



Exact results for the six-vertex model with domain wall boundary conditions and a partially reflecting end

Downloaded from: <https://research.chalmers.se>, 2022-06-26 15:13 UTC

Citation for the original published paper (version of record):

Hietala, L. (2022). Exact results for the six-vertex model with domain wall boundary conditions and a partially reflecting end. *Letters in Mathematical Physics*, 112(2).
<http://dx.doi.org/10.1007/s11005-022-01530-5>

N.B. When citing this work, cite the original published paper.



Exact results for the six-vertex model with domain wall boundary conditions and a partially reflecting end

Linnea Hietala^{1,2}

Received: 26 November 2021 / Revised: 21 February 2022 / Accepted: 23 February 2022
© The Author(s) 2022

Abstract

The trigonometric six-vertex model with domain wall boundary conditions and one partially reflecting end on a lattice of size $2n \times m$, $m \leq n$, is considered. The partition function is computed using the Izergin–Korepin method, generalizing the result of Foda and Zarembo from the rational to the trigonometric case. Thereafter, we specify the parameters in Kuperberg’s way to get a formula for the number of states as a determinant of Wilson polynomials. We relate this to a new type of alternating sign matrices, similar to how the six-vertex model with domain wall boundary conditions is related to normal alternating sign matrices. In an appendix, we compute the partition function again, showing that it is also possible to find it with the method of Foda and Wheeler.

Keywords Six-vertex model · Domain wall boundary conditions · Reflecting end · Partition function · Triangular K -matrix

Mathematics Subject Classification 82B23 · 05A15 · 33C45

1 Introduction

The first example of a six-vertex (6V) model was the ice-model, where all states have the same weight. This and some other special cases of the 6V model with periodic

Affiliation 1: I no longer work at this department.

✉ Linnea Hietala
linnea.hietala@math.uu.se

¹ Department of Mathematical Sciences, Chalmers University of Technology and University of Gothenburg, 412 96 Gothenburg, Sweden

² Department of Mathematics, Uppsala University, Box 480, 751 06 Uppsala, Sweden

boundary conditions were solved in 1967 by Lieb [1]. The same year, Sutherland [2] solved the general case.

One of the first nontrivial examples of fixed boundaries were the domain wall boundary conditions (DWBC) [3]. In 1996, Zeilberger [4] proved the alternating sign matrix conjecture of Mills, Robbins and Rumsey [5], which gives a formula for the number of alternating sign matrices (ASMs). There is a bijection between the ASMs and the states of the 6V model with DWBC. Izergin [6, 7] showed that the partition function of the 6V model with DWBC can be expressed as a determinant, which Kuperberg [8] used to give another proof of the alternating sign matrix conjecture.

Tsuchiya [9] used the Izergin–Korepin method to obtain a determinant formula for the partition function of the 6V model with one diagonal reflecting end and DWBC on the three other sides on a lattice of size $2n \times n$. Kuperberg [10] used this to give a formula for the number of the corresponding UASMs. The UASMs are alternating sign matrices with U-turns on one side and generalize the vertically symmetric alternating sign matrices (VSASMs).

Foda and Wheeler [11] found a determinant formula for the partition function of the 6V model with partial DWBC on a lattice of size $m \times n$, which generalizes the determinant formula of Korepin and Izergin. Foda and Zarembo [12] found the corresponding generalization of Tsuchiya's determinant formula in the rational case. They obtained a determinant formula for the rational 6V model on a lattice with $2n \times m$ sites, $m \leq n$, with DWBC and where the reflecting end has a triangular K -matrix. These boundary conditions are called DWBC with a partially reflecting end. Pozsgay [13] used the homogeneous limit of Tsuchiya's $2n \times n$ determinant to compute overlaps (i.e. inner products) between (off-shell) Bethe states and certain simple product states, such as the Néel states. In a similar way, Foda and Zarembo used their rational $2n \times m$ determinant formula to compute overlaps between Bethe states and more general objects which they call partial Néel states.

Foda and Wheeler commented that in the case of partial DWBC on a lattice of size $n \times m$, it is not obvious if and how one could count ASM-like objects with Kuperberg's specialization, due to phases that vary between different states, coming from the trigonometric weights. However, in the present paper, we find that in the case of DWBC and partial reflection, it is possible to count the states, since similar phases do not appear in this case. The reason for this is the alternating orientations of the lines.

Counting ASMs can be generalized to x -enumerations. In an x -enumeration of ASMs, each state is counted with a weight x^k , where k is the number of -1 's in the ASM. A formula in the general case is not known, but in some special cases, $x = 1, 2$ and 3 , there are closed expressions [4, 5, 8]. Colomo and Pronko [14] obtained a simplified treatment of x -enumerations by rewriting the Hankel determinant representation of the partition function of the 6V model with DWBC in terms of orthogonal polynomials. The method can be used to find a formula for the x -enumerations for those x where the underlying orthogonal polynomials belong to the Askey scheme of hypergeometric orthogonal polynomials. This works for the 1-, 2- and 3-enumerations. In the original case of 1-enumerations, the orthogonal polynomials are continuous Hahn polynomials.

In this paper, we study the trigonometric 6V model with DWBC on three sides and one partially reflecting end, on a lattice with $2n \times m$ sites, $m \leq n$. We first find a determinant formula for the partition function. Then we specialize the parameters in Kuperberg’s manner to finally find a formula that counts the number of states of the model.

At first, in Sect. 2, we introduce the model. In Sect. 3 we follow the Izergin–Korepin method to obtain a determinant formula for the partition function in Theorem 3.5, i.e. the trigonometric generalization of what Foda and Zarembo did in the rational case. The alternative method of Foda and Wheeler to find the determinant formula is presented in “Appendix A”.

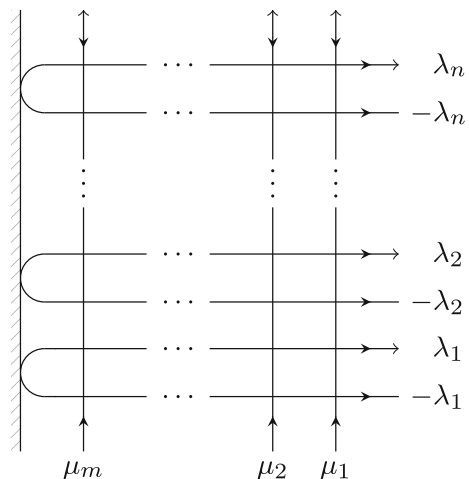
In Sect. 4, we find a formula for the number of states in terms of the partition function. We connect this to the enumeration of a type of generalized UASMs. The objective of Sect. 5 is to specialize the parameters in the determinant formula in Kuperberg’s way. We rewrite the partition function following the ideas of Colomo and Pronko [14]. The determinant can be represented by a matrix consisting of a Hankel matrix part and a Vandermonde matrix part, and the underlying orthogonal polynomials are Wilson polynomials. Then in Theorem 5.1, we finally write down a determinant formula counting the number of states of the 6V model with DWBC and one partially reflecting end.

2 Preliminaries

Consider a square lattice with $2n \times m$ lines, where the horizontal lines are connected pairwise at the left side, as in Fig. 1. Each such pair of horizontal lines can be thought of as one single line turning at a wall on the left side. We assign a spin ± 1 to each edge. A lattice with a spin assigned to each edge is called a state.

In order to assign weights to the states, we give each line an orientation. We choose a positive direction, which goes upwards for the vertical lines, to the left for the lower part

Fig. 1 The 6V model with DWBC and one partially reflecting end. The parameters λ_i and μ_j are the spectral parameters



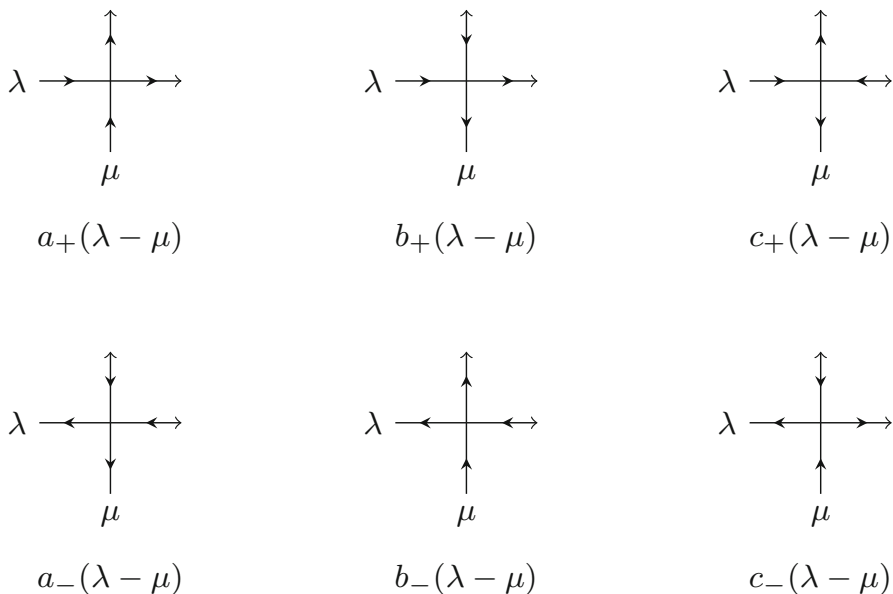


Fig. 2 The possible vertices and their vertex weights for the 6V model. The orientation of the lines through a vertex is indicated with an arrow in the end of each line. Given the orientation up and to the right, the spins are indicated by an arrow halfway the edge, where right and up are positive spins, and left and down are negative spins. The vertex weights also depend on the spectral parameters λ and μ . In a state of the 6V model with reflecting end, the oriented vertices are tilted 90 degrees counterclockwise on every second row

of the horizontal double line, and to the right for the upper part. The positive direction is indicated by an arrow at the end of a line. Graphically, spin +1 corresponds to an arrow pointing in the positive direction of the line, and spin -1 corresponds to an arrow pointing in the opposite direction. At each vertex, the so-called ice rule must hold, which demands that two arrows must be pointing inwards to the vertex and two arrows must be pointing outwards. Because of the ice rule, there are only six types of possible vertices, see Fig. 2.

To each vertical line, we assign a spectral parameter μ_j , and to each horizontal double line, we assign a spectral parameter which is $-\lambda_i$ on the lower part of the double line and shifts to λ_i on the upper part. In Fig. 1, we write these parameters at the lines. Also define a fixed boundary parameter $\zeta \in \mathbb{C}$, associated with the reflecting wall at the turns.

Define $f(x) = 2 \sinh(x)$ and let $\gamma \notin 2\pi i\mathbb{Z}$ be a fixed parameter. Then define local weights

$$a_{\pm}(\lambda) = 1, \quad b_{\pm}(\lambda) = e^{\mp\gamma} \frac{f(\lambda)}{f(\lambda + \gamma)}, \quad c_{\pm}(\lambda) = e^{\pm\lambda} \frac{f(\gamma)}{f(\lambda + \gamma)}, \quad (1)$$

$$k_{\pm}(\lambda, \zeta) = e^{\zeta \mp \lambda} f(\zeta \pm \lambda), \quad k_c(\lambda, \zeta) = \varphi f(2\lambda), \quad (2)$$

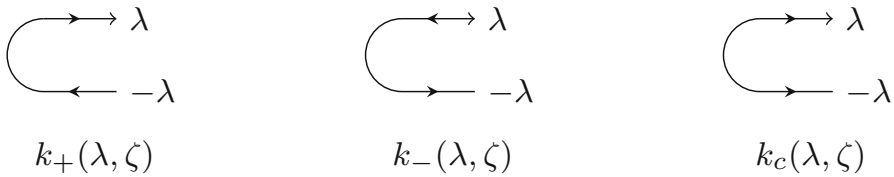


Fig. 3 The possible boundary configurations and boundary weights for the triangular reflecting end. The weights depend on the spectral parameter λ as well as on a boundary parameter ζ

to each vertex and each turn as in Figs. 2 and 3. Here φ is a fixed number. We call the turns k_{\pm} a ‘positive’ and ‘negative’ turn, respectively, and k_c can be seen as a turn with creation of arrows. These choices of weights satisfy the Yang–Baxter equation and the reflection equation with a triangular K -matrix where $k_c(\lambda, \zeta)$ is an off-diagonal element (see Sect. 2.1). Observe that k_c does not depend on ζ , and for $\varphi = 0$, we have diagonal reflection. Sometimes, we will refer to a ‘ w vertex’, where w is one of a_{\pm} , b_{\pm} or c_{\pm} , meaning a vertex with spin configurations corresponding to weight $w(\lambda)$, for some λ . Similarly, a ‘ k_{\pm} turn’ or ‘ k_c turn’ will refer to a turn with weight $k_{\pm}(\lambda, \zeta)$ or $k_c(\lambda, \zeta)$, respectively. The terms ‘positive’ or ‘negative turn’ will sometimes also be used to refer to a k_{\pm} turn, respectively.

The local weight at a vertex with the positive directions up and to the right depends on the spins of the surrounding edges, as well as on the difference between the spectral parameters on the incoming lines from the left and the bottom. Because of the reflecting ends, we need to differentiate between the vertices on the left-oriented and the right-oriented horizontal lines. The vertices in the right-oriented rows are depicted in Fig. 2, and the vertices in the left-oriented rows are the same, tilted 90 degrees counterclockwise, as in Fig. 4b. The (local) weight of the vertex in Fig. 4a is $w(\lambda_i - \mu_j)$, and for the vertex in Fig. 4b, the weight is $w(\mu_j - (-\lambda_i)) = w(\lambda_i + \mu_j)$, where w is one of a_{\pm} , b_{\pm} or c_{\pm} . The (local) boundary weight at each turn depends on the spin on the turning edge, but also on the spectral parameter λ_i of the line going through the turn, and on the fixed boundary parameter ζ , as in Fig. 3. The weight of a state is the product of all local weights of the vertices and the turns.

