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Magnetic nature of wolframite MgReO_4

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Abstract. Rhenium oxides belonging to the family $A\text{ReO}_4$ where A is a metal cation, exhibit interesting electronic and magnetic properties. In this study we have utilized the muon spin rotation/relaxation ($\mu^+\text{SR}$) technique to study the magnetic properties of the MgReO_4 compound. To the best of our knowledge, this is the first investigation reported on this interesting material, that is stabilized in a wolframite crystal structure using a special high-pressure synthesis technique. Bulk magnetic studies show the onset of an antiferromagnetic (AF) long range order, or a possible singlet spin state at $T_{C1} \approx 90$ K, with a subtle second high-temperature transition at $T_{C2} \approx 280$ K. Both transitions are also confirmed by heat capacity (C_p) measurements. From our $\mu^+\text{SR}$ measurements, it is clear that the sample enters an AF order below $T_{C1} = T_N \approx 85$ K. We find no evidence of magnetic signal above T_N , which indicates that T_{C2} is likely linked to a structural transition. Further, via sensitive zero field (ZF) $\mu^+\text{SR}$ measurements we find evidence of a spin reorientation at $T_{\text{Cant}} \approx 65$ K. This points towards a transition from a collinear AF into a canted AF order at low temperature, which is proposed to be driven by competing magnetic interactions.

1. Introduction

Rhenium oxides [1, 2, 3] of the type $A\text{ReO}_4$, where A is usually a first-row transition metal and Re is found in an oxidation state lower than 7+, are not very common and the scientific literature presents only few examples of such materials [4, 5, 6]. Compounds with Re in an oxidation state 6+ are expected to manifest unusual physical properties due to the strong spin-orbit coupling of the Re $5d^1$ electrons [7, 8, 9], which could introduce quantum fluctuations and lead to multipolar magnetic ordering. The very first material of this kind, MnReO_4 , was synthesized in the 1970's together with MgReO_4 and ZnReO_4 [4]. When first reported, it was suggested that the three samples possessed a wolframite structure, based on the refinement of the unit cell and in qualitative agreement with analogous tungstates and molybdates. While the exact crystal structure, with the refinement of atomic positions, has not been reported yet for MgReO_4 and ZnReO_4 , a recent x-ray diffraction study confirmed the wolframite structure for MnReO_4 [10]. The latter material also manifested an anisotropic electrical resistivity, which



is an attractive feature for potential applications in the development of electrical devices [11]. The wolframite-type structure can be described as a distorted hexagonal close packing of O atoms with each of the A and Re atoms occupying one-fourth of the octahedral voids. The distorted octahedra, containing cations of only one type, form zig-zag chains extended along the *c*-axis. Octahedra of the same type are joined by edges, octahedra of different type are joined by corners. The wolframite structure of AReO_4 is displayed in Fig. 1.

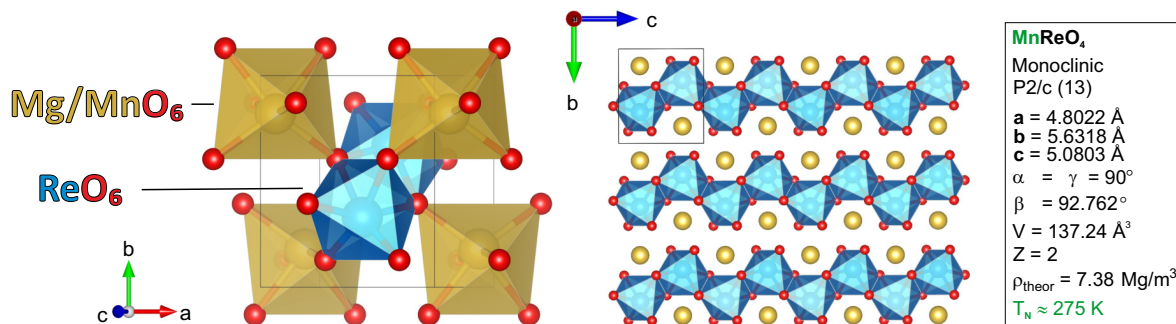


Figure 1. The wolframite crystal structure of MgReO_4 and room temperature crystallographic lattice parameters of the closely related MnReO_4 compound, which display an antiferromagnetic transition at $T_N \approx 275$ K [10].

