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# Linking ice and gas in the Coronet cluster in Corona Australis

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## ABSTRACT

**Context.** During the journey from the cloud to the disc, the chemical composition of the protostellar envelope material can be either preserved or processed to varying degrees depending on the surrounding physical environment.

**Aims.** This work aims to constrain the interplay of solid (ice) and gaseous methanol (CH<sub>3</sub>OH) in the outer regions of protostellar envelopes located in the Coronet cluster in Corona Australis (CrA), and assess the importance of irradiation by the Herbig Ae/Be star R CrA. CH<sub>3</sub>OH is a prime test case as it predominantly forms as a consequence of the solid-gas interplay (hydrogenation of condensed CO molecules onto the grain surfaces) and it plays an important role in future complex molecular processing.

**Methods.** We present 1.3 mm Submillimeter Array (SMA) and Atacama Pathfinder Experiment (APEX) observations towards the envelopes of four low-mass protostars in the Coronet cluster. Eighteen molecular transitions of seven species were identified. We calculated CH<sub>3</sub>OH gas-to-ice ratios in this strongly irradiated cluster and compared them with ratios determined towards protostars located in less irradiated regions such as Serpens SVS 4 in Serpens Main and the Barnard 35A cloud in the  $\lambda$  Orionis region.

**Results.** The CH<sub>3</sub>OH gas-to-ice ratios in the Coronet cluster vary by one order of magnitude (from  $1.2 \times 10^{-4}$  to  $3.1 \times 10^{-3}$ ) which is similar to less irradiated regions as found in previous studies. We find that the CH<sub>3</sub>OH gas-to-ice ratios estimated in these three regions are remarkably similar despite the different UV radiation field intensities and formation histories.

**Conclusions.** This result suggests that the overall CH<sub>3</sub>OH chemistry in the outer regions of low-mass envelopes is relatively independent of variations in the physical conditions and hence that it is set during the prestellar stage.

**Key words.** ISM: molecules – stars: protostars – astrochemistry – molecular processes – ISM: individual objects: R CrA

## 1. Introduction

Solar-type stars and planets form inside molecular clouds, when dense cores undergo gravitational collapse. While collapsing, the gas and dust constituting the clouds are assembled into infalling envelopes, streamers, and circumstellar discs, which supply the fundamental ingredients for planet formation (see e.g. Pineda et al. 2023). Recent observations (e.g. Andrews et al. 2018; Harsono et al. 2018; Keppler et al. 2018; Segura-Cox et al. 2020) suggest that planets form earlier than previously thought (a few  $10^3$  yr; Tychoniec et al. 2020) and in tandem with their host stars (Alves et al. 2020). In this context, it is still unclear whether the composition of the molecular cloud material is preserved when becoming part of planets or, instead, entirely thermally processed losing the natal cloud's chemical fingerprint (see e.g. Jørgensen et al. 2020; van Dishoeck & Bergin 2020 for recent reviews).

To address this query, it is necessary to study how the environment impacts the physical and chemical properties of embedded protostars and their discs. For instance, external irradiation (Winter et al. 2020) and cosmic ray ionisation (Kuffmeier et al. 2020) might shape the chemical and physical evolution of forming low-mass stars and their discs (van Terwisga et al. 2020; Haworth et al. 2021 and references therein) or leading to less massive discs (Cazzoletti et al. 2019). In this paper we investigate the variations of the methanol gas and ice towards deeply embedded sources in Corona Australis (CrA), to study the effects of external irradiation on the gas-to-ice ratios.

CH<sub>3</sub>OH is the most suitable molecule for the main aim of this work because it is abundantly detected in both solid and gas phases (Boogert et al. 2015; McGuire 2022). Most importantly, it is primarily formed in the solid state (Watanabe & Kouchi 2002; Fuchs et al. 2009; Qasim et al. 2018; Simons et al. 2020;

**Table 1.** Overview of the Coronet cluster objects.

Object	RA (J2000)	Dec (J2000)	Designation <sup>(a)</sup>	Other names <sup>(b)</sup>
IRS5A	19:01:48.03	−36:57:22.2	Class I YSO	CrA-19
IRS5N	19:01:48.46	−36:57:14.7	Class I YSO	CrA-20
HH100 IRS1	19:01:50.56	−36:58:08.9	Herbig-Haro object	–
IRS1	19:01:50.68	−36:58:09.7	Class I YSO	VSS15, TS 2.6
R CrA	19:01:53.67	−36:57:08.0	Herbig Ae/Be star	–
SMM1C	19:01:55.29	−36:57:17.0	Class I YSO	B9, Brown 9
IRS7A	19:01:55.32	−36:57:21.9	Class I YSO	IRS7W, IRS7
SMM1A	19:01:55.60	−36:57:43.0	Pre-stellar core	–
IRS7B	19:01:56.40	−36:57:28.3	Class I YSO	SMM1B, IRS7E
SMM2	19:01:58.54	−36:57:08.5	Class I YSO	CrA-43, WMB55
T CrA	19:01:58.78	−36:57:49.9	T Tauri star	–