On the three sides without reflecting end, we impose the domain wall boundary conditions, with outgoing spin arrows to the right and ingoing arrows on the upper and lower boundaries. The ice rule implies that $n \geq m$ and that there are $n - m$ turns of type k_c .

The model described above is the six-vertex (6V) model of size $2n \times m$ with DWBC and one partially reflecting end. We want to find a determinant formula for the partition function

$$Z_{n,m}(\lambda, \mu) = \sum_{\text{state}} \text{weight}(\text{state}) \tag{3}$$

of this model, generalizing Tsuchiya’s [9] partition function for $m = n$. This also generalizes the results of Foda and Zarembo [12] from the rational to the trigonometric case.

Fig. 4 The different vertex weights depending on the direction of the row in the 6V model with reflecting end, with spectral parameters λ_i and μ_j



2.1 The Yang–Baxter equation and the reflection equation

Define V as a two-dimensional complex vector space. To each line of the lattice, we associate a copy of V . Given a parameter $\lambda \in \mathbb{C}$, define operators $R(\lambda) \in \text{End}(V \otimes V)$ by

$$R(\lambda) = \begin{pmatrix} a_+(\lambda) & 0 & 0 & 0 \\ 0 & b_+(\lambda) & c_-(\lambda) & 0 \\ 0 & c_+(\lambda) & b_-(\lambda) & 0 \\ 0 & 0 & 0 & a_-(\lambda) \end{pmatrix},$$

with the weights parametrized as in (1). The operator is called the R -matrix and satisfies the Yang–Baxter equation (YBE) on $V_1 \otimes V_2 \otimes V_3$ (where V_i are copies of V), i.e.

$$R_{12}(\lambda_1 - \lambda_2)R_{13}(\lambda_1 - \lambda_3)R_{23}(\lambda_2 - \lambda_3) = R_{23}(\lambda_2 - \lambda_3)R_{13}(\lambda_1 - \lambda_3)R_{12}(\lambda_1 - \lambda_2),$$

where the indices indicate on which spaces the R -matrix acts, e.g.

$$R_{12}(\lambda_1 - \lambda_2) = R(\lambda_1 - \lambda_2) \otimes \text{Id},$$

and similarly for R_{23} and R_{13} . The YBE is depicted in Fig. 5.

To describe the reflecting boundary, define a triangular operator $K(\lambda, \zeta) \in \text{End}(V)$, by

$$K(\lambda, \zeta) = \begin{pmatrix} k_+(\lambda, \zeta) & k_c(\lambda, \zeta) \\ 0 & k_-(\lambda, \zeta) \end{pmatrix},$$

with entries parametrized as in (2). The operator is called the K -matrix and satisfies the reflection equation for the above R -matrix on $V_0 \otimes V_{0'}$ [15], i.e.

$$\begin{aligned} R_{00'}(\lambda - \lambda')K_0(\lambda, \zeta)R_{0'0}(\lambda + \lambda')K_{0'}(\lambda', \zeta) \\ = K_{0'}(\lambda', \zeta)R_{0'0}(\lambda + \lambda')K_0(\lambda, \zeta)R_{00'}(\lambda - \lambda'), \end{aligned}$$

where $K_0(\lambda, \zeta) = K(\lambda, \zeta) \otimes \text{Id}$ and $K_{0'}(\lambda, \zeta) = \text{Id} \otimes K(\lambda, \zeta)$, see Fig. 6.

3 Izergin–Korepin method

Foda and Zarembo [12] used the Izergin–Korepin method [3, 6] to find the partition function (3) in the rational case, i.e. where $f(x) = x$, which in turn means that all

Fig. 5 The Yang–Baxter equation

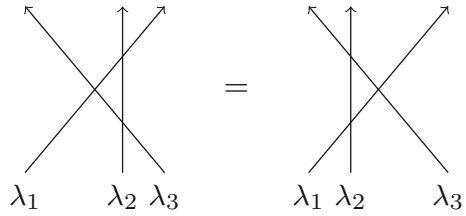
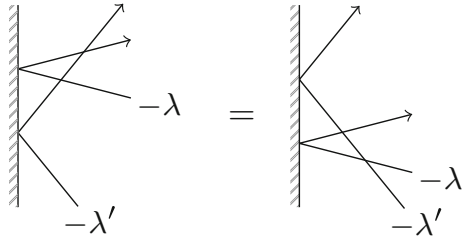


Fig. 6 The reflection equation



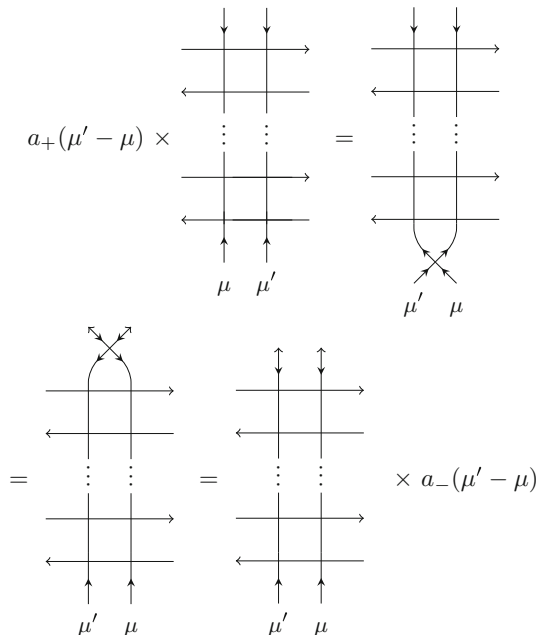
weights are rational. We follow the Izergin–Korepin procedure to find the determinant formula for the partition function in the trigonometric case.

Lemma 3.1 $Z_{n,m}(\lambda, \mu)$ is symmetric in λ_i and μ_j separately.

To prove this, we will use the so-called train argument.

Proof Consider two adjacent vertical lines with spectral parameters μ and μ' . Insert an extra vertex below the lattice, as in Fig. 7. Since we have DWBC, this will be a

Fig. 7 The partition function is symmetric in the μ_j 's. Because of the DWBC, an extra vertex can be moved through the lattice by using the YBE



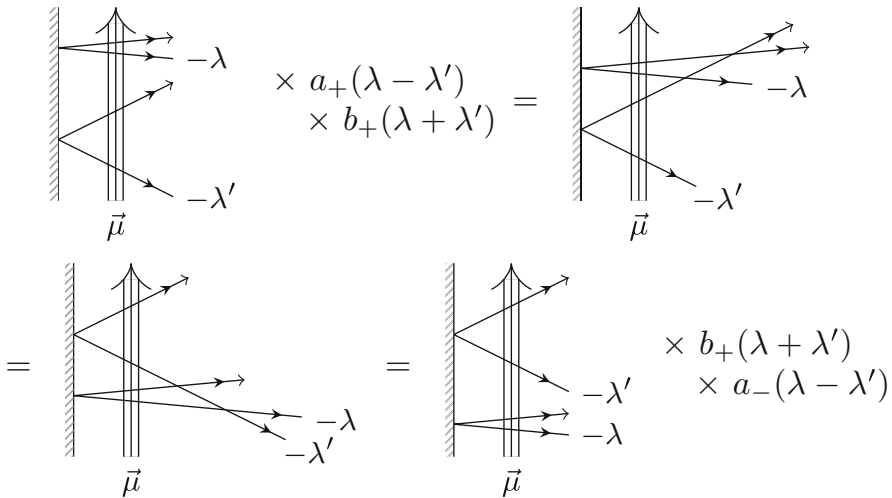


Fig. 8 The partition function is symmetric in the λ_i 's. Two adjacent double rows can switch places by adding extra vertices that can be moved through the lattice by using the YBE and the reflection equation. The triple arrows should be understood as n vertical arrows

vertex with weight $a_+(\mu - \mu')$. By the YBE, the extra vertex can be moved through the whole lattice to end up on the top, where it can be removed. On the top, the extra vertex has the weight $a_-(\mu - \mu')$. Since $a_+(\mu) = a_-(\mu) \neq 0$ for all μ , these factors cancel. Hence, this procedure switches μ and μ' .

To prove the symmetry in the λ_i 's, we add two extra vertices on the right, as in Fig. 8. The vertices can then be pulled through each other in a similar way as before, using the YBE and the reflection equation. Because of the boundary conditions and the ice rule, the extra vertices give rise to two extra factors on each side of the equation, see Fig. 8. Again since $a_+(\lambda) = a_-(\lambda) \neq 0$, the extra factors cancel. In this manner, we can pairwise switch spectral parameters on adjacent lines, which proves the lemma. \square

Lemma 3.2 For any $1 \leq j \leq m$, the function

$$e^{(2n-2)\mu_j} \prod_{i=1}^n \prod_{j=1}^m [f(\lambda_i + \mu_j + \gamma)f(\lambda_i - \mu_j + \gamma)] Z_{n,m}(\lambda, \mu)$$

is a polynomial of degree $2n - 1$ in $e^{2\mu_j}$.

Proof Multiply each vertex weight (1) by $f(\lambda + \gamma)$. This is the same as to multiply the partition function by

$$\prod_{i=1}^n \prod_{j=1}^m [f(\lambda_i + \mu_j + \gamma)f(\lambda_i - \mu_j + \gamma)].$$

Then, the new weights are

$$\hat{a}_\pm(\lambda) = f(\lambda + \gamma), \quad \hat{b}_\pm(\lambda) = e^{\mp\gamma} f(\lambda), \quad \hat{c}_\pm(\lambda) = e^{\pm\lambda} f(\gamma).$$

For a given j , all vertices involving μ_j lie along the same vertical line. Along a vertical line, the spin must be changed at least once from up spin on the bottom of the lattice to down spin on the top. Each vertical line thus contains at least one $\hat{c}_+(\lambda + \mu)$ or one $\hat{c}_-(\lambda - \mu)$ vertex, since these are the only combinations changing the spin in this way, see Fig. 2. These weights each contain a factor e^{μ_j} . Each additional \hat{c}_\pm vertex yields a factor $e^{\pm\mu_j}$. Each \hat{a}_\pm and \hat{b}_\pm vertex gives rise to a factor $e^{\mu_j+x} - e^{-\mu_j-x}$, where x possibly contains λ_i and γ . There can be a maximum of $2n - 1$ vertices of type \hat{a}_\pm and \hat{b}_\pm involving μ_j . We multiply $\prod_{i=1}^n \prod_{j=1}^m [f(\lambda_i + \mu_j + \gamma)f(\lambda_i - \mu_j + \gamma)] Z_{n,m}(\lambda, \mu)$ by $e^{(2n-2)\mu_j}$ to get a polynomial of degree $2n - 1$ in $e^{2\mu_j}$. \square

Lemma 3.3 $Z_{n,m}(\lambda, \mu)$ satisfies recurrence relations when setting $\mu_k = \pm\lambda_l$, namely,

$$Z_{n,m}(\lambda, \mu)|_{\mu_k=\pm\lambda_l} = e^{-n\gamma} e^{\zeta\pm\lambda_l} f(\zeta \mp \lambda_l) \prod_{i=1}^n \frac{f(\lambda_i + \lambda_l)}{f(\lambda_i + \lambda_l + \gamma)} Z_{n-1,m-1}(\hat{\lambda}_l, \hat{\mu}_k),$$

where

$$\hat{\lambda}_l = (\lambda_1, \dots, \hat{\lambda}_l, \dots, \lambda_n), \quad \hat{\mu}_k = (\mu_1, \dots, \hat{\mu}_k, \dots, \mu_m),$$

and $\hat{\lambda}_l$ indicates that the variable λ_l is omitted, and similarly $\hat{\mu}_k$ indicates that the variable μ_k is omitted.

Proof The vertices marked with a dot in Fig. 9 can only be b_+ or c_\pm vertices, because of the ice rule and the DWBC. Specialize $\mu_1 = -\lambda_1$ as in the left lattice. Now if the vertex at the bottom right is of type b_+ , it has zero weight and does not contribute to the partition function. Hence, the vertex must be a c_+ vertex. Then, the ice rule determines the spins on the two horizontal bottom lines and the rightmost vertical line, i.e. the vertices within the frozen region (the area marked with a dotted line) are determined uniquely. The part of the lattice outside the frozen region is then a lattice of size $2(n - 1) \times (m - 1)$ with DWBC. This yields a recursion relation for the partition function. Likewise, specializing $\mu_1 = \lambda_n$, as in the right lattice, forces the vertex marked with a dot to be a c_- vertex, and the ice rule determines the rest of the vertex weights within the frozen region. Lemma 3.1 yields that the recursion relations hold for $\mu_k = \pm\lambda_l$ for any $1 \leq l \leq n, 1 \leq k \leq m$. \square

Lemma 3.4 For $m = 0$, the partition function is

$$Z_{n,0}(\lambda) = \varphi^n \prod_{i=1}^n f(2\lambda_i).$$

Proof For $m = 0$, there are no vertical lines. In this case, the partition function is just a product of n turns of type k_c . \square

Because of the ice rule, $n \geq m$, so $n = 0$ is only possible if the system has no lines at all. It makes little sense to think of a system with no lines, but for the sake of the

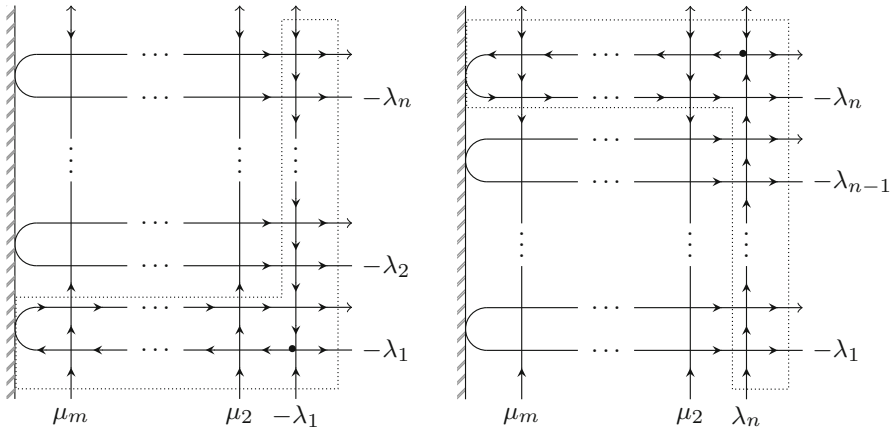


Fig. 9 Specializing $\mu_j = \pm\lambda_i$ determines the vertices inside the frozen region (the area marked with a dotted line). The part of the lattice outside the frozen region is then a lattice of size $2(n - 1) \times (m - 1)$ with DWBC. This yields a recursion relation for the partition function

following induction argument, we define $Z_{0,0} = 1$ as the empty product, which is in line with the above lemma.

