips in the spectra at around 340 and 460 nm correspond to the $4f \rightarrow 5d_2$ and $4f \rightarrow 5d_1$ (absorption) transitions of Ce 3+, which are denoted by $\lambda_{5d_2}$ and $\lambda_{5d_1}$, respectively (Table 1). The magnitude of these dips increases generally with increasing Ce 3+ concentration, as expected from the Beer–Lambert law. The corresponding emission spectra [Fig. 4(b and c)] were measured for excitation at 340 and 454 nm, respectively, and show one asymmetric broad band between approximately 500 and 700 nm. The emission spectra are almost identical for excitation at 340 nm and 454 nm, suggesting that the electronic-vibrational crossing from the 5d 2 (higher) to 5d 1 (lower) orbital excited states of Ce 3+ is a very fast process, followed by a change of (local) structural dynamics arising from vibrational relaxation. Therefore, this change of structural dynamics around the luminescent Ce 3+ ions has no significant effect on the subsequent 5d 1 $\rightarrow$ 4f emission process of Ce 3+.

Using the respective energies of the $4f \rightarrow 5d_2$ and $4f \rightarrow 5d_1$ absorption transitions of Ce 3+ (associated with $\lambda_{5d_2}$ and $\lambda_{5d_1}$, respectively), and the energy positions of the bottom of the CB ($E_l$) and the 4f level of Ce 3+ ($E_{5d_1}$), we constructed the VRBE diagram of Ce 3+ in YAG, see Table 1. This analysis shows that the thermal ionization energy for the 5d 1 level ($\Delta E_{i,5d_1}$) is about 5–8 times larger than that for the 5d 2 level ($\Delta E_{i,5d_2}$), see Fig. 1(a) and Table 1. Both $\Delta E_{i,5d_1}$ and $\Delta E_{i,5d_2}$ increase slightly with increasing Ce 3+ concentration, as an effect of a lowering of the 5d 2 and 5d 2 levels (red-shifting effect). This shows that the thermal ionization process is very weakly dependent on the Ce 3+ concentration for $<3$ mol% Ce 3+.

$\Delta E_{i,5d_2}$ takes on values in the range of 1.11–1.13 eV for variable Ce 3+ concentration. For the lowest Ce 3+ concentrations (0.033 and 0.05 mol%), i.e., when the Ce 3+ concentration is so low so that the Ce 3+ ions in all likelihood can be considered as isolated luminescent centers, $\Delta E_{i,5d_2}$ is quite similar to $\Delta E$ (1.06–1.11 eV, see Table 1). In effect, this suggests that for these Ce 3+ concentrations, thermal quenching is mainly, if not fully, governed by thermal ionization. For higher Ce 3+ concentrations (0.2–3 mol%), $\Delta E_{i,5d_2}$ deviates significantly from $\Delta E$. This deviation increases systematically with increasing Ce 3+ concentration and points toward that another or other quenching processes, which have activation energies lower than $\Delta E$, increasingly dominate the overall thermal quenching of luminescence.

Further support for the contribution of more than one quenching process to the observed thermal quenching behaviour
Table 1  Compilation of thermal quenching data for YAG:Ce$^{3+}$. VRBE parameters ($E_C$, bottom of conduction band, $E_{4f}$, 4f level of Ce$^{3+}$, and $\Delta E_{5d4f}/\Delta E_{i,5d_i}$: energy difference between $E_C$ and the 5d$_i$/5d$_j$ of Ce$^{3+}$), PL spectral data ($\lambda_{4f}$, maximum of the 5d$_i$/5d$_j$/emission band of Ce$^{3+}$), thermal quenching measures ($\Delta E$, $\Delta E_{TL}$ and $\Delta E_{ph}$: activation energies obtained from decay time, TL, and FELIX experiments, respectively), $T_{50\%}$: thermal quenching temperature, and $\eta_{loss}$: loss of luminescence efficiency with respect to the one at 80 K, and PLQY, for YAG:Ce$^{3+}$ with varying Ce$^{3+}$ concentration and average distance between Ce$^{3+}$ ions ($R_{Ce}$).