From a comparison among bond-lengths in different Re and Mn oxides, it was suggested that the electronic configuration in MnReO_4 would be $\text{Mn}^{2+}\text{Re}^{6+}\text{O}_4^{2-}$ with a half filled Mn- $3d^5$ and Re- $5d^1$ system [10] as in $\text{Ba}_{2-x}\text{Sr}_x\text{MnReO}_6$ [12]. For MgReO_4 and ZnReO_4 , the definite oxidation state 2+ for both Mg and Zn ions simply leads to a Re- $5d^1$ system with a delocalized nature of the valence electron [4].

These materials represent a potentially interesting study case and, even though they were already known for decades, their physical properties have not been extensively studied and, for MgReO_4 and ZnReO_4 , no further report beyond the original synthesis paper have been published. In this work we present the very first investigation of the magnetic nature of the wolframite insulator MgReO_4 carried out by both bulk methods as well as muon spin rotation/relaxation technique ($\mu^+\text{SR}$). The occurrence of a static long range magnetic ordering was clearly observed in the muon spectra and the order parameter of the magnetic phase transition was determined. Finally, by the unique sensitivity of the zero field $\mu^+\text{SR}$ protocol we are able to reveal the onset of a possible spin canting inside the antiferromagnetic phase as temperature is further decreased.

2. Experimental Methods

A high pressure synthesis technique ($p = 6$ GPa at 1300°C for 1 hr) was adopted at the National Institute for Materials Science (NIMS) in Japan to produce high-quality powder samples of MgReO_4 . The $\mu^+\text{SR}$ measurements were performed at the General Purpose Surface-Muon Instrument (GPS) of the Swiss Muon Source ($S\mu\text{S}$) at the Paul Scherrer Institute (PSI). Approximately 200 mg of powder was loaded into an envelope made of $30\ \mu\text{m}$ thick aluminum-coated Mylar tape, which was in turn sealed with thin kapton tape to create a hermetic seal from air and moisture. The entire sample handling was conducted inside a helium glove box to avoid any sample degradation. The envelope was suspended in the muon beam using a fork-type sample holder and loaded in a flow cryostat (temperature range: 1.5 K to 300 K). Zero field (ZF) and weak transfer field (wTF) protocols were used to systematically acquire $\mu^+\text{SR}$ time spectra for a set of different temperatures. From the ZF experiment it was possible to investigate the

sample's internal magnetic environment and outline the magnetic order parameter, while the wTF setup allowed to estimate the magnetic transition temperature and the magnetic volume fractions in the sample through the application of an external magnetic field of 50 G, with its direction oriented transversally with respect to the initial direction of the muon spin polarization. More details about the μ^+ SR technique and experimental setup can be found elsewhere [13, 14].

Initial bulk characterizations were carried out on the sample prior to the μ^+ SR experiment. Figure 2(a) shows the temperature dependence of magnetic susceptibility (χ) and inverse χ of MgReO_4 , measured under a magnetic field of 5 kG, together with a corrected curve, which accounts for the presence of magnetic impurities. As the temperature decreases from 320 K, the corrected χ (red line) increases with a change in its slope ($d\chi/dT$) at $T_{C2} = 280$ K, and reaches a local maximum at $T_{C1} = 90$ K. Then, χ abruptly decreases below T_{C1} and levels off to a constant value below about 20 K. This behavior indicates that MgReO_4 enters into either an antiferromagnetically (AF) ordered state or a spin singlet state below T_{C1} , while the nature of the magnetic state in the temperature range between T_{C1} and T_{C2} is not clear. Heat capacity (C_p) measurements [Fig. 2(b-d)] also show the presence of two transitions at $T_{C1} = 90$ K and at $T_{C2} = 275$ K with no temperature hysteresis. The latter is a possible indication for a second order-type phase transition. The heat capacity curve also allowed to estimate the sample's Debye temperature as $\Theta_D = 120$ K.