Lemmas 3.2, 3.3 and 3.4 together determine $Z_{n,m}(\lambda, \mu)$. A polynomial of degree $2n - 1$ is uniquely determined by its values in $2n$ distinct points. Starting from the case $m = 0$, we can hence establish $Z_{n,m}(\lambda, \mu)$ as a function of μ_j by induction, using Lemma 3.3.

Theorem 3.5 *For the 6V model with DWBC and a partially reflecting end on a lattice of size $2n \times m$, $m \leq n$, the partition function is*

$$\begin{aligned}
 Z_{n,m}(\lambda, \mu) &= \varphi^{n-m} e^{(\binom{m}{2}-nm)\gamma} f(\gamma)^m \prod_{i=1}^m [e^{\mu_i+\zeta} f(\mu_i - \zeta)] \prod_{i=1}^n f(2\lambda_i) \\
 &\times \frac{\prod_{i=1}^n \prod_{j=1}^m [f(\mu_j + \lambda_i) f(\mu_j - \lambda_i)]}{\prod_{1 \leq i < j \leq m} [f(\mu_j + \mu_i) f(\mu_j - \mu_i)] \prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]} \\
 &\times \det_{1 \leq i, j \leq n} M, \tag{4}
 \end{aligned}$$

where M is an $n \times n$ matrix with

$$M_{ij} = \begin{cases} F(\lambda_i, \mu_j), & \text{if } i \leq m, \\ h((n - i)(2\lambda_j + \gamma)), & \text{if } m < i < n, \\ 1, & \text{if } i = n \text{ and } m \neq n, \end{cases}$$

where

$$F(\lambda_i, \mu_j) = \frac{1}{f(\mu_j + \lambda_i) f(\mu_j - \lambda_i) f(\mu_j + (\lambda_i + \gamma)) f(\mu_j - (\lambda_i + \gamma))},$$

$$f(x) = 2 \sinh(x) \text{ and } h(x) = 2 \cosh(x).$$

To prove this, it is enough to check that the right-hand side of (4) satisfies the properties stated in Lemmas 3.1 through 3.4.

Proof Let $\tilde{Z}_{n,m}(\lambda, \mu)$ be the right-hand side of (4). The symmetry in the λ_i 's and μ_j 's is obvious, so the condition in Lemma 3.1 holds for (4). Proving that the statements in Lemmas 3.3 and 3.4 hold for (4) is straightforward: put $\lambda_l = \mu_k$ into (4). Then all terms with $f(\mu_k - \lambda_l) = 0$ vanish. The only terms left are the terms coming from the part of the determinant containing $f(\mu_k - \lambda_l)$, since the corresponding factor from the numerator cancels. The terms left are

$$\begin{aligned} & \tilde{Z}_{n,m}(\lambda, \mu)|_{\mu_k=\lambda_l} \\ &= \frac{(-1)^{k+l} e^{-n\gamma} f(\gamma) e^{\mu_k+\zeta} f(\mu_k - \zeta) f(2\lambda_l)}{f(\mu_k + (\lambda_l + \gamma)) f(\mu_k - (\lambda_l + \gamma))} \\ & \times \prod_{\substack{1 \leq i \leq n \\ i \neq l}} [f(\mu_k + \lambda_i) f(\mu_k - \lambda_i)] \prod_{\substack{1 \leq j \leq m \\ j \neq k}} [f(\mu_j + \lambda_l) f(\mu_j - \lambda_l)] \\ & \times \frac{1}{\prod_{1 \leq i < l} [f(\lambda_i - \lambda_l) f(\lambda_i + \lambda_l + \gamma)] \prod_{l < i \leq n} [f(\lambda_l - \lambda_i) f(\lambda_l + \lambda_i + \gamma)]} \\ & \times \frac{1}{\prod_{1 \leq i < k} [f(\mu_k + \mu_i) f(\mu_k - \mu_i)] \prod_{k < i \leq m} [f(\mu_i + \mu_k) f(\mu_i - \mu_k)]} \\ & \times \tilde{Z}_{n-1,m-1}(\hat{\lambda}_l, \hat{\mu}_k). \end{aligned}$$

Simplifying this, we can conclude that the recurrence relation in Lemma 3.3 holds. The proof is similar for $\tilde{Z}_{n,m}(\lambda, \mu)|_{\mu_k=-\lambda_l}$.

For $m = 0$, the right-hand side of (4) is

$$\tilde{Z}_{n,0}(\lambda) = \frac{\varphi^n \prod_{i=1}^n f(2\lambda_i)}{\prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]} \det_{1 \leq i, j \leq n} M,$$

where M is an $n \times n$ matrix with

$$M_{ij} = \begin{cases} h((n - i)(2\lambda_j + \gamma)), & \text{for } 1 \leq i \leq n - 1, \\ 1, & \text{for } i = n. \end{cases}$$

Now let

$$D = \prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]$$

and put $y_i = h(2\lambda_i + \gamma)$. Then, we can write $D = \prod_{1 \leq i < j \leq n} (y_i - y_j)$. By row operations, we can rewrite $\det M$ as a determinant of a Vandermonde matrix $\tilde{M}_{ij} = y_j^{n-i}$. The determinant of a Vandermonde matrix is exactly D . Hence, the condition in Lemma 3.4 holds.

The only thing left is to prove the condition of Lemma 3.2. Observe that both the determinant and the denominator of (4) are 0 when $\mu_i = \pm\mu_j$. It is easy to see that $e^{2(m-1)\mu_i}$ times the denominator is a polynomial of order $2(m - 1)$ in $e^{2\mu_i}$ and is hence determined by its value in $2(m - 1)$ points (except for a constant factor). Since the determinant and the denominator share zeroes, the denominator is a factor in the determinant. This proves that $\tilde{Z}_{n,m}(\lambda, \mu)$ is a Laurent polynomial in e^{μ_i} . To determine the order, we check what happens with the determinant when $\mu_i \rightarrow \infty$. For this part of the proof, let $x_i = e^{-2\mu_i}$ and $y_j = h(2\lambda_j + \gamma)$. Let

$$\hat{F}(x_i, y_j) = \frac{1}{f(\mu_i + \lambda_j)f(\mu_i - \lambda_j)f(\mu_i + (\lambda_j + \gamma))f(\mu_i - (\lambda_j + \gamma))}.$$

Then, each entry of the determinant involving μ_i is

$$\hat{F}(x_i, y_j) = \frac{x_i^2}{(1 - e^{-\gamma}x_iy_j + e^{-2\gamma}x_i^2)(1 - e^{\gamma}x_iy_j + e^{2\gamma}x_i^2)}.$$

We can see that y_j always stands together with x_i , so in the Taylor series expansion, the power of x_i is always bigger than the power of y_j . The Taylor series expansion around $x_i = 0$ is

$$\hat{F}(x_i, y_j) = x_i^2 \left(\sum_{\substack{l \leq n-m-1 \\ l \leq k}} A_{k,l}x_i^k y_j^l + \sum_{\substack{n-m \leq l \\ l \leq k}} A_{k,l}x_i^k y_j^l \right),$$

for some constants $A_{k,l}$. The sum on the left can be written in terms of the last $n - m$ rows of the matrix. Hence in the determinant, the left sum can be removed by row operations, and we can change the entries $\hat{F}(x_i, y_j)$ in the determinant to

$$x_i^2 \sum_{l=n-m}^k A_{k,l}x_i^k y_j^l.$$

Letting $x_i \rightarrow 0$ (i.e. letting $\mu_i \rightarrow \infty$), $A_{n-m,n-m}x_i^{n-m+2}y_j^{n-m}$ is the leading term of the determinant, so the degree of the determinant in the variable $e^{2\mu_i}$ is $m - n - 2$. The degree of $e^{(2n-2)\mu_i} \prod_{i=1}^n \prod_{j=1}^m [f(\lambda_i + \mu_j + \gamma)f(\lambda_i - \mu_j + \gamma)] \tilde{Z}_{n,m}(\lambda, \mu)$ in $e^{2\mu_i}$ is

$$3n - m + 1 + \deg_{e^{2\mu_i}}(\det M) = 2n - 1$$

which proves the condition in Lemma 3.2. □

For $m = n$, the determinant formula in Theorem 3.5 coincides with Tsuchiya’s determinant formula.

In [12], another way to find the partition function is suggested, by starting from Tsuchiya’s determinant for $2n \times n$ lattices and step by step removing the extra vertical

lines by letting the corresponding variables $\mu_j \rightarrow \infty$, following Foda and Wheeler [11]. We prove Theorem 3.5 with this method in ‘‘Appendix A’’.

4 Kuperberg’s specialization of the parameters

In this section, we specialize the parameters in the definition of the partition function (3). By doing this, we find a way to count the total number of states of the model in terms of the partition function.

First, we will state and prove the $2n \times m$ generalization of Lemma 3.1 in [16]. For any given state, let $v(w)$ denote the number of vertices of type w .

Lemma 4.1 *For any given state of the 6V model with DWBC and one partially reflecting end on a $2n \times m$ lattice, we have*

$$v(b_+) - v(b_-) = \binom{n + 1}{2} - \binom{n - m + 1}{2} \quad \text{and} \quad v(c_+) - v(c_-) = m - 2v(k_-).$$

Proof We follow the same reasoning as in [16]. For the first result, we count the number of different spin arrows. To be able to keep track of the arrow flow through the system, in particular where the arrows go up and down, we think of a k_+ turn as one left, one up and one right arrow, and likewise we think of a k_- turn as one right, one down and one left arrow. To get the number of up arrows, we count the number of the different vertices and turns with up arrows, and similarly for the down arrows. To not count the edges twice, we count the vertices on every second row and add the edges from the boundaries. Hence, we get two different equations, depending on which rows we choose to count. Denote by w^N a vertex of type w on the upper part of a double line and similarly w^S for a vertex of type w on the lower part (N for north, S for south). We get the equations

$$v(\uparrow) = 2v(a_+^S) + 2v(b_+^S) + v(c_+^S) + v(c_-^S) + v(k_+), \tag{5}$$

$$v(\uparrow) = 2v(a_+^N) + 2v(b_-^N) + v(c_+^N) + v(c_-^N) + v(k_+) + m, \tag{6}$$

$$v(\downarrow) = 2v(a_-^S) + 2v(b_-^S) + v(c_+^S) + v(c_-^S) + v(k_-) + m, \tag{7}$$

and

$$v(\downarrow) = 2v(a_-^N) + 2v(b_+^N) + v(c_+^N) + v(c_-^N) + v(k_-). \tag{8}$$

To get the number of left arrows, we count the number of the different vertices with left arrows, plus the number of all k_+ and k_- turns, then divide everything by two. The number of left arrows is

$$v(\leftarrow) = v(a_-^N) + v(b_-^N) + v(a_+^S) + v(b_-^S) + \frac{v(c_+) + v(c_-) + v(k_+) + v(k_-)}{2}. \tag{9}$$

Similarly, the number of right arrows is

$$\begin{aligned}
 v(\rightarrow) &= v(a_+^N) + v(b_+^N) + v(a_-^S) + v(b_+^S) + \frac{v(c_+) + v(c_-) + v(k_+) + v(k_-)}{2} \\
 &\quad + v(k_c) + n.
 \end{aligned}
 \tag{10}$$

Adding Eqs. (5), (8) and two times (10), and subtracting (7), (6) and two times (9) yields

$$4[v(b_+) - v(b_-)] = 2v(\rightarrow) - 2v(\leftarrow) - 2n + 2m - 2v(k_c).
 \tag{11}$$

With the given boundary conditions, we have $v(k_c) = n - m$, and furthermore, we must have $v(\rightarrow) = 2n(m + 1) - \frac{m(m+1)}{2}$ and $v(\leftarrow) = \frac{m(m+1)}{2}$ to satisfy the ice rule. Insert this into (11). The first part of the lemma follows.

For the second part of the lemma, we observe that at a k_+ double row, the upper part of the double line starts and ends with a right arrow, and the lower part of the line starts with a left arrow and ends with a right arrow. The only vertex types that can change the arrows are the c_{\pm} vertices. Hence at each k_+ double row,

$$v(c_+^S) = v(c_-^S) + 1 \quad \text{and} \quad v(c_+^N) = v(c_-^N),$$

and similarly at each k_- double row,

$$v(c_+^S) = v(c_-^S) \quad \text{and} \quad v(c_+^N) = v(c_-^N) - 1.$$

At a k_c double row, each row starts and ends with right arrows, so

$$v(c_+^S) = v(c_-^S) \quad \text{and} \quad v(c_+^N) = v(c_-^N).$$

Since $v(c_+) = v(c_-)$ on the k_c rows, we only need to consider the k_+ and k_- rows. There are m rows of this type. Hence, we get the sum $v(c_+) = v(c_-) + m - 2v(k_-)$, which is the second result of the lemma. \square

Furthermore, we can count the number of a_{\pm} vertices.