<table>
<thead>
<tr>
<th>YAG:Ce$^{3+}$ (R$_{Ce}$) (Å)</th>
<th>Ce$^{3+}$ concentration (mol%)</th>
<th>VRBE parameters &amp; spectral data (exc. 340/454 nm)</th>
<th>PLQY (exc. 454 nm)</th>
<th>T$_{50%}$ (K)</th>
<th>$\eta_{loss}$ (%)</th>
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<tbody>
<tr>
<td>0.033$^a$ (36.5)</td>
<td>0.05$^b$ (31.7)</td>
<td>0.2 (20.0)</td>
<td>0.5$^c$ (14.8)</td>
<td>1 (11.8)</td>
<td>2 (9.5)</td>
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<tr>
<td>$E_C$ (eV)</td>
<td>$E_{4f}$ (eV)</td>
<td>$\lambda_{4f}$ (nm)</td>
<td>$\Delta E_{5d4f}$ (eV)</td>
<td>$\Delta E_{i,5d_i}$ (eV)</td>
<td>$\lambda_{emi}$ (nm)</td>
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<td>1.11</td>
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<td>0.55</td>
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<td>PLQY (exc. 454 nm)</td>
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$^a$ Reproduced from ref. 6. $^b$ Reproduced from ref. 26. $^c$ Reproduced from ref. 20. $^d$ Cited from ref. 19 and 27. $^e$ Derived by assuming that the $E_C$ and $E_{4f}$ values are independent of Ce$^{3+}$ concentration (for the studied range). $^f$ Estimated at 600 K. $^g$ Measured at 300 K. $^h$ Derived from the TL glow curves corresponding to trap3.

Fig. 3  Variable temperature PLQY of YAG:x%Ce$^{3+}$ with x = (a) 1, (b) 2 and (c) 3 for excitation at 454 nm and 340 nm, respectively; see Fig. S2 (ESI†) for luminescence spectra with respect to the one of the excitation source.

Fig. 4  (a) Diffuse reflectance/transmittance spectra of YAG:x%Ce$^{3+}$ ($x = 0.2$, 1, 2 and 3). The spectra have been vertically offset and smoothed (see Fig. S3, ESI† for the raw spectral). (b and c) Emission spectra of YAG:x%Ce$^{3+}$ ($x = 0.2$, 1, 2 and 3) for excitation at 340 and 454 nm, respectively. The spectra of the respective excitation sources, normalized to the band maxima, are shown by the grey areas in (a).
thermal stability of the PLQY (estimated from $T_{50\%}$) for excitation at 340 nm is higher than that at 454 nm, cf. Fig. 2(c). Since the photon irradiation at 340 nm corresponds to excitation to the $5d_2$ level, which is closer to the CB minimum than the $5d_1$ level, this should cause stronger thermal ionization at a given, sufficiently high, temperature and should be reflected in a lower thermal stability of PLQY. In effect, the reverse behaviour as observed here suggests that thermal ionization is not the predominating quenching process for these higher Ce$^{3+}$ concentrations (1–3 mol%). Rather, the systematic decrease of $T_{50\%}$ with increasing Ce$^{3+}$ concentration [Fig. 2(c)] suggests that either thermally activated energy transfer processes to luminescence killer centers, or/and Ce$^{3+}$ concentration induced nonradiative $5d \rightarrow 4f$ crossover-relaxation processes are at play.

2.3 Optical spectra and Ce–Ce distance analysis

As far as energy transfer processes between different Ce$^{3+}$ ions are concerned, their probability is thought to depend on the overlap between the absorption and emission bands and the distances between Ce$^{3+}$ ions, see Fig. 5(a and b). In the case of YAG:Ce$^{3+}$, the spectral overlap is as high as about 5% of the absorption/emission spectrum at RT [Fig. 5(c)]. It implies that 5% of excited energy has a certain probability of being energetically transferred to other Ce$^{3+}$ ions in YAG:Ce$^{3+}$.

Physically, energy transfer may occur as a consequence of electric dipole–dipole interactions of Ce$^{3+}$ ions with a rate $P_{oe}$. The critical distance $R_c$ for such energy transfer processes is defined as the distance between Ce$^{3+}$ ions for which $P_{oe}$ is equal to the radiative rate of Ce$^{3+}$. $R_c$ can be estimated using the following equation,10–12

$$R_c^6 = (6.3 \times 10^{27}) Q_{abs} \int \frac{F_{abs}(E) f_{emi}(E)}{E^4} dE. \quad (2)$$

Here, $Q_{abs}$ is the absorption strength of Ce$^{3+}$ and takes on a value of $5.55 \times 10^{-19}$ cm$^2$ eV$^{-1}$ as we have derived from the integrated absorption cross section of Ce$^{3+}$ (see Fig. S4, ESI†). $F_{abs}(E)$ and $f_{emi}(E)$ are the absorption and emission spectral bands that are normalized to unity, respectively, and $E$ is the photon energy.