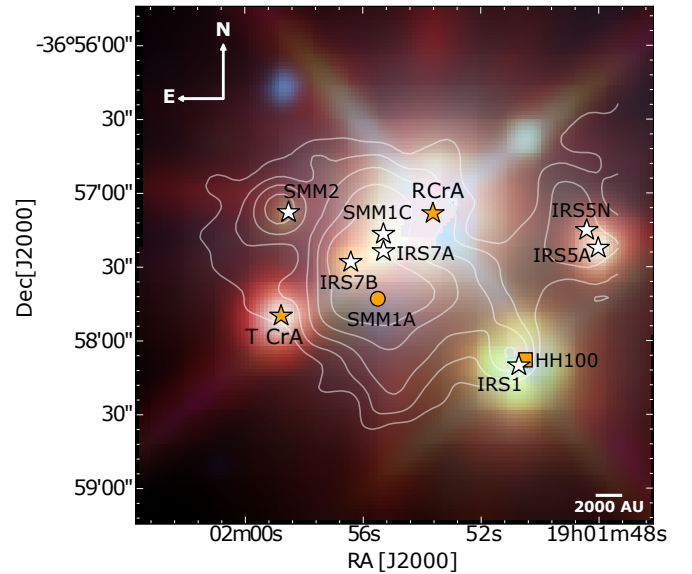
**Notes.** All the coordinates are taken from Peterson et al. (2011), except for SMM1A which comes from Nutter et al. (2005) and HH100 IRS1 from Boogert et al. (2008). <sup>(a)</sup>The YSO class is assigned from the spectral index  $\alpha$  reported in Peterson et al. (2011); the objects designated as ‘Class I YSO’ are Class I or younger. <sup>(b)</sup>From Peterson et al. (2011) and Lindberg & Jørgensen (2012).

Santos et al. 2022) as its gas-phase formation pathways are considerably less efficient (Roberts & Millar 2000; Garrod & Herbst 2006; Geppert et al. 2006). Finally, CH<sub>3</sub>OH is regarded as the gateway species for the formation of complex organic molecules both in the solid state (e.g. Öberg et al. 2009b; Chuang et al. 2016; Fedoseev et al. 2017) and in the gas phase (e.g. Shannon et al. 2014; Balucani et al. 2015).

Corona Australis is one of the nearest regions with ongoing star formation, located at a distance of  $149.4 \pm 0.4$  pc as estimated from *Gaia* Data Release 2 measurements (Galli et al. 2020). The most recent census of the cloud counts 393 young stellar object (YSO) candidates (Esplin & Luhman 2022): these are relatively evolved, mainly Class II and III sources, and they are predominantly concentrated in the most extinguished region, the ‘head’ (Peterson et al. 2011; Alves et al. 2014). The correlation between the star formation activity and the head-tail structure of CrA has been investigated by Dib & Henning (2019), who found that the spatial distribution of dense cores is a consequence of the physical conditions of the large-scale environment present at the time the cloud assembled. The youngest population of YSOs (Class 0/I) is situated in the Coronet cluster (Taylor & Storey 1984; Forbrich et al. 2007; Forbrich & Preibisch 2007), also associated with the luminous Herbig Ae/Be star R CrA (Peterson et al. 2011).

Figure 1 displays a section of the Coronet cluster and a summary of the principal cluster members is provided in Table 1. The Herbig Ae/Be R CrA (spectral type B5–B8; Gray et al. 2006; Bibo et al. 1992) is the brightest star in this very young cluster that has an estimated age of 0.5–1 Myr (Sicilia-Aguilar et al. 2011). Due to the variable nature of R CrA, its stellar mass and luminosity are uncertain (Mesa et al. 2019). Recent Very Large Telescope (VLT)/SPHERE observations of R CrA resolved a companion at a separation of  $\sim 0.156''$ – $0.184''$  and complex extended jet-like structures around the star (Mesa et al. 2019). A second variable star, the T Tauri star T CrA is present in the Coronet cluster, approximately  $30''$  to the south-east of R CrA (Herbig 1960; Taylor & Storey 1984). The region surrounding the two variable stars harbours seven identified Class 0/I YSOs (Table 1).