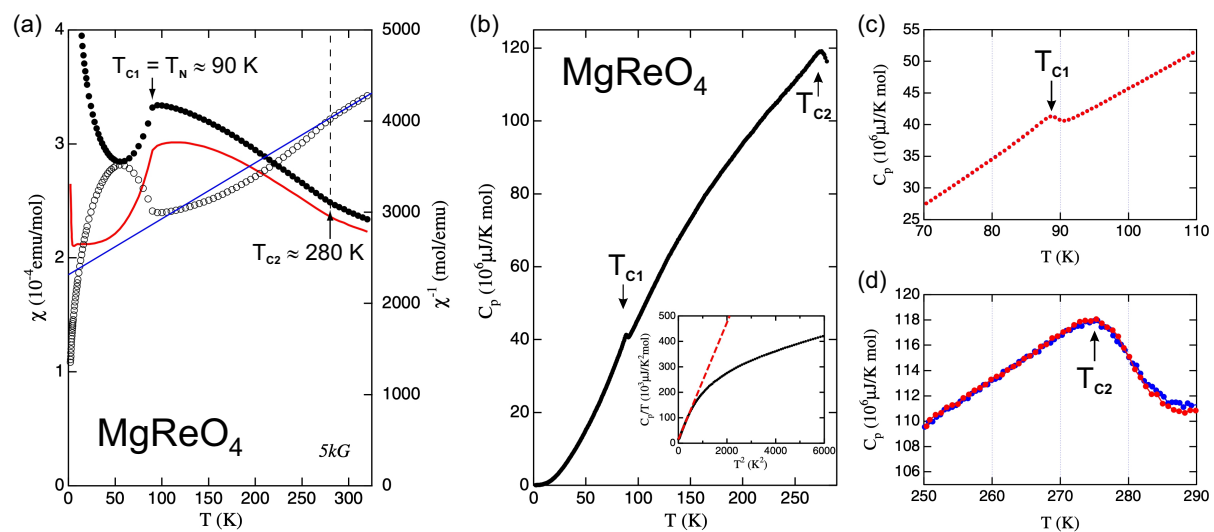


Figure 2. (a) Temperature dependencies of χ (solid circle) and inverse χ (open circle) measured with a magnetic field of 5 kG. A clear antiferromagnetic transition is revealed at $T_{C1} = T_N \approx 90$ K. The red line shows the corrected $\chi(T)$ curve obtained by subtraction of the magnetic impurities contribution. The blue line highlights the deviation of the susceptibility from linearity below T_{C2} . (b) Temperature dependence of the heat capacity (C_p) in MgReO_4 in the full temperature range measured on heating and cooling. The inset shows the fitting curve for the estimation of the Debye temperature. The magnification of $C_p(T)$ around the transition temperatures (c) $T_{C1} = T_N$ and (d) T_{C2} , show that there is no hysteresis between the data measured on heating (red circles) and cooling (blue circles).

3. Results and Discussion

The weak transverse field (wTF) μ^+ SR time spectra at selected temperatures are shown in Figure 3(a). The wTF data is very well fitted [black solid lines in Figure 3(a)] by a muon polarization function described as follows:

$$\begin{aligned} A_0 P_{\text{TF}}(t) &= A_{\text{TF}} \cos(2\pi\nu_{\text{TF}}t + \frac{\pi\phi}{180}) \cdot e^{(-\lambda_{\text{TF}}t)} \\ &+ A_{\text{tail}} \cdot e^{(-\lambda_{\text{tail}}t)} \\ &+ A_{\text{AF}} \cos(2\pi\nu_{\text{AF}}t + \frac{\pi\phi_{\text{AF}}}{180}) \cdot e^{(-\lambda_{\text{AF}}t)} \end{aligned} \quad (1)$$

The muon signal in the wTF regime consists of 3 terms: an oscillating signal with asymmetry A_{TF} reflecting the externally applied weak transverse field, a tail component, commonly found in powder samples (asymmetry A_{tail}) with nearly zero relaxation rate (λ_{tail}), and a fast relaxing term in the early time domain (AF component with asymmetry A_{AF}), which becomes relevant in the temperature range of interest for the magnetic ordering. The phase for the AF component, ϕ_{AF} here shown for completeness, is fixed to zero. By plotting the transverse field asymmetry as a function of temperature [Fig. 3 (b)], it is possible to obtain the magnetic transition temperature $T_N = (82.91 \pm 0.07)$ K (consistent with T_{C1} from susceptibility measurements) as the middle point of the sigmoid fitting function [continuous line in figure 3(b)].