Using the spectral data in Fig. 5(a), we obtain $R_c = 13$ Å. This value is in agreement with the one reported by Blasse12 on the basis of the same model [eqn (2)]. This implies that when the Ce$^{3+}$ concentration is higher than 0.7 mol% [Fig. 5(d)], the effect of energy-transfer induced concentration quenching is comparable with radiative transitions of Ce$^{3+}$. In addition, the energy transfer rate is found to increase upon elevating temperature due to the twin effect of an increased spectral overlap and a red-shift of the luminescence spectra,30 see Fig. S5 (ESI†). Therefore, we infer that in order to suppress thermally activated concentration quenching, the Ce$^{3+}$ concentration should be far lower than 0.7 mol%. A lowering of the Ce$^{3+}$ concentration can be also believed to reduce the amount of Ce$^{3+}$ induced defects such as anion vacancies and local structural distortions,7–9,23,31 which may act as luminescence killer centers, cf. Fig. 1(b), in a manner of, e.g., charge trapping.

Information about the defect structure and charge-trapping dynamics in YAG:Ce$^{3+}$ can be obtained from TL experiments.

2.4 Thermoluminescence glow curve measurements

Fig. 6 shows the TL glow curves for YAG:1–3%Ce$^{3+}$, as a function of excitation temperature $T_{ex}$ between 280 K and 700 K. For each material, we observe three distinct glow curve bands when $T_{ex} = 280$ K and their positions and intensity change significantly upon varying $T_{ex}$ see Fig. 6 top and bottom, respectively. When the intensity of the glow curve bands reaches their maxima at variable $T_{ex}$ the positions of the bands are determined, which are for example centred at
approximately \( T_m = 325, 415, \) and 560 K for 3% Ce\(^{3+}\), respectively, where \( T_m \) is referred to as the temperature corresponding to the maximum of a glow curve band. The glow curve bands correspond to different types of charge-trapping defects, here labelled as \( \text{trap1}, \text{trap2}, \) and \( \text{trap3} \), respectively. On the basis of the \( T_m \) values and first-order TL kinetics, we determine the energy difference between the respective trap and CB (here called the trap depth \( E_t \)) according to

\[
\beta E_t / k_B T_m = s \exp(-E_t / k_B T_m),
\]

where \( \beta \) is the heating rate \((2 \text{ K s}^{-1})\), \( s \) is a frequency factor \((1 \times 10^{11} \text{ s}^{-1} \text{ for YAG:Ce}^{3+})\),\(^{20} \) and \( k_B \) is the Boltzmann constant. For example, for 3% Ce\(^{3+}\), \( E_t \) takes on values of 0.76 eV (\( \text{trap1} \)), 0.98 eV (\( \text{trap2} \)) and 1.33 eV (\( \text{trap3} \)), respectively, cf. Fig. 7(a).

The activation energy for charge filling the respective trap, \( \Delta E_{\text{TL}} \), can be extracted from the \( 1/T \text{ex} \) dependence of the integrated intensity of the TL glow curve bands, where \( T \text{ex} \) should be below the temperature at which the trap is filled to each maximum. This analysis points toward a strict Arrhenius dependence with \( \Delta E_{\text{TL}} \) taking on values between 0.15 and 0.32 eV [Fig. 7(b)]. Since all \( \Delta E_{\text{TL}} \) values are significantly smaller than the activation energy for thermal ionization of the Ce\(^{3+} 5d_1\) electron into the CB \((\Delta E_{\text{i,5d1}} \approx 1.1 \text{ eV}, \text{ see Table 1})\), this indicates a higher probability of direct charge migration/trapping from the Ce\(^{3+}\) ions to defect states, in comparison with charge trapping through thermal ionization.

Interestingly, the energy position of \( \text{trap3} \) (1.3–1.6 eV below the CB) is lower than the Ce\(^{3+} 5d_1\) level (1.1 eV below the CB), which implies that the process of de-trapping charges at \( \text{trap3} \) and returning to the 5d\(_1\) level is energetically unfavourable due to the relatively large energy barrier of 0.2–0.5 eV. Moreover, we observe that for \( \text{trap3} \), which exhibits the strongest TL intensity among the three traps and, hence, the analysis leads to the most reliable results, \( \Delta E_{\text{TL}} \) decreases systematically with increasing Ce\(^{3+}\) concentration. In effect, this means that for the higher Ce\(^{3+}\) concentrations, the process of trapping the migrating charges of Ce\(^{3+}\) at \( \text{trap3} \) becomes more thermally
active at a given temperature, which is very likely attributed to enhanced energy transfer among Ce$^{3+}$ ions. In this scenario, we use $\Delta E_{\text{TL}}$ obtained from trap$3$ as the activation energy for thermally activated concentration quenching.