Corollary 4.2 *For any given state of the 6V model with DWBC and one partially reflecting end on a $2n \times m$ lattice, the number of vertices of type a (i.e. a_+ and a_- together) is*

$$v(a) = mn + \binom{m}{2} - m + 2[v(k_-) - v(b_-) - v(c_-)].$$

Proof We have that $v(a) = 2mn - v(b) - v(c)$. Use Lemma 4.1 to compute $v(b)$ in terms of $v(b_-)$, and $v(c)$ in terms of $v(c_-)$. \square

Now let $\gamma = 4\pi i/3$ so that e^γ becomes a third root of 1. Then, let $\lambda_i \rightarrow \gamma$ and $\mu_j \rightarrow 0$ in the partition function $Z_{n,m}(\lambda, \mu)$. In this limit, the weights (1) and (2) are

$$\begin{aligned}
 a_{\pm}(\gamma) &= 1, & b_+(\gamma) &= c_-(\gamma) = -e^{-\gamma}, & b_-(\gamma) &= c_+(\gamma) = -e^\gamma f(\gamma), \\
 k_+(\gamma) &= e^{2\zeta} - e^{-2\gamma}, & k_-(\gamma) &= e^{2\zeta} - e^{2\gamma}, & k_c(\gamma) &= \varphi f(2\gamma),
 \end{aligned}$$

and the partition function becomes

$$Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) = \sum_{\text{states}} a_+(\gamma)^{v(a_+)} a_-(\gamma)^{v(a_-)} b_+(\gamma)^{v(b_+)} b_-(\gamma)^{v(b_-)} \\ \times c_+(\gamma)^{v(c_+)} c_-(\gamma)^{v(c_-)} k_+(\gamma)^{v(k_+)} k_-(\gamma)^{v(k_-)} k_c(\gamma)^{v(k_c)}.$$

Here we have used that for $\gamma = 4\pi i/3$, it holds that $f(2\gamma) = -f(\gamma)$. Inserting the expressions for the weights and using Lemma 4.1 and Corollary 4.2, the partition function can be computed to be:

$$Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) = (-1)^{\binom{m}{2} - nm + n} \varphi^{n-m} e^{((\binom{m+1}{2} - nm)\gamma)} f(\gamma)^{n-m} \\ \times \sum_{k=0}^m N_k (e^{2\zeta} - e^{-2\gamma})^k \left(\frac{e^{2\zeta} - e^{2\gamma}}{e^{2\gamma}} \right)^{m-k},$$

where N_k is the number of states with $v(k_+) = k$, i.e. where k is the number of k_+ turns. Let $Z_{n,m,k}(\lambda, \mu)$ be the partition function for the same model where the number of k_+ turns is fixed to $v(k_+) = k$. Then,

$$N_k = (-1)^{\binom{m}{2} - nm + n} \varphi^{m-n} e^{(nm - \binom{m+1}{2})\gamma} f(\gamma)^{m-n} \\ \times \left(\frac{1}{e^{2\zeta} - e^{-2\gamma}} \right)^k \left(\frac{e^{2\gamma}}{e^{2\zeta} - e^{2\gamma}} \right)^{m-k} Z_{n,m,k}(\gamma, \dots, \gamma, 0, \dots, 0). \tag{12}$$

Finally,

$$A(m, n) := \sum_{k=0}^m N_k \tag{13}$$

counts the total number of states of the model.

4.1 Connection to alternating sign matrices

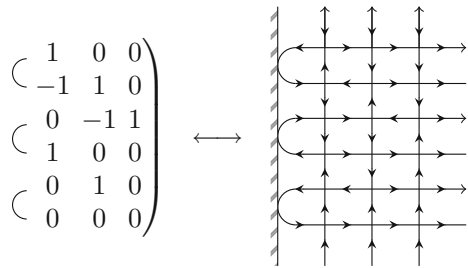
By specializing the parameters in the partition function for the 6V model with DWBC (without reflecting end), Kuperberg [8] was able to count the number of alternating sign matrices (ASMs). ASMs are matrices where the entries consist of 1, -1 and 0, such that the nonzero elements in each row and each column alternate in sign and such that the sum of the elements of any row or any column is 1. Hence, the first and last nonzero element of each row and each column is always 1.

In the case of DWBC and partially reflecting end, we also get a bijection to a type of ASM-like objects. Consider matrices of size $2n \times m$, $m \leq n$, consisting of elements 0, -1 and 1. Vertically and horizontally the nonzero elements alternate in sign. The sum of the elements of each column is 1, as for ASMs. Horizontally connect the rows pairwise on the left edge to form a double row as in Fig. 10. A double row may consist of only zeroes. If a row has any nonzero elements, the rightmost of these must be 1. Furthermore, the sum of the entries in a double row must be 0 or 1. Equivalently,

$$\begin{pmatrix} 1 & 0 \\ 0 & 0 \\ -1 & 1 \\ 0 & 0 \\ 1 & 0 \\ 0 & 0 \end{pmatrix} \quad
 \begin{pmatrix} 1 & 0 \\ 0 & 0 \\ 0 & 0 \\ -1 & 1 \\ 1 & 0 \\ 0 & 0 \end{pmatrix} \quad
 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & -1 & 1 \\ -1 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad
 \begin{pmatrix} 0 & 0 \\ 0 & 0 \\ 0 & 1 \\ 1 & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix}$$

Fig. 10 The three matrices on the left are counted by $A(m, n)$, whereas the rightmost matrix does not correspond to a state in the model considered in this paper

Fig. 11 The bijection between a UASM and a state of the 6V model with DWBC and a reflecting end



a double row can not consist of two rows both having 1 as their leftmost nonzero element.

Proposition 4.3 *The expression $A(m, n)$ yields the number of matrices of size $2n \times m$, $m \leq n$, described above.*

Proof There is a bijection between the states of the 6V model and matrices consisting of 0, 1 and -1 [8] (see Fig. 11), where a_{\pm} and b_{\pm} correspond to 0 in the matrix, and in the case of reflecting end, a c_{-} vertex in the upper part of a double row corresponds to 1, and c_{+} corresponds to -1 . On the lower part of a double row, c_{+} corresponds to 1, and c_{-} corresponds to -1 . Due to the ice rule and the boundary conditions, the rest of the vertices are determined uniquely.

The fact that there are more rows than columns allows for rows without c_{\pm} vertices, corresponding to rows with only zeroes in the matrix. If there are any c_{\pm} vertices on the lower part of a double row, k_{-} and k_c turns on that double row force the leftmost c_{\pm} vertex to be c_{-} , corresponding to -1 in the matrix, whereas a k_{+} turn forces it to be c_{+} , corresponding to 1 in the matrix. Similarly, if there are any c_{\pm} vertices on the upper part of the double row, k_{+} and k_c force the leftmost nonzero element in the matrix to be -1 , and k_{-} yields that the leftmost nonzero matrix element is 1. In this way, we can see that both rows cannot have 1 as their leftmost nonzero element. In a similar way, the DWBC impose that if a row has any nonzero element, then the rightmost of them must be 1.

□

From the discussion above, it follows that N_k (12) counts the number of matrices equivalent to states of the 6V model with exactly k positive turns. In the case $m = n$, the sum of the elements in each double row must be 1, and $A(m, n)$ counts the number of UASMs [10]. The case where $m = n$ and $k = 0$ corresponds to VSASMs.

5 Specialization in the determinant

In this section, we specialize the parameters in the determinant formula in the same way as in the previous section. We show that with this specialization of the variables, the partition function can be written as a determinant of Wilson polynomials, and finally we use this together with the results from the previous section to give a formula for the number of states of the model.

5.1 Moments and orthogonal polynomials

We start by rewriting the determinant formula of the partition function (4). Let $x_i = \lambda_i - \gamma$ and $y_j = \mu_j$. In the previous section, we let $\lambda_i \rightarrow \gamma$ and $\mu_i \rightarrow 0$, which corresponds to letting $x_i \rightarrow 0$ and $y_i \rightarrow 0$. For $1 \leq i \leq n$, multiply column i of the determinant by $f(2x_i)$ and for $1 \leq j \leq m$, multiply row j of the determinant by $f(2y_j)$. Then specify $\gamma = 4\pi i/3$. By row operations on the lower part of the determinant, we can then rewrite the partition function as

$$\begin{aligned}
 & Z_{n,m}(x_1 + \gamma, \dots, x_n + \gamma, y_1, \dots, y_m) \\
 &= \varphi^{n-m} e^{(\binom{m}{2} - nm)\gamma} f(\gamma)^m \frac{\prod_{i=1}^m [e^{y_i + \zeta} f(y_i - \zeta)] \prod_{i=1}^n f(2x_i + 2\gamma)}{\prod_{i=1}^n f(2x_i) \prod_{j=1}^m f(2y_j)} \\
 &\quad \times \frac{\prod_{i=1}^n \prod_{j=1}^m [f(y_j + (x_i + \gamma))f(y_j - (x_i + \gamma))]}{\prod_{1 \leq i < j \leq m} [f(y_j + y_i)f(y_j - y_i)] \prod_{1 \leq i < j \leq n} [f(x_i + x_j)f(x_i - x_j)]} \\
 &\quad \times \det_{1 \leq i, j \leq n} \tilde{M}, \tag{14}
 \end{aligned}$$

where

$$\tilde{M}_{ij} = \begin{cases} g(x_j, y_i), & \text{for } 1 \leq i \leq m, \\ f(2(n - i + 1)x_j), & \text{for } m + 1 \leq i \leq n, \end{cases}$$

with

$$g(x, y) = \frac{f(2x)f(2y)}{f(y + (x + \gamma))f(y - (x + \gamma))f(y + (x + 2\gamma))f(y - (x + 2\gamma))}.$$

It is apparent that $f(x) = \sinh x$ and $g(x, y)$ are odd functions. Furthermore, $x = 0$ is a zero of $f(x)$, and $x = 0$ and $y = 0$ are zeroes of $g(x, y)$. Hence, we can write the functions as $f(x) = x \hat{f}(x^2)$ and $g(x, y) = xy \hat{g}(x^2, y^2)$, where \hat{f} and \hat{g} are analytic

at $x = y = 0$. The partition function (14) can thus be written:

$$\begin{aligned} & Z_{n,m}(x_1 + \gamma, \dots, x_n + \gamma, y_1, \dots, y_m) \\ &= 2^{n-m} (n - m)! \varphi^{n-m} e^{\binom{m}{2} - nm} \gamma f(\gamma)^m \\ &\quad \times \frac{\prod_{i=1}^m [e^{y_i + \zeta} f(y_i - \zeta)] \prod_{i=1}^n f(2x_i + 2\gamma)}{\prod_{i=1}^n f(2x_i) \prod_{j=1}^m f(2y_j)} \\ &\quad \times \frac{\prod_{i=1}^n \prod_{j=1}^m [f(y_j + (x_i + \gamma)) f(y_j - (x_i + \gamma))] \prod_{j=1}^n x_j \prod_{k=1}^m y_k}{\prod_{1 \leq i < j \leq m} [f(y_j + y_i) f(y_j - y_i)] \prod_{1 \leq i < j \leq n} [f(x_i + x_j) f(x_i - x_j)]} \\ &\quad \times \det_{1 \leq i, j \leq n} \hat{M}, \end{aligned}$$

where

$$\hat{M}_{ij} = \begin{cases} \hat{g}(x_j^2, y_i^2), & \text{for } 1 \leq i \leq m, \\ \hat{f}((2(n - i + 1)x_j)^2), & \text{for } m + 1 \leq i \leq n. \end{cases}$$

Now regard x_i^2 and y_j^2 as our variables. Then in the limit where all $x_i, y_j \rightarrow 0$, the determinant can be written in terms of partial derivatives. Subtract the first column from the second and divide by $x_2^2 - x_1^2$; then in the limit $x_i \rightarrow 0$, this equals the first partial derivative. Similarly for the k th column, subtract the $(k - 1)$ th-order Taylor series expansion of each entry around $x_k^2 = 0$, which in the limit is a sum of the first $k - 1$ columns. Then divide by $\prod_{i=1}^{k-1} (x_k^2 - x_i^2)/k!$. By l'Hôpital's rule, this limit equals the $(k - 1)$ th partial derivative. A similar argument for the first m rows lets us write the entries as partial derivatives in y_i^2 as well. The method is described in detail in [7]. Hence in the limit, the partition function equals

$$\begin{aligned} & Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) \\ &= \lim_{x_i, y_j \rightarrow 0} 2^{n-m} (n - m)! \varphi^{n-m} e^{\binom{m}{2} - nm} \gamma f(\gamma)^m \\ &\quad \times \frac{\prod_{i=1}^m [e^{y_i + \zeta} f(y_i - \zeta)] \prod_{i=1}^n f(2x_i + 2\gamma)}{\prod_{i=1}^n f(2x_i) \prod_{j=1}^m f(2y_j)} \\ &\quad \times \frac{\prod_{i=1}^n \prod_{j=1}^m [f(y_j + (x_i + \gamma)) f(y_j - (x_i + \gamma))] \prod_{j=1}^n x_j \prod_{i=1}^m y_i}{\prod_{1 \leq i < j \leq m} [f(y_j + y_i) f(y_j - y_i)] \prod_{1 \leq i < j \leq n} [f(x_i + x_j) f(x_i - x_j)]} \\ &\quad \times \frac{\prod_{1 \leq i < j \leq n} (x_j^2 - x_i^2) \prod_{1 \leq i < j \leq m} (y_j^2 - y_i^2)}{\prod_{i=1}^n (i - 1)! \prod_{j=1}^m (j - 1)!} \end{aligned}$$

$$\times \det \begin{pmatrix} \hat{g}(x_1^2, y_1^2) & \frac{\partial}{\partial x_2^2} \hat{g}(x_2^2, y_1^2) & \cdots & \frac{\partial^{n-1}}{\partial x_n^{2(n-1)}} \hat{g}(x_n^2, y_1^2) \\ \frac{\partial}{\partial y_2^2} \hat{g}(x_1^2, y_2^2) & \frac{\partial^2}{\partial x_2^2 \partial y_2^2} \hat{g}(x_2^2, y_2^2) & \cdots & \frac{\partial^n}{\partial x_n^{2(n-1)} \partial y_2^2} \hat{g}(x_n^2, y_2^2) \\ \vdots & \vdots & \ddots & \vdots \\ \frac{\partial^{m-1}}{\partial y_m^{2(m-1)}} \hat{g}(x_1^2, y_m^2) & \frac{\partial^m}{\partial x_2^2 \partial y_m^{2(m-1)}} \hat{g}(x_2^2, y_m^2) & \cdots & \frac{\partial^{n+m-2}}{\partial x_n^{2(n-1)} \partial y_m^{2(m-1)}} \hat{g}(x_n^2, y_m^2) \\ \hat{f}((2(n-m)x_1)^2) & \frac{\partial}{\partial x_2^2} \hat{f}((2(n-m)x_2)^2) & \cdots & \frac{\partial^{n-1}}{\partial x_n^{2(n-1)}} \hat{f}((2(n-m)x_n)^2) \\ \vdots & \vdots & \ddots & \vdots \\ \hat{f}((2x_1)^2) & \frac{\partial}{\partial x_2^2} \hat{f}((2x_2)^2) & \cdots & \frac{\partial^{n-1}}{\partial x_n^{2(n-1)}} \hat{f}((2x_n)^2) \end{pmatrix}.$$