The weak transverse field (wTF) oscillations in the signal are not completely suppressed at base temperature, i.e. the temperature dependence of the wTF-asymmetry does not reach the zero value at base temperature. This means that we have a small fraction of background signal that is not magnetically ordered even at lowest temperature. Most likely this background signal is a combination of muons stopping in the sample holder/envelope and/or a partial degradation of the sample. The latter could happen since the material is sensitive to air and moisture. Two of the powers of the μ^+ SR technique are its high sensitivity to low ordered magnetic moments along with its possibility to acquire data under true zero externally applied magnetic fields. Here we also acquired temperature dependent data using the so-called zero field (ZF) protocol. Selected ZF time spectra are presented in Fig. 3(c), which reveal a clear oscillation at base-temperature ($T = 2$ K). The ZF spectra were fitted with two exponentially damped cosine functions, both having a phase $\phi = 0$:

$$\begin{aligned} A_0 P_{\text{ZF}}(t) &= A_{\text{AF1}} \cos(2\pi\nu_{\text{AF1}}t + \frac{\pi\phi_{\text{AF1}}}{180}) \cdot e^{(-\lambda_{\text{AF1}}t)} \\ &+ A_{\text{AF2}} \cos(2\pi\nu_{\text{AF2}}t + \frac{\pi\phi_{\text{AF2}}}{180}) \cdot e^{(-\lambda_{\text{AF2}}t)} \\ &+ A_{\text{Fast}} \cdot e^{(-\lambda_{\text{Fast}}t)} \\ &+ A_{\text{tail}} \cdot e^{(-\lambda_{\text{tail}}t)}. \end{aligned} \quad (2)$$

Such function indicates the presence of a long range magnetic ordering with a magnetic unit cell commensurate to the symmetry of the crystal lattice. In addition, a fast relaxing exponential term, consistent with the presence of magnetic impurities (possibly metallic Re), and a tail component, usually found in powder samples, are also found in the ZF muon signal. The clear establishment of a static long range magnetic ordering clearly rules out the possibility of dimerization and formation of a spin singlet ground state in MgReO_4 . The presence of two distinct frequencies ν_{AF1} and ν_{AF2} in the sample's internal field distribution usually denotes the presence of two nonequivalent crystallographic muon stopping sites. However, it is also possible that this is a result from one muon stopping site in combination with a complex magnetic structure (that could be either commensurate or incommensurate to the crystal lattice). With increasing temperature, it is found that there is a difference in the temperature dependencies of the two frequencies [see Fig. 3(d)]. While ν_{AF1} is rather abruptly suppressed around 65 K, ν_{AF2}