The number of charges in each trap has been also determined from the difference of the TL intensity between two adjacent $T_{\text{ex}}$, where $T_{\text{ex}}$ is higher than the temperature resulting in the strongest TL intensity, i.e. when $T_{\text{ex}}$ is sufficiently high to begin to empty the trapped charges from the trap states. From the combined analyses of $E_t$ and the number of trapped charges associated to trap$3$, the trap density of states (TDOS), representing the number of states (population) with respect to $E_t$, has been determined [Fig. 7(a)]. The TDOS of the deepest trap, trap$3$, can be approximated with a Gaussian distribution for each sample. The bandwidth of the TDOS distribution of trap$3$ generally increases upon increasing Ce$^{3+}$ concentration from 1 to 3 mol%. For the shallower traps, trap$2$ and trap$1$, the TDOS distribution cannot be determined because of the limited amount of data making the fitting analysis unreliable.

In comparison with the shallower defects that are often associated with non-vacancy defects such as antisite defects, the origin of the deepest trap trap$3$ is most probably related to oxygen vacancies, which are known to act as electron traps and are energetically favourable to be formed in a reducing environment, as in the case of our sample preparation. This statement is further supported by the vanished TL glow in YAG:Ce$^{3+}$ after high-temperature annealing in air. We thus infer that the broadening of the TDOS distribution for trap$3$ upon increasing Ce$^{3+}$ concentration may be caused by increased local structural disorder and distortions nearby the Ce$^{3+}$ ions. This local structural effect changes the coordination environments of the O vacancies and thus leads to a larger variation of the energy levels of the traps related to these vacancies. This effect becomes more pronounced as the Ce$^{3+}$ ions replace the smaller Y$^{3+}$ ions in the neighbourhood of O vacancies. As a result, this TDOS broadening effect may enhance the interactions between Ce$^{3+}$ ions and trap states and further increase the probability for charges to be trapped, which leads to the effect of lowering $\Delta E_{\text{TL}}$. We infer that the combined observations of a systematically increasing energy transfer rate (Fig. 5) and a broadening of the TDOS (Fig. 7), with a simultaneously decreasing $T_{\text{300}}$ [Fig. 2(c)] as a function of increasing Ce$^{3+}$ concentration, confirms the contribution from concentration quenching to the overall quenching behaviour of YAG:Ce$^{3+}$. However, we also notice the decreasing Debye temperature $\Theta_D$ upon increasing the Ce$^{3+}$ concentration from 0 to 3 mol%, which reflects a general softening of the YAG lattice. As $\Theta_D$ is a key measure for thermal stability/ downconversion efficiency of luminescence, the apparent correlation between $\Theta_D$ and $T_{\text{300}}$ implies the involvement of nonradiative 5d $\rightarrow$ 4f crossover relaxation in YAG:Ce$^{3+}$. We investigated this process using mode-selective vibrational excitation techniques.

### 2.5 Mode-selective vibrational excitation experiments

Fig. 8 summarizes the results from the two-laser experiment combining a pulsed tunable infrared (IR) laser from FELIX with a pulsed 454 nm blue laser. The aim of this experiment is to investigate the effect of coherent large-amplitude excitation of IR active vibrational modes (phonons) on the luminescence decay time, $\tau$, and thus to get information about possible electron–phonon coupling mechanisms associated with nonradiative crossover relaxation processes. Fig. 8(a) shows $\tau$ as a function of the wavenumber of the IR irradiation over the range $\approx 460$–900 cm$^{-1}$ for YAG:3%Ce$^{3+}$. The lower part of the figure shows the IR spectrum of YAG:3%Ce$^{3+}$, as well as the total energy of an IR macro-pulse; Fig. S7 (ESI†) shows a scheme of the experimental setup. As can be observed, $\tau$ is overall shortened, from a few nanoseconds up to several nanoseconds as a function of increasing temperature from 300 K to 450 K, especially in the high wavenumber region (650–900 cm$^{-1}$) where the IR irradiation is relatively strong. The solid lines in Fig. 8(a) are the simulated $\tau$ estimated by considering (only) the heating effect of the IR pre-pulse excitation, as discussed in the Methods section. The simulated $\tau$ mimics quite well the energy profile of the IR macro-pulse, with a relatively weak effect in the regions of low IR absorbance, e.g. 600–650 cm$^{-1}$ and 850–900 cm$^{-1}$.

Fig. 8(b) shows the experimentally determined $\tau$ subtracted by the simulated $\tau$, here denoted as $\Delta \tau$, which is the response of the mode-selective excitation of IR active phonons on $\tau$. Clearly, the selective excitation of the three highest-frequency phonons at 698, 724, and 789 cm$^{-1}$, which are assigned to different asymmetric bending motions of the local CeO$_6$ dodecahedra, depopulates the light-emitting levels of Ce$^{3+}$ in a nonradiative way. Moreover, the magnitude of $\Delta \tau$ increases progressively with increasing temperature of the sample.