To cancel the zeroes in the denominators, the prefactors can be simplified by

$$\begin{aligned} & \lim_{x_i, y_j \rightarrow 0} \left(\frac{\prod_{i=1}^n x_i \prod_{j=1}^m y_j}{\prod_{i=1}^n f(2x_i) \prod_{j=1}^m f(2y_j)} \right. \\ & \quad \times \left. \frac{\prod_{1 \leq i < j \leq n} (x_j^2 - x_i^2) \prod_{1 \leq i < j \leq m} (y_j^2 - y_i^2)}{\prod_{1 \leq i < j \leq n} [f(x_i + x_j) f(x_i - x_j)] \prod_{1 \leq i < j \leq m} [f(y_j + y_i) f(y_j - y_i)]} \right) \\ & = \frac{(-1)^{\binom{n}{2}}}{2^{n^2+n+m^2+m}}. \end{aligned}$$

Moreover,

$$f(2lx) = 4lx \left(1 + \frac{(2lx)^2}{3!} + \frac{(2lx)^4}{5!} + \cdots \right),$$

so

$$\hat{f}(x) = 2 \left(1 + \frac{x}{3!} + \frac{x^2}{5!} + \cdots \right),$$

and

$$\frac{\partial^k}{\partial x^{2k}} \hat{f}((2lx)^2)|_{x=0} = \frac{2(2l)^{2k} k!}{(2k + 1)!}.$$

For $\gamma = 4\pi i/3$,

$$g(x, y) = \frac{f(x - y)}{f(3x - 3y)} - \frac{f(x + y)}{f(3x + 3y)}.$$

Following [14], we use Fourier transforms to see that

$$\frac{f(x)}{f(3x)} = \frac{1}{2\sqrt{3}} \int_{-\infty}^{\infty} e^{-ixt} \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt.$$

From this, it follows that

$$\frac{f(x - y)}{f(3x - 3y)} - \frac{f(x + y)}{f(3x + 3y)} = \frac{2}{\sqrt{3}} \int_0^{\infty} \sin(xt) \sin(yt) \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt.$$

By Taylor expansion and Lebesgue’s dominated convergence theorem, we get

$$g(x, y) = xy \frac{2}{\sqrt{3}} \sum_{k,l=0}^{\infty} \left(\frac{(-1)^{k+l} x^{2k} y^{2l}}{(2k + 1)!(2l + 1)!} \int_0^{\infty} t^{2k+2l+2} \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt \right),$$

for $|x|, |y| < \pi/3$, so

$$\hat{g}(x^2, y^2) = \frac{2}{\sqrt{3}} \sum_{k,l=0}^{\infty} \left(\frac{(-1)^{k+l} x^{2k} y^{2l}}{(2k + 1)!(2l + 1)!} \int_0^{\infty} t^{2k+2l+2} \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt \right).$$

Hence,

$$\begin{aligned} & \left. \frac{\partial^{k+l}}{\partial x^{2k} \partial y^{2l}} \hat{g}(x^2, y^2) \right|_{x=y=0} \\ &= \frac{2(-1)^{k+l} k! l!}{\sqrt{3}(2k + 1)!(2l + 1)!} \int_0^{\infty} t^{2k+2l+2} \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt. \end{aligned}$$

Let c_k denote the moment

$$c_k = \int t^k d\mu(t),$$

where

$$\int f(t) d\mu(t) = \int_0^{\infty} f(t^2) t^2 \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt.$$

The corresponding inner product is

$$\langle p(t), q(t) \rangle = \int_0^{\infty} p(t^2) q(t^2) t^2 \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} dt,$$

for polynomials p and q .

In the limit $x_i, y_j \rightarrow 0$, the partition function thus equals

$$\begin{aligned} Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) &= (-1)^{\binom{n}{2}+mn} \prod_{i=1}^n \frac{1}{(2i - 1)!} \prod_{j=1}^m \frac{1}{(2j - 1)!} \\ &\times \frac{2^{n-2m-m^2-n^2}}{3^{m/2}} (n - m)! \varphi^{n-m} e^{\binom{m}{2}-mn} \gamma f(\gamma)^{m+2mn} f(2\gamma)^n (1 - e^{2\zeta})^m \\ &\times \det \begin{pmatrix} c_0 & -c_1 & \dots & (-1)^{n+1} c_{n-1} \\ -c_1 & c_2 & \dots & (-1)^{n+2} c_n \\ \vdots & \vdots & \ddots & \vdots \\ (-1)^{m+1} c_{m-1} & (-1)^{m+2} c_m & \dots & (-1)^{n+m} c_{n+m-2} \\ 1 & (2(n - m))^2 & \dots & (2(n - m))^{2(n-1)} \\ \vdots & \vdots & \ddots & \vdots \\ 1 & 2^2 & \dots & 2^{2(n-1)} \end{pmatrix}. \end{aligned}$$

We can write the entries c_k in terms of inner products. Since

$$(-1)^{\frac{n(n-1)}{2} + \lfloor \frac{n}{2} \rfloor} = 1$$

for all n , the partition function above becomes

$$\begin{aligned}
 Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) &= (-1)^{mn-m(m-1)/2} \prod_{i=1}^n \frac{1}{(2i-1)!} \prod_{j=1}^m \frac{1}{(2j-1)!} \\
 &\times \frac{2^{n-2m-m^2-n^2}}{3^{m/2}} (n-m)! \varphi^{n-m} e^{\binom{m}{2}-nm} \gamma f(\gamma)^{2mn+m} f(2\gamma)^n (1-e^{2\zeta})^m \\
 &\times \det \begin{pmatrix} \langle 1, 1 \rangle & \langle t, 1 \rangle & \cdots & \langle t^{n-1}, 1 \rangle \\ \langle 1, t \rangle & \langle t, t \rangle & \cdots & \langle t^{n-1}, t \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle 1, t^{m-1} \rangle & \langle t, t^{m-1} \rangle & \cdots & \langle t^{n-1}, t^{m-1} \rangle \\ 1 & -(2(n-m))^2 & \cdots & (-1)^{n-1} (2(n-m))^{2(n-1)} \\ \vdots & \vdots & \ddots & \vdots \\ 1 & -2^2 & \cdots & (-1)^{n-1} 2^{2(n-1)} \end{pmatrix}.
 \end{aligned}$$

Similar determinants show up, for example, in [14]. By row and column operations, the determinant can be rewritten, and the partition function is

$$\begin{aligned}
 Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) &= (-1)^{mn-m(m-1)/2} \prod_{j=1}^n \frac{1}{(2j-1)!} \prod_{j=1}^m \frac{1}{(2j-1)!} \\
 &\times \frac{2^{n-2m-m^2-n^2}}{3^{m/2}} (n-m)! \varphi^{n-m} e^{\binom{m}{2}-nm} \gamma f(\gamma)^{2mn+m} f(2\gamma)^n (1-e^{2\zeta})^m \\
 &\times \det \begin{pmatrix} \langle p_0(t), q_0(t) \rangle & \langle p_1(t), q_0(t) \rangle & \cdots & \langle p_{n-1}(t), q_0(t) \rangle \\ \langle p_0(t), q_1(t) \rangle & \langle p_1(t), q_1(t) \rangle & \cdots & \langle p_{n-1}(t), q_1(t) \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle p_0(t), q_{m-1}(t) \rangle & \langle p_1(t), q_{m-1}(t) \rangle & \cdots & \langle p_{n-1}(t), q_{m-1}(t) \rangle \\ p_0(-(2(n-m))^2) & p_1(-(2(n-m))^2) & \cdots & p_{n-1}(-(2(n-m))^2) \\ \vdots & \vdots & \ddots & \vdots \\ p_0(-4^2) & p_1(-4^2) & \cdots & p_{n-1}(-4^2) \\ p_0(-2^2) & p_1(-2^2) & \cdots & p_{n-1}(-2^2) \end{pmatrix},
 \end{aligned}$$

where $p_k(t)$ and $q_k(t)$ are arbitrary monic polynomials of degree k . We can choose $\{p_i = q_i\}_{i=1}^n$ orthogonal, i.e.

$$\langle p_k(t), p_l(t) \rangle = \int p_k(t)p_l(t)d\mu(t) = h_k\delta_{kl}, \tag{15}$$

for some h_k . Then, most entries in the upper part of the matrix become 0, and we can simplify the expression further. We also rearrange the rows. Then, the partition function is

$$\begin{aligned} & Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) \\ &= \frac{2^{n-2m-m^2-n^2}}{3^{m/2}} (n-m)! \varphi^{n-m} e^{\binom{m}{2}-nm}\gamma f(\gamma)^{2mn+m} f(2\gamma)^n (1-e^{2\xi})^m \\ &\times (-1)^{\binom{n}{2}+m} \prod_{j=1}^n \frac{1}{(2j-1)!} \prod_{j=1}^m \frac{1}{(2j-1)!} \prod_{i=0}^{m-1} \langle p_i(t), p_i(t) \rangle \\ &\times \det \begin{pmatrix} p_m(-2^2) & p_{m+1}(-2^2) & \cdots & p_{n-1}(-2^2) \\ p_m(-4^2) & p_{m+1}(-4^2) & \cdots & p_{n-1}(-4^2) \\ \vdots & \vdots & \ddots & \vdots \\ p_m(-(2(n-m))^2) & p_{m+1}(-(2(n-m))^2) & \cdots & p_{n-1}(-(2(n-m))^2) \end{pmatrix}. \end{aligned} \tag{16}$$

5.2 Wilson polynomials

Colomo and Pronko [14] found a way to write the determinant formula of the partition function of the 6V model with DWBC in terms of orthogonal polynomials. In this section, we show that we can rewrite the determinant formula (4) in terms of orthogonal polynomials as well.

We can rewrite the weight $\mu(t)$ in terms of the gamma function, defined as the analytic continuation of the integral

$$\Gamma(z) = \int_0^\infty x^{z-1} e^{-x} dx,$$

which is defined only for $\text{Re}(z) > 0$. It can be shown that the gamma function satisfies the following identities:

$$|\Gamma(1 + ix)|^2 = \frac{\pi x}{\sinh(\pi x)}$$

and

$$\prod_{k=0}^{l-1} \Gamma\left(x + \frac{k}{l}\right) = (2\pi)^{\frac{l-1}{2}} l^{\frac{1}{2}-lx} \Gamma(lx). \tag{17}$$

Hence, we have

$$\begin{aligned} \mu(t) &= t^2 \frac{\sinh(\pi t/6)}{\sinh(\pi t/2)} \\ &= \frac{3^2}{2^2 \pi^3} \left| \frac{\Gamma(it/6 + 1/3)\Gamma(it/6 + 1/2)\Gamma(it/6 + 2/3)\Gamma(it/6 + 1)}{\Gamma(it/3)} \right|^2. \end{aligned} \tag{18}$$

This is the orthogonality weight for a known family of polynomials, namely the Wilson polynomials. These are defined in terms of generalized hypergeometric series as (see, e.g. [17])

$$\frac{W_k(x^2; a, b, c, d)}{(a + b)_k(a + c)_k(a + d)_k} = {}_4F_3 \left(\begin{matrix} -k, k + a + b + c + d - 1, a + ix, a - ix \\ a + b, a + c, a + d \end{matrix} \middle| 1 \right),$$

where the orthogonality condition reads

$$\begin{aligned} &\frac{1}{2\pi} \int_0^\infty W_k(x^2; a, b, c, d) W_l(x^2; a, b, c, d) \\ &\quad \times \left| \frac{\Gamma(ix + a)\Gamma(ix + b)\Gamma(ix + c)\Gamma(ix + d)}{\Gamma(2ix)} \right|^2 dx \\ &= \frac{\Gamma(k + a + b)\Gamma(k + a + c)\Gamma(k + a + d)\Gamma(k + b + c)\Gamma(k + b + d)\Gamma(k + c + d)}{\Gamma(2k + a + b + c + d)} \\ &\quad \times (k + a + b + c + d - 1)_k k! \delta_{kl}, \end{aligned} \tag{19}$$

and where $(a)_k = a(a + 1) \cdots (a + k - 1)$, for $k \geq 1$, and $(a)_0 = 1$ is the rising factorial. Comparing this to (15) and (18), we choose the parameters $a = 1/3, b = 1/2, c = 2/3, d = 1$ and $x = t/6$. Then, the Wilson polynomials are:

$$\begin{aligned} &W_k \left(\left(\frac{t}{6} \right)^2; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \\ &= (5/6)_k (4/3)_k k! \sum_{j=0}^k \frac{(-k)_j (3/2 + k)_j (1/3 + it/6)_j (1/3 - it/6)_j}{(5/6)_j (4/3)_j (j!)^2}. \end{aligned} \tag{20}$$