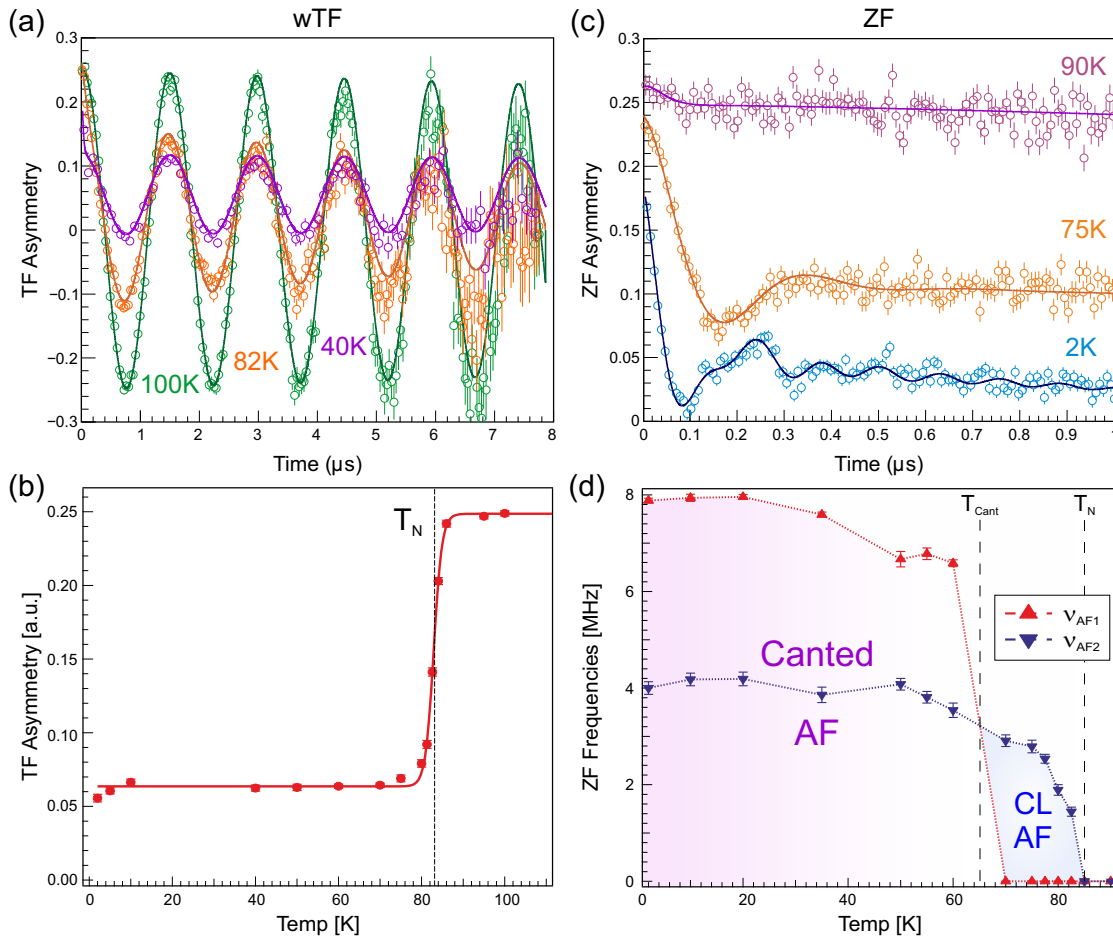


Figure 3. (a) wTF muon time spectra at different temperatures. The black lines are fits to equation 1. (b) Temperature dependence of the transverse field asymmetry in MgReO_4 . The continuous line is a fit to the sigmoid function and the dashed line underlines the magnetic transition temperature. (c) ZF raw muon time spectra in the short time domain for different temperatures. The black lines are fits to equation 2. For clarity of display, the ZF spectra are shifted along the y axis. (d) Temperature dependence of the internal field frequencies. The dotted lines are guides to the eye and dashed vertical lines indicate the antiferromagnetic (AF) order transition ($T_N \approx 85$ K) into a collinear (CL) AF spin order, along with the spin reorientation temperature ($T_{Cant} \approx 65$ K) below which the spins arrange themselves in a proposed canted AF order.

survives up until $T_N \approx 85$ K, displaying a clear second order-like phase transition. Since the wTF data don't show any additional magnetic transitions between 60 K and 70 K the possibility of a magnetic phase separation in the sample can be completely ruled out (determination of magnetic volume fractions is another power of the μ^+ SR technique). The change from one to two frequencies at 65 K could be the result of a structural transition resulting in a change in the muon stopping sites. However, from the C_p data shown in Fig. 2(b-d), there are no indications of a structural transition in this temperature range. Another, here more likely, scenario is that a spin reorientation occurs, creating a complex magnetic structure below 65 K. The zero phase of the ZF oscillations tells us that it should be a commensurate spin structure. Consequently, this suggests that a canted AF structure is formed below $T_{Cant} \approx 65$ K, while a collinear (CL)

AF phase is present for temperatures $T_{\text{Cant}} < T < T_N$ [see Fig. 3(d)].