To further investigate this temperature dependence, we take the average of $\Delta \tau$, here denoted as $\Delta \tau_{\text{ave}}$, over the wavenumber range 650–860 cm$^{-1}$. $\Delta \tau_{\text{ave}}$ follows an Arrhenius dependence with a characteristic activation energy $\Delta E_{\text{ph}} = 0.23$ eV, see Fig. 8(c) and Table 1. A priori, $\Delta E_{\text{ph}}$ could relate to either one or a combination of the three possible thermal quenching processes (Fig. 1). However, because $\Delta E_{\text{ph}}$ is far lower than $\Delta E_{5d1}$, vibrationally stimulated ionization of Ce$^{3+}$ ions seems unlikely. In contrast, $\Delta E_{\text{ph}}$ is closer to $\Delta E_{\text{TL}}$, suggesting that $\Delta E_{\text{ph}}$ relates to vibrationally stimulated concentration quenching. In addition, it is also possible to relate to vibrationally stimulated 5d $\rightarrow$ 4f crossover via electron–phonon interactions in the material. Based on the 5d $\rightarrow$ 4f crossover model, $\Delta E_{\text{ph}}$ thus corresponds to the energy barrier for thermal quenching activated by mode-selective vibrational excitation, see Fig. 1(c).

One may note that, whereas the activation of vibrational modes for 5d $\rightarrow$ 4f crossover relaxation is obvious, concentration quenching should as well be promoted by the increased probability of energy resonance between Ce$^{3+}$ ions with the help of vibrational excitation. Further, we note that the heating of the sample, due to the increase of applied temperature, has the effect of activating the modes in the low wavenumber region ($<300$ cm$^{-1}$), which primarily relate to localized motions of Ce$^{3+}$ ions. This suggests that the vibrationally induced thermal quenching may need the excitation of multiple types of phonons.
2.6 Overall picture of the thermal quenching mechanisms

To sum up the results so far, the three quenching processes appear to be all involved in thermal quenching of luminescence, at least for YAG:3%Ce³⁺ since the magnitude of $D_E = 0.32 \text{ eV}$ is between $D_{E_{i,5d1}} = 1.13 \text{ eV}$ and $D_{E_{TL}} = 0.15 \text{ eV}$ or $D_{E_{ph}} = 0.23 \text{ eV}$. Under the assumptions that the radiative rate $G_r$ is temperature independent, and that the nonradiative rates of the three quenching processes increase exponentially with temperature, eqn (1) can be modified as follows, by

$$\tau(T) = 1/[G_r + G_{n1} \exp(-E_{n1}/k_BT) + G_{n2} \exp(-E_{n2}/k_BT)]. \quad (4)$$

Here, $G_{n1}$ is the nonradiative rate for thermal ionization, and $G_{n2}$ reflects a combined nonradiative rate for concentration quenching and 5d → 4f crossover. Note that the latter two processes are here treated as one, because their activation energies are too close to each other to be distinguished. Therefore, $E_{n1}$ equals $D_{E_{5d1}}$ and $E_{n2}$ equals the average of $D_{E_{TL}}$ and $D_{E_{ph}}$. A free fit of eqn (4) to $\tau(T)$ of YAG:3%Ce³⁺ [Fig. 2[a]] yields $G_r = 1.61 \times 10^7 \text{ s}^{-1}$, $G_{n1} = 4.24 \times 10^{15} \text{ s}^{-1}$ and $G_{n2} = 7.72 \times 10^8 \text{ s}^{-1}$. Using these results, we have calculated the temperature dependence of the ratios of the radiative/nonradiative transition rates to the total transition rate. The result shows that the temperature dependence of the calculated ratio for the radiative transition is in good agreement with that of the experimental $\tau(T)$ (Fig. 9). Most importantly, we observe that the combined effect of concentration quenching and 5d → 4f crossover quenching contributes predominantly to the thermal quenching up to about 600 K, whereas at higher temperatures it...
is dictated by thermal ionization. Interestingly, a similar result is obtained for YAG:1%Ce$^{3+}$ [Fig. 9(b)] using the same method under the assumption that $\Delta E_{\text{ph}}$ is the same for 1 and 3 mol% Ce$^{3+}$ dopant concentrations. However, it is noticed that the most significant and important difference between the two samples is the contribution from combined concentration and crossover quenching, which clearly shows that lowering the Ce$^{3+}$ concentration improves greatly the thermal resistance against thermal quenching especially when the operating temperature is below 600 K, see Fig. 9.

2.7 Generation and evaluation of artificial white light

Emission spectra of artificial white light generated from a blue LED and YAG:x%Ce$^{3+}$ (x = 1 and 3) at the temperatures 80 K, 600 K, and 860 K, converted from the emission spectra to the right.