These are polynomials in t^2 of degree k , with leading coefficient

$$\kappa_k = \frac{(-1)^k (3/2 + k)_k}{6^{2k}}. \tag{21}$$

The polynomials can hence be written

$$W_k \left(\left(\frac{t}{6} \right)^2; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) = \kappa_k p_k(t^2),$$

where $p_k(t)$ is a monic polynomial of degree k in t . The right-hand side of (19) with the above choices of parameters is

$$\frac{\Gamma(k + 5/6)\Gamma(k + 1)\Gamma(k + 4/3)\Gamma(k + 7/6)\Gamma(k + 3/2)\Gamma(k + 5/3)}{\Gamma(2k + 5/2)} \times (k + 3/2)_k k! \delta_{kl} = \frac{\pi^2 (k + 3/2)_k k! (6k + 4)! (2k + 1)!}{2^{2k-1} 3^{6k+9/2} (4k + 3)!} \delta_{kl}, \tag{22}$$

where we have used (17) and the fact that $\Gamma(j + 1) = j!$ if j is a nonnegative integer. We insert (19)–(22) into (15) and get

$$\langle p_k(t), p_l(t) \rangle = \frac{2^{2k+1} k! (6k + 4)! (2k + 1)!}{3^{2k+3/2} (3/2 + k)_k (4k + 3)!} \delta_{kl}.$$

Hence,

$$\prod_{j=0}^{m-1} \langle p_j(t), p_j(t) \rangle = \frac{2^{2m^2-m}}{3^{m^2+m/2}} \prod_{j=1}^m \frac{(6j - 2)! ((2j - 1)!)^2 (2j - 2)!}{(4j - 3)! (4j - 1)!},$$

where we used that

$$(x)_k = \frac{\Gamma(x + k)}{\Gamma(x)}. \tag{23}$$

Insert $p_k(t^2) = W_k \left(\left(\frac{t}{6}\right)^2; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) / \kappa_k$ into the partition function (16), which then reads

$$\begin{aligned} & Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) \\ &= \frac{2^{n-n^2-3m+m^2}}{3^{m^2+m}} (n - m)! \varphi^{n-m} e^{\binom{m}{2} - nm} \gamma f(\gamma)^{2mn+m} f(2\gamma)^n (1 - e^{2\zeta})^m \\ &\times (-1)^{\binom{n}{2} + m} \prod_{i=m+1}^n \frac{1}{(2i - 1)!} \prod_{j=1}^m \frac{(6j - 2)! (2j - 2)!}{(4j - 3)! (4j - 1)!} \\ &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) / \kappa_{m+j-1} \right). \end{aligned}$$

We factor out κ_{m+j-1} from each column of the determinant and use (23) and get

$$\begin{aligned}
 & Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) \\
 &= \frac{2^{n^2-n-m^2-m}}{3^{2m^2-n^2+n}} (n-m)! \varphi^{n-m} e^{\binom{m}{2}-nm}\gamma f(\gamma)^{2mn+m} f(2\gamma)^n (1-e^{2\zeta})^m \\
 &\times (-1)^{\binom{m+1}{2}} \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \prod_{j=m+1}^n \frac{1}{(j-1)!} \\
 &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right). \tag{24}
 \end{aligned}$$

5.3 A formula for the number of states

Next, we rewrite the partition function as a sum over the number of positive turns, k , to be able to identify the terms with the terms in the formulas from Sect. 4. We write

$$1 - e^{2\zeta} = -\frac{e^\gamma f(\gamma)}{f(2\gamma)} \left((e^{2\zeta} - e^{-2\gamma}) + \frac{e^{2\zeta} - e^{2\gamma}}{e^{2\gamma}} \right),$$

and by using the binomial theorem we get

$$(1 - e^{2\zeta})^m = \left(-\frac{e^\gamma f(\gamma)}{f(2\gamma)} \right)^m \sum_{k=0}^m \binom{m}{k} (e^{2\zeta} - e^{-2\gamma})^k \left(\frac{e^{2\zeta} - e^{2\gamma}}{e^{2\gamma}} \right)^{m-k}.$$

For $\gamma = 4\pi i/3$, we have $f(2\gamma) = -f(\gamma)$ and $f(\gamma)^2 = -3$. The partition function (24) thus becomes

$$\begin{aligned}
 & Z_{n,m}(\gamma, \dots, \gamma, 0, \dots, 0) = \sum_{k=0}^m \binom{m}{k} (e^{2\zeta} - e^{-2\gamma})^k \left(\frac{e^{2\zeta} - e^{2\gamma}}{e^{2\gamma}} \right)^{m-k} \\
 &\times (-1)^{\binom{m+1}{2}+mn+n} \frac{2^{n^2-n-m^2-m}}{3^{2m^2-n^2+n-mn}} (n-m)! \varphi^{n-m} e^{\binom{m+1}{2}-nm}\gamma f(\gamma)^{m+n} \\
 &\times \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \prod_{j=m+1}^n \frac{1}{(j-1)!} \\
 &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right).
 \end{aligned}$$

The terms $(e^{2\zeta} - e^{-2\gamma})^k \left(\frac{e^{2\zeta} - e^{2\gamma}}{e^{2\gamma}} \right)^{m-k}$ are linearly independent for different k 's as functions of ζ , since the k th term has a zero of degree $m - k$ in $\zeta = \gamma$. Therefore,

we can fix a k as in the previous section, and we get

$$\begin{aligned}
 Z_{n,m,k}(\gamma, \dots, \gamma, 0, \dots, 0) &= \binom{m}{k} (e^{2\xi} - e^{-2\gamma})^k \left(\frac{e^{2\xi} - e^{2\gamma}}{e^{2\gamma}} \right)^{m-k} \\
 &\times (-1)^{\binom{m+1}{2} + mn + n} \frac{2^{n^2 - n - m^2 - m}}{3^{2m^2 - n^2 + n - mn}} (n - m)! \varphi^{n-m} e^{\binom{m+1}{2} - nm} \gamma f(\gamma)^{m+n} \\
 &\times \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \prod_{j=m+1}^n \frac{1}{(j-1)!} \\
 &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right). \tag{25}
 \end{aligned}$$

Now we can go back to N_k (12), the number of states where k is the number of k_+ turns. We insert (25) and get the following theorem.

Theorem 5.1 *For the 6V model with DWBC and a partially reflecting end, the number of states with exactly k turns of type k_+ is*

$$\begin{aligned}
 N_k &= \binom{m}{k} \frac{2^{n^2 - n - m^2 - m} (n - m)!}{3^{2m^2 - m - n^2 + n - mn}} \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \prod_{j=m+1}^n \frac{1}{(j-1)!} \\
 &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right). \tag{26}
 \end{aligned}$$

As a corollary, we get that the total number of states (13) of the model is

$$\begin{aligned}
 A(m, n) &= \frac{2^{n^2 - n - m^2} (n - m)!}{3^{2m^2 - m - n^2 + n - mn}} \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \prod_{j=m+1}^n \frac{1}{(j-1)!} \\
 &\times \det_{1 \leq l, j \leq n-m} \left(W_{m+j-1} \left(-\frac{l^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right).
 \end{aligned}$$

From Proposition 4.3, it thus follows that the number of matrices described in Sect. 4.1 is also given by this expression.

5.4 An alternative expression for the number of states

In this section, we derive another way to present the expression (26). We can insert the formula (20) for $W_{m+j-1} \left(-\frac{k^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right)$ into the determinant

$$D = \det_{1 \leq k, j \leq n-m} \left(W_{m+j-1} \left(-\frac{k^2}{9}; \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, 1 \right) \right).$$

Factor out $(5/6)_{m+j-1}(4/3)_{m+j-1}(m+j-1)!$ from each column j . We can rewrite the sum in each determinant entry

$$\sum_{l=0}^{m+j-1} \frac{(1-m-j)_l(1/2+m+j)_l(1/3-k/3)_l(1/3+k/3)_l}{(5/6)_l(4/3)_l(l!)^2}$$

of row k and column j to go from 0 to $n-1$, since the terms disappear when $l > m+j-1$. We then get

$$D = \prod_{j=m}^{n-1} (5/6)_j(4/3)_j j! \times \det_{1 \leq k, j \leq n-m} \left(\sum_{l=0}^{n-1} \frac{(1-m-j)_l(1/2+m+j)_l(1/3+k/3)_l(1/3-k/3)_l}{(5/6)_l(4/3)_l(l!)^2} \right).$$

We use linearity of the rows and write the determinant as a sum of determinants. The above becomes

$$D = \prod_{j=m}^{n-1} (5/6)_j(4/3)_j j! \sum_{l_1, l_2, \dots, l_{n-m}=0}^{n-1} \left(\prod_{i=1}^{n-m} \frac{((1-i)/3)_{l_i}((1+i)/3)_{l_i}}{(5/6)_{l_i}(4/3)_{l_i}(l_i!)^2} \times \det_{1 \leq k, j \leq n-m} ((1-m-j)_{l_k}(1/2+m+j)_{l_k}) \right).$$

Since $(x)_j/(x)_i = (x+i)_{j-i}$, for $j \geq i$, we can write

$$(1-m-j)_{l_k}(1/2+m+j)_{l_k} = \frac{(1-n)_{l_k}(m+3/2)_{l_k}}{(1-n)_{n-m-j}(m+3/2)_{j-1}}(1-n+l_k)_{n-m-j}(m+3/2+l_k)_{j-1},$$

and we can factor out all factors that depend either only on row k or only on column j . The determinant D now becomes

$$D = \prod_{j=m}^{n-1} (5/6)_j(4/3)_j j! \times \sum_{l_1, l_2, \dots, l_{n-m}=0}^{n-1} \left(\prod_{i=1}^{n-m} \frac{((1-i)/3)_{l_i}((1+i)/3)_{l_i}(1-n)_{l_i}(m+3/2)_{l_i}}{(5/6)_{l_i}(4/3)_{l_i}(l_i!)^2(1-n)_{i-1}(m+3/2)_{i-1}} \times \det_{1 \leq k, j \leq n-m} ((1-n+l_k)_{n-m-j}(m+3/2+l_k)_{j-1}) \right).$$

Each element of the alternant matrix

$$\hat{D} = \det_{1 \leq k, j \leq n-m} ((1-n+l_k)_{n-m-j}(m+3/2+l_k)_{j-1})$$

is a polynomial in l_k of degree $n - m - 1$. This is a special case of a bigger family of determinants, see further [18, Lemma 3]. The determinant is

$$\hat{D} = C \prod_{1 \leq i < j \leq n-m} (l_i - l_j)$$

for some C not depending on the l_k 's. To find C , we put $l_k = -k - m - 1/2$ which makes the matrix triangular. Then,

$$\prod_{j=0}^{n-m-1} (-1)^j j! (-1/2 - n - m - j)_{n-m-1-j} = C \prod_{1 \leq i < j \leq n-m} (j - i),$$

so

$$C = \prod_{j=1}^{n-m} (5/2 + 2n - 2j)_{j-1}.$$

Hence, the determinant D becomes

$$\begin{aligned} D &= \prod_{j=m}^{n-1} (5/6)_j (4/3)_j j! \prod_{j=1}^{n-m} \frac{(5/2 + 2n - 2j)_{j-1}}{(1-n)_{j-1} (m + 3/2)_{j-1}} \\ &\times \sum_{l_1, l_2, \dots, l_{n-m}=0}^{n-1} \left(\prod_{i=1}^{n-m} \frac{((1-i)/3)_{l_i} ((1+i)/3)_{l_i} (1-n)_{l_i} (m + 3/2)_{l_i}}{(5/6)_{l_i} (4/3)_{l_i} (l_i!)^2} \right. \\ &\times \left. \prod_{1 \leq i < j \leq n-m} (l_i - l_j) \right). \end{aligned}$$

Series of similar type appear, for example, in [19, 20].

Inserting the above into (26) of the last section, we get an expression for the number of states with exactly k turns of type k_+ as an $(n - m)$ -fold hypergeometric sum,

$$\begin{aligned} N_k &= \binom{m}{k} \frac{2^{n^2-n-m^2-m} (n-m)!}{3^{2m^2-m-n^2+n-mn}} \prod_{j=1}^n \frac{(2j-2)!}{(4j-3)!} \prod_{j=1}^m \frac{(6j-2)!}{(4j-1)!} \\ &\times \prod_{j=m+1}^n (5/6)_{j-1} (4/3)_{j-1} \prod_{j=1}^{n-m} \frac{(5/2 + 2n - 2j)_{j-1}}{(1-n)_{j-1} (m + 3/2)_{j-1}} \\ &\times \sum_{l_1, l_2, \dots, l_{n-m}=0}^{n-1} \left(\prod_{i=1}^{n-m} \frac{((1-i)/3)_{l_i} ((1+i)/3)_{l_i} (1-n)_{l_i} (m + 3/2)_{l_i}}{(5/6)_{l_i} (4/3)_{l_i} (l_i!)^2} \right. \\ &\times \left. \prod_{1 \leq i < j \leq n-m} (l_i - l_j) \right). \end{aligned}$$

Acknowledgements I would like to thank my supervisor Hjalmar Rosengren for his support and encouragement throughout the whole research process. I am also very thankful for support from and discussions with my co-supervisor Jules Lamers. Thanks also to the anonymous reviewers.

Funding Open access funding provided by Uppsala University. This paper was written during my time as a Ph.D. student at University of Gothenburg and Chalmers University of Technology.

Data availability Data sharing is not applicable to this article as no datasets were generated or analysed during the study.