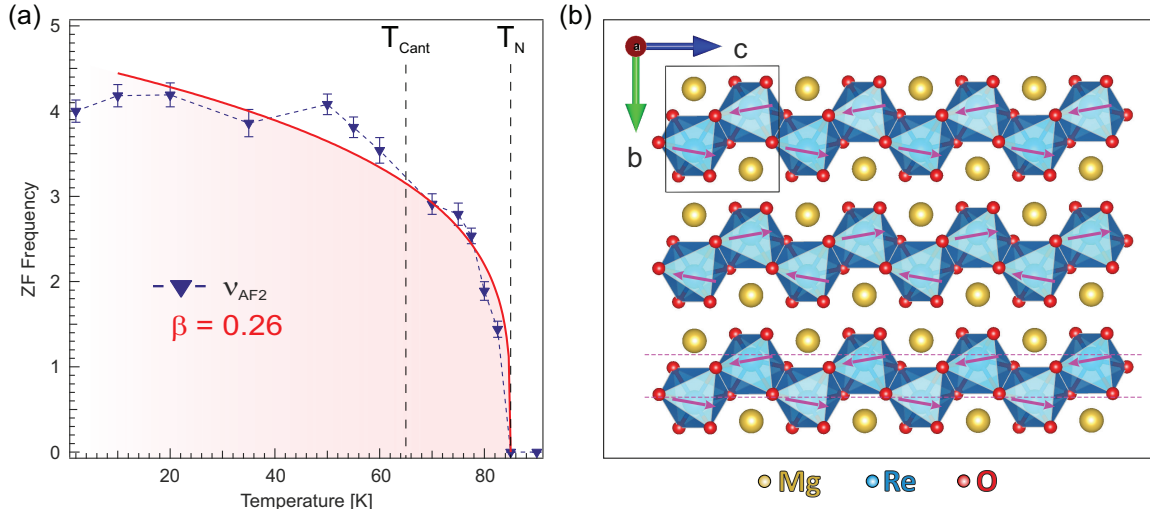


Figure 4. (a) Order parameter of the magnetic phase transition represented by the slower oscillation frequency (ν_{AF2}). The red solid line is a fit to a mean-field theory power law: $\nu \propto (1 - \frac{T}{T_N})^\beta$, which yields $\beta = 0.26$. (b) Schematically presented canted spin structure (canting is exaggerated for display purposes) that is potentially established in MgReO_4 below $T_{\text{Cant}} \approx 65$ K. Dashed horizontal pink lines are guides to the eye to emphasize the canting of the spins away from the collinear spin structure.

The occurrence of canted AF orders is usually a result of competing magnetic interactions, in qualitative agreement with other systems showing a similar behavior [15, 16, 17, 18, 19, 3]. No detailed structural refinement of the atomic positions in MgReO_4 is reported. However, assuming a structure similar to the sister compound MnReO_4 [10], the Re spins in the zig-zag chain could be subjected to exchange interactions with nearest neighbor as well as next nearest neighbor spins. This is similar to the double tungstates analogues of MgReO_4 [20]. The temperature dependence of the ZF- μ^+ SR frequencies reflects the order parameter of the magnetic phase transition in spin ordered materials [21, 22, 18]. By fitting the second order-like muon frequency ν_{AF2} in MgReO_4 , with a mean field theory power law $\nu \propto (1 - \frac{T}{T_N})^\beta$, it is possible to obtain an indication about the nature of magnetic transition in this material [Fig. 4(a)]. The value of the critical exponent obtained from the fit, $\beta = 0.26 \pm 0.03$, suggests that the magnetism in MgReO_4 cannot be categorized under any specific universality class [23]. This could perhaps be taken as an indication for non-conventional magnetic nature of this material. It could of course simply be a result of how we fit the order parameter. In Fig. 4(a) we show the fit to the entire temperature range. However, even if we select to fit only the data in the vicinity of the magnetic phase transition ($T = 60 - 90$ K), we actually obtain the same value $\beta = 0.26$. It is nevertheless clear that we do not have enough temperature point around T_N to make a robust fit nor any definitive statements, and more detailed investigations will therefore be necessary.

Combining the presence of competing magnetic interactions and the itinerant nature of the $\text{Re}^{6+} 5d^1$ electrons, leads to the conclusion that formation of a CL AF order below $T_N \approx 85$ K which is re-arrange into a canted AF spin order below $T_{\text{Cant}} \approx 65$ K is possible in MgReO_4 . As already mentioned, this scenario is also coherent with spin orders that are commensurate to the crystal structure, which is expected from the zero phase in the cosine fit of the ZF data. A very crude schematic representation of one of the possible spin configurations in MgReO_4 , compatibly with these observations, is displayed in Fig. 4(b).

Finally, we do not find any oscillations or clear magnetic signal above the AF transition temperature $T_{C1} = T_N \approx 85$ K. Consequently, the anomaly observed in the susceptibility measurements $[\chi(T)]$ as well as heat capacity $[C_p(T)]$ at $T_{C2} \approx 280$ K (see Fig. 2) is most likely not related to an additional magnetically ordered state. It is instead more likely related to a structural transition, which could easily be checked via, e.g. X-ray diffraction (XRD).

4. Conclusion

In this work we present the the first study of the magnetic properties of the wolframite insulator MgReO_4 , investigated by both bulk methods as well as the μ^+ SR technique. From the detailed analysis of sensitive zero-field μ^+ SR data we find that a collinear antiferromagnetic (AF) order is first formed below $T_N \approx 85$ K. Such spin order is then transformed via a spin-reorientation into a canted AF phase below $T_{\text{Cant}} \approx 65$ K. Our study clearly reveal that such a scenario is very likely, considering the itinerant nature of the $\text{Re}^{6+} 5d^1$ electrons and indications for the presence of competing magnetic interactions. Nevertheless, further experimental investigations (e.g. neutron diffraction), as well as modeling of muon stopping sites, would be needed to reveal the true magnetic ground state and detailed spin structure of this interesting material. Finally, computer modeling could be utilized for extracting the theoretical magnetic exchange interactions ($J/J'/\dots$) in order to also clarify the possible formation of multipolar magnetic ordering in MgReO_4 .

5. Acknowledgments

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