![Fig. 10](image.png)

Fig. 10: Left: CIE 1931 coordinates of artificial white light (inset) generated from a 450 nm blue LED and YAG:x%Ce$^{3+}$ (x = 1 and 3) at the temperatures 80 K, 600 K, and 860 K, converted from the emission spectra to the right.

is found to be much stronger for the higher Ce$^{3+}$ coordinates in Fig. 10 left. Especially, upon increasing the bluish white towards pure blue, as shown by their CIE 1931 coordinates in Fig. 10 left. The color-rendering index (CRI) of the artificial white light luminescence increases greatly with the decrease in intensity in the green-yellow region (500–600 nm) upon elevating temperature from 80 K to 860 K (Fig. 10 right).

Thermal quenching is often a major concern for artificial white light phosphors and is tied to the presence of defects in the host crystal lattice. A general strategy of enhancing the structural rigidity by decreasing the Ce$^{3+}$ concentration shows the potential to improve the resistance against thermal quenching of luminescence, since fewer phonon modes are activated. In particular, the activation of high-frequency phonons, which are mainly associated with asymmetric bending motions of the local CeO$_6$ moieties, shows the effect of inducing nonradiative relaxation through most probably 5d $\rightarrow$ 4f crossover processes. For temperatures higher than 600 K, the thermal quenching is instead primarily dictated by thermal ionization, which can be suppressed by increasing the ionization energy via different cation co-substitution strategies.
4.3 Photoluminescence measurements

PL emission spectra of the YAG:Ce\(^{3+}\) powder samples were measured for temperatures between 80 K and 860 K, using the same setup as used for the diffuse reflectance spectroscopy measurements (see above) but with different excitation sources. For excitation in the blue region, a pulsed laser (DeltaDiode-450L, HORIBA Scientific), with a peak wavelength at 454 nm and a repetition rate of 100 MHz and a pulse width of 80 ps, was used. For excitation in the UV region, we used a continuous-wave UV LED (M340D3, Thorlabs) with a nominal wavelength of 340 nm, powered by 5 V. Note that the emission spectra of YAG:0.2\%Ce\(^{3+}\) were measured using an optical fiber probe close to the sample instead of using the integrating sphere that led to very weak emission intensity. For PL decay curve measurements, also at temperatures between 80 K and 860 K, we used the pulsed 454 nm laser (DeltaDiode-450L, HORIBA Scientific), but with the repetition rate adjusted to 500 kHz. For the detection of the emitted light, we used a photosensor (H10721-20, Hamamatsu) coupled to a 500 nm longpass filter (A10033-62, Hamamatsu). The temperature was controlled by a Linkam THMS 600 stage.

4.4 Ce–Ce distances analysis

Ce–Ce distances in YAG:Ce\(^{3+}\) were determined by calculating the probability distribution of the distances between two Ce\(^{3+}\) ions that occupy any on two Y sites (with an equal probability of site occupation) in a super cell comprising \(n\) YAG unit cells, where \(n = 1, 2, 3, \ldots\) and 7 correspond to Ce\(^{3+}\) concentrations of 8.3, 0.1, 0.31, \ldots\) and 0.024 mol%, respectively [Fig. 5(b)] and using crystallographic data of YAG as published in ref. 42. The mean value of the probability distribution of the Ce\(^{3+}\) distances is taken as the average distance between Ce\(^{3+}\) ions (\(R_{\text{Ce}}\)) for a corresponding Ce\(^{3+}\) concentration. Using the relation between \(R_{\text{Ce}}\) and Ce\(^{3+}\) concentration, which can be well fitted to a power function, we can estimate \(R_{\text{Ce}}\) for any specific Ce\(^{3+}\) concentration in YAG:Ce\(^{3+}\), see Fig. 5(b) and Table 1. Interestingly, the result using this combined structural and mathematical model is very similar to the one derived from the relation 

\[
R_{\text{Ce}} = \left[ \frac{3}{4(4n)} \right]^{1/3}, \quad \text{where } N \text{ is the density of Ce}^{3+} \text{ in YAG:Ce}^{3+}.
\]

Using the same crystallographic data as above,\(^{42}\) for which \(N = 1.3867 \times 10^{26} \text{ ions m}^{-3}\) for 1\% Ce\(^{3+}\) that scales linearly with Ce\(^{3+}\) concentration, we obtain \(R_{\text{Ce}} = 37.4, 32.5, 20.5, 15.1, 12.0, 9.5\) and 8.3 Å for 0.033, 0.05, 0.2, 0.5, 1, 2 and 3 mol\% Ce\(^{3+}\) doping, respectively (see Table 1 for comparison).