Declarations

Conflict of interest There is no conflict of interests.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article’s Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article’s Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit <http://creativecommons.org/licenses/by/4.0/>.

Appendix A: Foda–Wheeler method

Another way to find the partition function is suggested in [12], by starting from Tsuchiya’s determinant for the partition function on a $2n \times n$ lattice and step by step removing the extra vertical lines by letting the corresponding variables $\mu_j \rightarrow \infty$, following Foda and Wheeler [11].

To do this, we need to choose weights that behave well in the limit. The weights that we have chosen in Sect. 2 will do the job. In the lattice, the local weights are either $w(\lambda - \mu)$ for vertices on the upper part of a double line, or $w(\mu + \lambda)$ for vertices on the lower part, where w is one of a_{\pm} , b_{\pm} and c_{\pm} . When letting $\mu_i \rightarrow \infty$, the local weights become

$$\begin{aligned} \lim_{\mu \rightarrow \infty} a_{\pm}(\lambda - \mu) &= \lim_{\mu \rightarrow \infty} a_{\pm}(\lambda + \mu) = 1, \\ \lim_{\mu \rightarrow \infty} b_{-}(\lambda - \mu) &= e^{2\gamma}, \\ \lim_{\mu \rightarrow \infty} b_{+}(\lambda - \mu) &= 1, \\ \lim_{\mu \rightarrow \infty} b_{-}(\lambda + \mu) &= 1, \\ \lim_{\mu \rightarrow \infty} b_{+}(\lambda + \mu) &= e^{-2\gamma}, \\ \lim_{\mu \rightarrow \infty} c_{-}(\lambda - \mu) &= -e^{\gamma}(e^{\gamma} - e^{-\gamma}), \\ \lim_{\mu \rightarrow \infty} c_{+}(\lambda - \mu) &= -\lim_{\mu \rightarrow \infty} e^{\gamma}(e^{\gamma} - e^{-\gamma}) \frac{1}{e^{2\mu-2\lambda}}, \end{aligned}$$

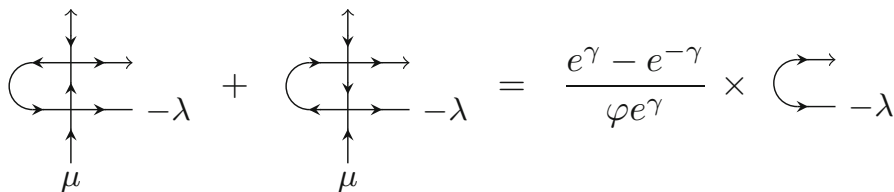


Fig. 12 A k_c turn comes from the sum of a k_+ and k_- turn, when removing the vertical line from the one size bigger lattice

$$\lim_{\mu \rightarrow \infty} c_+(\lambda + \mu) = \frac{e^\gamma - e^{-\gamma}}{e^\gamma},$$

$$\lim_{\mu \rightarrow \infty} c_-(\lambda + \mu) = \lim_{\mu \rightarrow \infty} \left(\frac{e^\gamma - e^{-\gamma}}{e^\gamma} \right) \frac{1}{e^{2\mu + 2\lambda}}.$$

We start from Tsuchiya’s determinant for $2n \times n$ lattices,

$$Z_n(\lambda, \mu) = e^{-\binom{n+1}{2}\gamma} f(\gamma)^n \prod_{i=1}^n [e^{\mu_i + \zeta} f(\mu_i - \zeta) f(2\lambda_i)]$$

$$\times \frac{\prod_{i=1}^n \prod_{j=1}^n [f(\mu_j + \lambda_i) f(\mu_j - \lambda_i)]}{\prod_{1 \leq i < j \leq n} [f(\mu_j + \mu_i) f(\mu_j - \mu_i) f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]} \det_{1 \leq i, j \leq n} M, \tag{A1}$$

where M is an $n \times n$ matrix with

$$M_{ij} = \frac{1}{f(\mu_i + \lambda_j) f(\mu_i - \lambda_j) f(\mu_i + (\lambda_j + \gamma)) f(\mu_i - (\lambda_j + \gamma))}.$$

Tsuchiya considered a diagonal K -matrix, but we have a triangular K -matrix. This does not affect the determinant formula in the $2n \times n$ case, since the ice rule implies that there cannot be any k_c turns, coming from the off-diagonal entry in the K -matrix.

We first consider the 2×1 lattice with reflecting end, since the observations that can be made for this case are important in the general case. This partition function consists of two terms (see Fig. 12),

$$c_+(\lambda + \mu)k_+(\lambda)b_+(\lambda - \mu) + b_+(\lambda + \mu)k_-(\lambda)c_-(\lambda - \mu),$$

which in the limit is $\frac{e^\gamma - e^{-\gamma}}{e^\gamma} (k_+(\lambda) - k_-(\lambda))$. It is easy to check that $k_+(\lambda) - k_-(\lambda) = \frac{k_c(\lambda)}{\phi}$. Hence, letting $\mu \rightarrow \infty$ yields that the partition function becomes $\frac{e^\gamma - e^{-\gamma}}{\phi e^\gamma} k_c(\lambda)$. This makes sense in the picture as well, where removing the vertical line means that both the k_+ and the k_- turn in the 2×1 lattice become a k_c turn in the 2×0 lattice. Conversely, the k_c turn in the smaller lattice can come from removing the vertical line from a state with a k_+ turn, as well as from a state with a k_- turn.

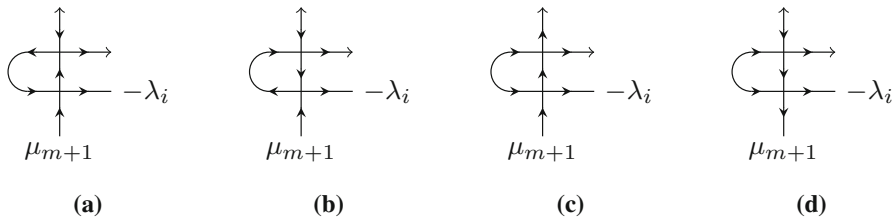


Fig. 13 The k_c turns in a lattice of size $2n \times m$ come from either one of these configurations in the lattice of size $2n \times (m + 1)$, when letting $\mu_{m+1} \rightarrow \infty$

Lemma A.1 Taking the limits $\mu_j \rightarrow \infty$, for $m + 1 \leq j \leq n$, in Tsuchiya’s determinant formula $Z_n(\lambda, \mu)$, one after each other, starting with $\mu_n \rightarrow \infty$, we obtain

$$\begin{aligned}
 & Z_{n,m}(\lambda, \mu_1, \dots, \mu_m) \\
 &= \frac{\varphi^{n-m}}{(1 - e^{-2(n-m)\gamma}) \dots (1 - e^{-4\gamma})(1 - e^{-2\gamma})} \lim_{\mu_n, \mu_{n-1}, \dots, \mu_{m+1} \rightarrow \infty} Z_n(\lambda, \mu).
 \end{aligned}
 \tag{A2}$$

Proof Consider a lattice of general size. In each state, the leftmost vertical line has an odd number of c_{\pm} vertices, because of the DWBC and the ice rule. The c_{\pm} vertices appear as products $c_+(\lambda_i + \mu_m)c_+(\lambda_j - \mu_m)$ or $c_-(\lambda_i + \mu_m)c_-(\lambda_j - \mu_m)$, where both $i = j$ or $i \neq j$ are possible, except for one single $c_+(\lambda_k + \mu_m)$ or $c_-(\lambda_k - \mu_m)$ vertex which is not paired up. In the limit $\mu_m \rightarrow \infty$, the nonzero contribution to $Z_{n,m}(\lambda, \mu_1, \dots, \mu_m)$ comes from the states with exactly one c_{\pm} vertex in the m th column, so we only need to consider these states. In this case, each state contains exactly one of the configurations in Fig. 13a or 13b.

Each state has $n - m$ turns of type k_c . For $m < n$, each state of size $2n \times m$ comes from letting $\mu_{m+1} \rightarrow \infty$ in states of size $2n \times (m + 1)$. In the smaller lattice, all turns are the same as in the bigger lattice, except for at one double row k , where the k_c turn corresponds to either a k_+ turn or a k_- turn in a bigger lattice configuration, just as in the case $n = 1$. Each k_+ or k_- turn in the smaller lattice comes from one of the configurations in Fig. 14 in the bigger lattice. In the limit, each state in Fig. 14, although consisting of two vertices and one turn, also has total weights $k_+(\lambda_i)$ or $k_-(\lambda_i)$, respectively. We are interested in how the partition function for the smaller lattice differs from the bigger lattice, and therefore, we only need to consider the k_c turns in the $2n \times m$ lattice, see Fig. 15. All turns of type k_c below row k come from turns of the type in Fig. 13c. All turns of type k_c above row k come from turns of the type in Fig. 13d. Each turn of type Fig. 13c has weight $e^{-2\gamma}k_c(\lambda_i)$, and each turn of type Fig. 13d has weight $k_c(\lambda_i)$ in the limit.

The row k can be chosen in $n - m$ different ways, so we need to sum over all these possibilities. Putting all of the above together, we conclude that

$$\begin{aligned}
 & Z_{n,m+1}(\lambda, \mu_1, \dots, \mu_{m+1}) \rightarrow \frac{e^\gamma - e^{-\gamma}}{\varphi e^\gamma} \left(1 + \frac{1}{e^{2\gamma}} + \dots + \frac{1}{e^{2(n-m-1)\gamma}} \right) \\
 & \times Z_{n,m}(\lambda, \mu_1, \dots, \mu_m),
 \end{aligned}
 \tag{A3}$$

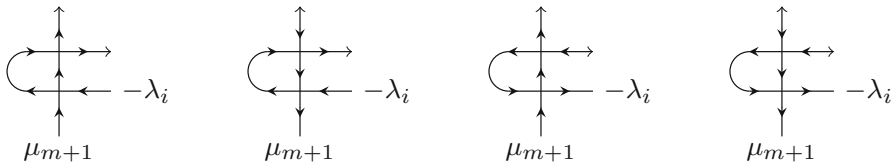
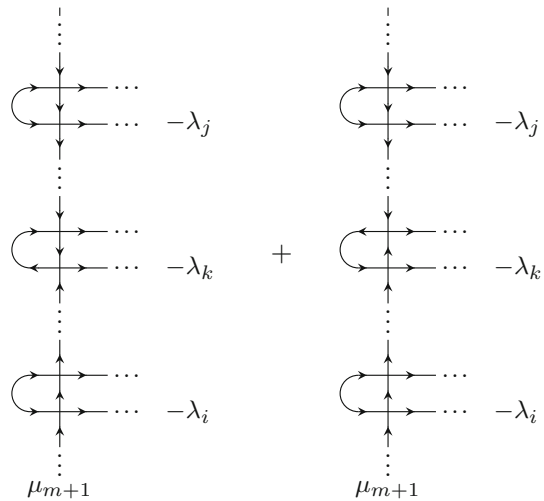


Fig. 14 The turns that are not k_c turns in a lattice of size $2n \times m$ come from either one of these configurations in the lattice of size $2n \times (m + 1)$, when letting $\mu_{m+1} \rightarrow \infty$

Fig. 15 The k_c turns of a state of size $2n \times m$ come from letting $\mu_{m+1} \rightarrow \infty$, i.e. removing the leftmost vertical line, in all states of size $2n \times (m + 1)$ where all turns are the same as in the smaller lattice, except at some row k , where the k_c turn in the smaller lattice corresponds to a k_+ or a k_- turn at row k in the bigger lattice. In the bigger lattice, all k_c double rows below row k are forced to be of the type in Fig. 13c and all k_c double rows above are of the type in Fig. 13d



as $\mu_{m+1} \rightarrow \infty$. By iteration, the lemma follows. □

We now prove the determinant formula by induction.

Proposition A.2 *For the 6V model with DWBC and a partially reflecting end on a lattice of size $2n \times m$, $m \leq n$, the partition function is*

$$\begin{aligned}
 Z_{n,m}(\lambda, \mu_1, \dots, \mu_m) &= \frac{\varphi^{n-m} e^{-\binom{n+1}{2}\gamma} f(\gamma)^n}{(1 - e^{-2(n-m)\gamma}) \dots (1 - e^{-4\gamma})(1 - e^{-2\gamma})} \\
 &\times \frac{\prod_{i=1}^m [e^{\mu_i + \zeta} f(\mu_i - \zeta)] \prod_{i=1}^n f(2\lambda_i) \prod_{i=1}^n \prod_{j=1}^m [f(\mu_j + \lambda_i) f(\mu_j - \lambda_i)]}{\prod_{1 \leq i < j \leq m} [f(\mu_j + \mu_i) f(\mu_j - \mu_i)] \prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]} \\
 &\times \det \begin{pmatrix} F(\lambda_1, \mu_1) & \dots & F(\lambda_n, \mu_1) \\ \vdots & \ddots & \vdots \\ F(\lambda_1, \mu_m) & \dots & F(\lambda_n, \mu_m) \\ C_{n-m-1}(\lambda_1) & \dots & C_{n-m-1}(\lambda_n) \\ \vdots & \ddots & \vdots \\ C_0(\lambda_1) & \dots & C_0(\lambda_n) \end{pmatrix}, \tag{A4}
 \end{aligned}$$

where

$$F(\lambda_i, \mu_j) = \frac{1}{f(\mu_j + \lambda_i)f(\mu_j - \lambda_i)f(\mu_j + (\lambda_i + \gamma))f(\mu_j - (\lambda_i + \gamma))},$$

and

$$C_k(\lambda_i) = \sum_{\substack{k_1, k_2, k_3, k_4 \geq 0 \\ k_1 + k_2 + k_3 + k_4 = k}} e^{(k_1 + k_2 - k_3 - k_4)v_i + (k_1 - k_2 + k_3 - k_4)\gamma} \tag{A5}$$

for $v_i = 2\lambda_i + \gamma$.