4.5 Thermoluminescence glow curve measurements

For the TL measurements, the excitation of charge storage was achieved by photon irradiation at different excitation temperatures (\(T_{\text{ex}} = 280–700\) K), using a continuous-wave blue LED at 450 nm for 2 min, which was immediately followed by fast cooling (100 K min\(^{-1}\)) to 280 K. After a thermal equilibration time of 1 min (at 280 K), TL glow curves were measured by detecting the emitted light as a function of increasing temperature from 280 K to 700 K with a heating rate of 2 K s\(^{-1}\). The TL signal was monitored using a photosensor (H10721-20, Hamamatsu) and an oscilloscope (DSO-X 2022A, Agilent Technologies). TL glow curves unbiased of the effect of thermal quenching of luminescence were determined by taking the ratio of the measured TL signal and the temperature dependent luminescence intensity.\(^{20}\)

The samples, in the form of powders, were held in crucibles of aluminum and their temperature was controlled using a Linkam THMS 600 heating stage.

4.6 Mode-selective vibrational excitation measurements

The mode-selective vibrational excitation experiments were performed at the free electron laser facility FELIX at Radboud University, The Netherlands.\(^{45}\) FELIX generates a pulsed, monochromatic beam of photons, which is tuneable in a wide IR range, e.g. between 100 cm\(^{-1}\) and 3600 cm\(^{-1}\), thus making it possible to excite selectively IR active lattice vibrational modes (phonons) of YAG:Ce\(^{3+}\).\(^{16}\) For this purpose, the wavelength of the IR light was scanned over the range 11–22 \(\mu\)m (\(\approx 460–900\) cm\(^{-1}\)) with a step size of 0.1 \(\mu\)m, while detecting in situ the PL decay curve of YAG:3\%Ce\(^{3+}\) using a photosensor (H10721-20, Hamamatsu) and a 500 nm long-pass filter in front upon pulsed excitation at 454 nm (DeltaDiode-450L, HORIBA Scientific), see Fig. S7 (ESI\(^{†}\)) for the experimental setup. In this way, features in the PL decay characteristics specifically associated with the excitation of specific IR active phonons were measured. Of importance here, the FELIX light pulses consist of macro-pulses with a pulse width of 5–10 \(\mu\)s and a repetition rate of 10 Hz, and for which each macro-pulse consists of a 1 GHz train of micro-pulses of a few ps duration (Fig. S7, ESI\(^{†}\)). The 454 nm excitation was triggered in the middle of each macro-pulse.

The sample was prepared as a disc-shaped pellet containing approximately 20 wt\% of YAG:3\%Ce\(^{3+}\) in 50 mg CsI, and was held in a custom-made heating block. It should be noted that CsI is transparent in the wavenumber range of interest (460–900 cm\(^{-1}\)) and hence does not contribute to any peaks in the vibrational spectrum. CsI is also transparent in the visible region and hence does not contribute to any effects related to the 454 nm excitation and the emission of YAG:Ce\(^{3+}\).

To suppress the decrease of the intensity of the IR light associated with the absorption of air molecules, the measurements were performed inside a custom-made enclosure with flowing N\(_{2}\). Measurements were taken at the temperatures 300 K, 350 K, 380 K, 400 K and 450 K.

4.7 Thermal simulations

In order to determine the effect of mode-selective vibrational excitation on the PL decay curve characteristics of YAG:3\%Ce\(^{3+}\), one has to account for any local (pre-)heating effects due to the IR irradiation on the sample. To evaluate such heating effects, we simulated the temperature of YAG:3\%Ce\(^{3+}\), by mimicking the experimental conditions. Specifically, we considered the effects of

(i) temperature increase due to the mode-selective vibrational excitation by the irradiation with monochromatic IR light, and

(ii) temperature decrease through spontaneous heat dissipation in the sample, as an effect of any thermal gradient in the system (sample + surrounding).
To determine the spatial and temporal distribution of thermal energy resulting from (i) and (ii), respectively, the bulk of the disc-shaped sample, made of YAG:3%Ce3+ and CsI (diluting agent), whose dimensions measured approximately 7 mm ϕ (xy-plane) × 0.35 mm thickness (z-direction), was simulated by meshed grids of 0.2 mm (x-axis) × 0.2 mm (y-axis) × 0.04 mm (z-axis), coordinated at a position (x, y, z) in the sample, and where the thickness of the sample was calculated by assuming that the mixture of YAG:3%Ce3+ and CsI was an ideal solid solution. The respective densities (ρ) for YAG:Ce3+ and CsI (4.56 and 4.51 mg mm−3, respectively) were assumed to be temperature independent, and CsI was assumed to be completely transparent in the IR and visible regions.

For the simulation of the IR irradiation, the IR laser beam was assumed to be irradiating in the vertical direction (i.e. in the z-direction) on the sample. The cross section of the intensity of the IR beam (xy-plane) was set to be a 2D Gaussian distribution with a full width at half-maximum (FWHM) of 1 mm (Fig. S7, ESI†) with an intensity that decreases along the z-direction according to the Beer–Lambert law,28 which was determined by the IR absorbance of YAG:3%Ce3+ (Fig. S8, ESI†). An IR macro-pulse was then simulated by ten pulses of 1 μs duration, separated by an infinitesimal time interval. Together with the measured energy of one IR macro-pulse [Fig. 8(a)] bottom, a 3D energy distribution of the simulated IR macro-pulse at any specific time (t) in the interval of 0–10 μs inside the sample at any variable position, EIR(x,y,z,t), was calculated.

By using this computational setting, the thermal effect due to mode-selective vibrational excitation (i) was simulated by48

\[
\frac{dT}{dt} = \frac{E_{IR}(x, y, z, t)}{C(T)m},
\]

where T is the temperature of the sample unit (i.e. meshed grid) at (x, y, z), m is the mass of YAG:Ce3+ contained in the sample unit and C(T) is the heat capacity of YAG:Ce3+, see Fig. S9 (ESI†).49 For the thermal effect of the spontaneous heat dissipation in the sample (ii), we instead used the following equation based on the Fourier's law and the law of energy conservation,48,50

\[
\frac{dT}{dt} = \frac{k_{eff}(T)}{\rho_{eff}C_{eff}(T)} \nabla^2 T.
\]

Here, \(k_{eff}(T)\), \(C_{eff}(T)\) and \(\rho_{eff}\) are the effective thermal conductivity, heat capacity and density of the sample, respectively (Fig. S9, ESI†).49,51,52 The effective thermal and materials properties were calculated by weighing the same properties of YAG:Ce3+ and CsI based on their masses and considering their average effect on heat conduction, because the diffusion of thermal energy acted on both materials in the sample.

Using as input in the simulations \(T = 300\) K, 350 K, 380 K, 400 K, and 450 K, that matched the experimental conditions, eqn (5) and (6) were solved. In effect, we obtained the evolution of temperature of the sample units throughout one IR macro-pulse, which could be further spatially averaged to represent the in situ temperature of the entire sample. Of particular interest, the in situ temperature at the time when the sample was excited by the 454 nm blue laser (i.e. after about 5 μs of IR irradiation, see Fig. S7, ESI†) was determined and the result was plotted as a function of the IR wavenumber (Fig. S6, ESI†). From this result, the temperature increment \(\Delta T\) averaged over the range of 650–860 cm\(^{-1}\), where the thermal effect of the mode-selective vibrational excitation was most effective, was calculated and used for the estimation of non-thermal effects of mode-selective vibrational excitation [Fig. 8(c)].

The heating effect due to the IR irradiation (as described above) on the luminescence decay time of YAG:Ce3+ was further simulated. In each PL detection, the decay curve was assumed to be the result of an averaging effect of the light emitted from YAG:Ce3+ powder particles that were distributed in different parts of the pellet sample, which were characterised by different temperatures, as discussed above. Therefore, these YAG:Ce3+ particles corresponded to different τ values that could be extracted from the experimental τ–T relation [Fig. 2(a)]. Furthermore, the contributions from YAG:Ce3+ particles situated at different positions needed to be weighed by the absorption of the 454 nm laser, i.e. the energy distribution of the 454 nm laser beam inside the sample. This energy distribution was defined as a 2D Gaussian distribution with a FWHM of 2 mm (z-direction) and 0.7 mm (y-direction) (Fig. S7, ESI†), whose intensity decreased along the z-direction based on the Beer–Lambert law. The absorption cross section of the sample at 454 nm was experimentally measured to be \(2.95 \times 10^{-20}\) cm². Although this cross section value was relatively small compared to the one obtained from the single crystal film of YAG:0.2%Ce3+ (Fig. S4, ESI†), it should reflect better the real sample condition, e.g. light scattering events in a powder made sample. To this end, the average of weighed τ values was used to represent the simulated τ varied upon the IR heating effect [Fig. 8(a)].

4.8 Generation and evaluation of artificial white light

Artificial white light was generated by the irradiation with a blue LED (450 nm, output optical power = 1 mW) on YAG:Ce3+, at the temperatures 80 K, 600 K, and 860 K. The emitted white light was detected using an integrating sphere (IS200-4, Thorlabs) connected to an UV-VIS spectrometer (USB2000+, Ocean Optics). The temperature of YAG:Ce3+ was set by a Linkam THMS 600 heating stage.

Conflicts of interest

There are no conflicts to declare.

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