Proof We start from Tsuchiya’s determinant formula (A1). Let P_{n-m} be the claim that (A4) holds. We know that P_0 is true, since this is Tsuchiya’s determinant. Assume P_{n-m} . We will prove that P_{n-m+1} also holds. From (A3), we have

$$Z_{n,m-1}(\lambda, \mu_1, \dots, \mu_{m-1}) = \frac{\varphi}{1 - e^{-2(n-m+1)\gamma}} \lim_{\mu_m \rightarrow \infty} Z_{n,m}(\lambda, \mu_1, \dots, \mu_m)$$

Insert (A4). Absorb factors that have to do with μ_m into the m th row of the matrix,

$$\begin{aligned} Z_{n,m-1}(\lambda, \mu_1, \dots, \mu_{m-1}) &= \frac{\varphi^{n-m+1} e^{-\binom{n+1}{2}\gamma} f(\gamma)^n}{(1 - e^{-2(n-m+1)\gamma}) \dots (1 - e^{-4\gamma})(1 - e^{-2\gamma})} \\ &\times \frac{\prod_{i=1}^{m-1} [e^{\mu_i + \zeta} f(\mu_i - \zeta)] \prod_{i=1}^n f(2\lambda_i) \prod_{i=1}^n \prod_{j=1}^{m-1} [f(\mu_j + \lambda_i)f(\mu_j - \lambda_i)]}{\prod_{1 \leq i < j \leq m-1} [f(\mu_j + \mu_i)f(\mu_j - \mu_i)] \prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j)f(\lambda_i + \lambda_j + \gamma)]} \\ &\times \lim_{\mu_m \rightarrow \infty} \det \begin{pmatrix} F(\lambda_1, \mu_1) & \dots & F(\lambda_n, \mu_1) \\ \vdots & \ddots & \vdots \\ F(\lambda_1, \mu_{m-1}) & \dots & F(\lambda_n, \mu_{m-1}) \\ d_1^m(\epsilon_m) & \dots & d_n^m(\epsilon_m) \\ C_{n-m-1}(\lambda_1) & \dots & C_{n-m-1}(\lambda_n) \\ \vdots & \ddots & \vdots \\ C_0(\lambda_1) & \dots & C_0(\lambda_n) \end{pmatrix}, \end{aligned}$$

where

$$d_j^k = \frac{e^{\mu_k + \zeta} f(\mu_k - \zeta) \prod_{i=1}^n [f(\mu_k + \lambda_i)f(\mu_k - \lambda_i)]}{\prod_{i=1}^{k-1} [f(\mu_k + \mu_i)f(\mu_k - \mu_i)]} F(\lambda_j, \mu_k).$$

Write $v_i = 2\lambda_i + \gamma$. Then

$$F(\lambda_i, \mu_j) = \frac{e^{-4\pi}}{(1 - e^{-2\mu_j + v_i + \gamma})(1 - e^{-2\mu_j + v_i - \gamma})(1 - e^{-2\mu_j - v_i + \gamma})(1 - e^{-2\mu_j - v_i - \gamma})}.$$

By Taylor expansion,

$$F(\lambda_i, \mu_j) = \sum_{k=0}^{\infty} C_k(\lambda_i) e^{-2(k+2)\mu_j},$$

with $C_k(\lambda_i)$ as defined in (A5).

Consider

$$d_j^m = \frac{e^{\mu_m + \zeta} f(\mu_m - \zeta) \prod_{i=1}^n [f(\mu_m + \lambda_i) f(\mu_m - \lambda_i)]}{\prod_{i=1}^{m-1} [f(\mu_m + \mu_i) f(\mu_m - \mu_i)]} \sum_{l=0}^{\infty} C_l(\lambda_j) e^{-2(l+2)\mu_m}.$$

By row reduction in the determinant, all terms with $C_l(\lambda_j)$, for $0 \leq l \leq n - m - 1$, can be removed from the sum. The elements of row m become

$$\hat{d}_j^m = \frac{e^{\mu_m + \zeta} f(\mu_m - \zeta) \prod_{i=1}^n [f(\mu_m + \lambda_i) f(\mu_m - \lambda_i)]}{\prod_{i=1}^{m-1} [f(\mu_m + \mu_i) f(\mu_m - \mu_i)]} \sum_{l=n-m}^{\infty} C_l(\lambda_j) e^{-2(l+2)\mu_m}.$$

Now we take the limit,

$$\begin{aligned} \lim_{\mu_m \rightarrow \infty} \hat{d}_j^m &= \lim_{\mu_m \rightarrow \infty} \frac{(1 - e^{2\zeta - 2\mu_m}) \prod_{i=1}^n [(1 - e^{-2\mu_m - 2\lambda_i})(1 - e^{-2\mu_m + 2\lambda_i})]}{\prod_{i=1}^{m-1} [(1 - e^{-2\mu_m - 2\mu_i})(1 - e^{-2\mu_m + 2\mu_i})]} \\ &\quad \times e^{2(n-m+2)\mu_m} \sum_{l=n-m}^{\infty} C_l(\lambda_j) e^{-2(l+2)\mu_m} \\ &= C_{n-m}(\lambda_j), \end{aligned}$$

so P_{n-m+1} is true. □

The last thing left to do is to simplify the last $n - m$ rows of the determinant.

Theorem A.3 (Theorem 3.5) *For the 6V model with DWBC and a partially reflecting end on a lattice of size $2n \times m$, $m \leq n$, the partition function is*

$$\begin{aligned} Z_{n,m}(\boldsymbol{\lambda}, \mu_1, \dots, \mu_m) &= \varphi^{n-m} e^{\binom{m}{2} - nm} \gamma f(\gamma)^m \\ &\quad \times \frac{\prod_{i=1}^m [e^{\mu_i + \zeta} f(\mu_i - \zeta)] \prod_{i=1}^n f(2\lambda_i) \prod_{i=1}^n \prod_{j=1}^m [f(\mu_j + \lambda_i) f(\mu_j - \lambda_i)]}{\prod_{1 \leq i < j \leq m} [f(\mu_j + \mu_i) f(\mu_j - \mu_i)] \prod_{1 \leq i < j \leq n} [f(\lambda_i - \lambda_j) f(\lambda_i + \lambda_j + \gamma)]} \\ &\quad \times \det \begin{pmatrix} F(\lambda_1, \mu_1) & \cdots & F(\lambda_n, \mu_1) \\ \vdots & \ddots & \vdots \\ F(\lambda_1, \mu_m) & \cdots & F(\lambda_n, \mu_m) \\ h(2(n - m - 1)(\lambda_1 + \gamma/2)) & \cdots & h(2(n - m - 1)(\lambda_n + \gamma/2)) \\ \vdots & \ddots & \vdots \\ h(2(\lambda_1 + \gamma/2)) & \cdots & h(2(\lambda_n + \gamma/2)) \\ 1 & \cdots & 1 \end{pmatrix}, \end{aligned}$$

where

$$F(\lambda_i, \mu_j) = \frac{1}{f(\mu_j + \lambda_i)f(\mu_j - \lambda_i)f(\mu_j + (\lambda_i + \gamma))f(\mu_j - (\lambda_i + \gamma))},$$

and where $f(x) = 2 \sinh x$ and $h(x) = 2 \cosh x$.

Remark A.1 In Theorem 3.5, the determinant is expressed in a more compact form.

Proof We start from the determinant (A4) from the last proposition. Focus on the entries of the lower part of the determinant, defined in (A5) as

$$C_k(\lambda_j) = \sum_{\substack{k_1, k_2, k_3, k_4 \geq 0 \\ k_1 + k_2 + k_3 + k_4 = k}} e^{(k_1 + k_2 - k_3 - k_4)v_j + (k_1 - k_2 + k_3 - k_4)\gamma}.$$

These are clearly Laurent polynomials of degree k in e^{v_j} , and they are even in v_j . We have already seen that $C_0(\lambda_j) = 1$. For $k > 0$, the leading coefficient is

$$\sum_{\substack{k_1, k_2 \geq 0 \\ k_1 + k_2 = k}} e^{(k_1 + k_2)v_j + (k_1 - k_2)\gamma} = \frac{f((k + 1)\gamma)}{f(\gamma)}.$$

Thus by row reduction in the determinant, we can replace $C_k(\lambda_j)$ by

$$\frac{f((k + 1)\gamma)}{f(\gamma)}(e^{kv_j} + e^{-kv_j})$$

for $k > 0$. Switching back to the variables λ_j and factoring out $\frac{f((k+1)\gamma)}{f(\gamma)}$ from each row of the lower part of the determinant yields the desired result. \square

References

1. Lieb, E.H.: Residual entropy of square ice. *Phys. Rev.* **162**(1), 162–172 (1967). <https://doi.org/10.1103/physrev.162.162>
2. Sutherland, B.: Exact solution of a two-dimensional model for hydrogen-bonded crystals. *Phys. Rev. Lett.* **19**(3), 103–104 (1967). <https://doi.org/10.1103/physrevlett.19.103>
3. Korepin, V.E.: Calculation of norms of Bethe wave functions. *Commun. Math. Phys.* **86**(3), 391–418 (1982). <https://doi.org/10.1007/bf01212176>
4. Zeilberger, D.: Proof of the alternating sign matrix conjecture. *Electron. J. Combin.* **3**, R13 (1996). <https://doi.org/10.37236/1271>. [arXiv:math/9407211](https://arxiv.org/abs/math/9407211)
5. Mills, W.H., Robbins, D.P., Rumsey, H.: Alternating sign matrices and descending plane partitions. *J. Combin. Theory Ser. A* **34**, 340–359 (1983). [https://doi.org/10.1016/0097-3165\(83\)90068-7](https://doi.org/10.1016/0097-3165(83)90068-7)
6. Izergin, A.G.: Partition function of the six-vertex model in a finite volume. *Soviet Phys. Dokl.* **32**, 878–879 (1987)
7. Izergin, A.G., Coker, D.A., Korepin, V.E.: Determinant formula for the six-vertex model. *J. Phys. A* **25**, 4315–4334 (1992). <https://doi.org/10.1088/0305-4470/25/16/010>
8. Kuperberg, G.: Another proof of the alternating sign matrix conjecture. *Int. Math. Res. Not.* **1996**(3), 139–150 (1996). [arXiv:math/9712207](https://arxiv.org/abs/math/9712207)

9. Tsuchiya, O.: Determinant formula for the six-vertex model with reflecting end. *J. Math. Phys.* **39**, 5946–5951 (1998). <https://doi.org/10.1063/1.532606>. [arXiv:solv-int/9804010](https://arxiv.org/abs/9804010)
10. Kuperberg, G.: Symmetry classes of alternating sign matrices under one roof. *Ann. Math. (2)* **156**(3), 835–866 (2002). <https://doi.org/10.2307/3597283>. [arXiv:math/0008184](https://arxiv.org/abs/math/0008184)
11. Foda, O., Wheeler, M.: Partial domain wall partition functions. *J. High Energy Phys.* **1207**(7), 186 (2012). [https://doi.org/10.1007/jhep07\(2012\)186](https://doi.org/10.1007/jhep07(2012)186). [arXiv:1205.4400v2](https://arxiv.org/abs/1205.4400v2)
12. Foda, O., Zarembo, K.: Overlaps of partial Néel states and Bethe states. *J. Stat. Mech. Theory Exp.* **2016**(2), 023107 (2016). <https://doi.org/10.1088/1742-5468/2016/02/023107>. [arXiv:1512.02533](https://arxiv.org/abs/1512.02533)
13. Pozsgay, B.: Overlaps between eigenstates of the XXZ spin-1/2 chain and a class of simple product states. *J. Stat. Mech. Theory Exp.* **2014**(6), P06011 (2014). <https://doi.org/10.1088/1742-5468/2014/06/P06011>. [arXiv:1309.4593](https://arxiv.org/abs/1309.4593)
14. Colomo, F., Pronko, A.G.: The role of orthogonal polynomials in the six-vertex model and its combinatorial applications. *J. Phys. A* **39**(28), 9015–9033 (2006). <https://doi.org/10.1088/0305-4470/39/28/S15>. [arXiv:math-ph/0602033](https://arxiv.org/abs/math-ph/0602033)
15. Sklyanin, E.K.: Boundary conditions for integrable quantum systems. *J. Phys. A* **21**(10), 2375–2389 (1988). <https://doi.org/10.1088/0305-4470/21/10/015>
16. Hietala, L.: A combinatorial description of certain polynomials related to the XYZ spin chain. *SIGMA* **16**, 101 (2020). <https://doi.org/10.3842/SIGMA.2020.101>. [arXiv:2004.09924v2](https://arxiv.org/abs/2004.09924v2)
17. Koekoek, R., Swarttouw, R.F.: The Askey-scheme of hypergeometric orthogonal polynomials and its q -analogue. Report 98–17, Delft University of Technology, Faculty of Technical Mathematics and Informatics (1998)
18. Krattenthaler, C.: Advanced determinant calculus. In: Foata, D., Han, G.N. (eds.) *The Andrews Festschrift*, pp. 349–426. Springer, Berlin (2001). https://doi.org/10.1007/978-3-642-56513-7_17
19. Gustafson, R., Krattenthaler, C.: Determinant evaluations and $U(n)$ extensions of Heine's ${}_2\phi_1$ -transformations. In: Ismail, M.E.H., Masson, D.R., Rahman, M. (eds.) *Special Functions, q -Series and Related Topics*, Fields Institute Communications, vol. 14, pp. 83–90. American Mathematical Society, Providence (1997)
20. Schlosser, M.: Summation theorems for multidimensional basic hypergeometric series by determinant evaluations. *Discrete Math.* **210**(1–3), 151–169 (2000). [https://doi.org/10.1016/S0012-365X\(99\)00125-9](https://doi.org/10.1016/S0012-365X(99)00125-9)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.