The Coronet cluster members have been at the centre of active multi-wavelength research, mostly aimed at the characterisation of the properties of YSOs at very early stages of stellar



**Fig. 1.** Three-colour image of the Coronet cluster (WISE 3.4  $\mu\text{m}$ , blue, 4.6  $\mu\text{m}$ , green, and 12  $\mu\text{m}$ , red, bands; Wright et al. 2010) overlaid with the SCUBA 850  $\mu\text{m}$  density flux from Nutter et al. (2005); contours are in decreasing steps starting at the peak flux ( $3.7 \text{ Jy beam}^{-1}$ ) and subtracting 30% from the previous level. The white stars mark the positions of Class 0/I YSOs in the Coronet region (Peterson et al. 2011; Nutter et al. 2005), whereas the orange stars indicate the Herbig Ae/Be star R CrA and the T Tauri star T CrA. The pre-stellar core candidate SMM1A is indicated with an orange circle and the orange square represents the Herbig Haro object HH100.

evolution. Some of the most studied objects are IRS7B, IRS7A, SMM1A, SMM1C and SMM2 (Nutter et al. 2005; Groppi et al. 2007; Miettinen et al. 2008; Chen & Arce 2010; Peterson et al. 2011). The spectral energy distributions (SEDs) of the cluster members have been investigated by Groppi et al. (2007), suggesting that SMM1C is a Class 0 YSO and IRS7B is a transitional Class 0/I object. IRS7A and SMM2 are likely Class I sources (Peterson et al. 2011) and SMM1A is classified as a pre-stellar core candidate (Nutter et al. 2005; Chen & Arce 2010).

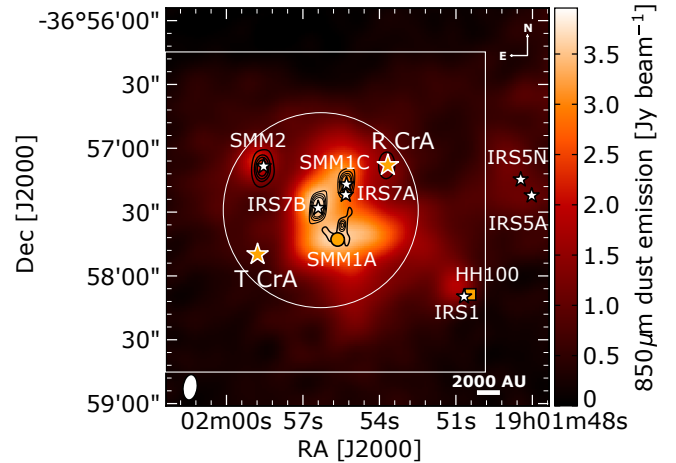
IRS1, IRS5N and IRS5A are suggested as Class I YSO based on sub-millimetre and infrared observations (Peterson et al. 2011).

Along with the identification and the evolutionary stage designation of the YSOs in the Coronet cluster, their chemical evolution has been the subject of a large number of studies over the last decades. The line-rich spectra of IRS7B and IRS7A have been investigated with the Atacama Pathfinder EXperiment (APEX) single-dish telescope by Schöier et al. (2006), who reported different kinetic temperatures for formaldehyde ( $\text{H}_2\text{CO}$ ) and methanol ( $>30$  K for  $\text{H}_2\text{CO}$  and  $\approx 20$  K for  $\text{CH}_3\text{OH}$ ). Lindberg & Jørgensen (2012) suggested that the high temperatures in the region ( $>30$  K) traced by the emission of  $\text{H}_2\text{CO}$ , are caused by external irradiation from the Herbig Ae/Be star R CrA. With the purpose of analysing the impact of external irradiation on the molecular inventory of low-mass protostars, systematic unbiased line surveys have been carried out firstly towards IRS7B with the Atacama Submillimeter Telescope Experiment (ASTE) by Watanabe et al. (2012), and extended to other Coronet cluster members using APEX by Lindberg et al. (2015). These systematic surveys confirmed that external irradiation can affect the chemical nature of protostars, particularly by enhancing the abundances of Photon-Dominated Regions (PDRs) tracers (such as CN,  $\text{C}_2\text{H}$ , and  $\text{c-C}_3\text{H}_2$ ). High-resolution Atacama Large Millimeter/submillimeter Array (ALMA) observations found a low column density of  $\text{CH}_3\text{OH}$  in the inner hot regions of the protostellar envelopes which could either be due to the suppressed  $\text{CH}_3\text{OH}$  formation in ices at the higher temperatures in the region or due to the envelope structures on small scales being affected by the presence of a disc (Lindberg et al. 2014). The latter interpretation is supported by recent observations (Artur de la Villarmois et al. 2018; van Gelder et al. 2022) and models (Nazari et al. 2022b).

Along with the gas-phase molecular species, the composition of interstellar ices towards a few objects in the Coronet cluster (IRS5A, IRS5B, HH100, IRS1, IRS7A, IRS7B) has been modelled (Siebenmorgen & Gredel 1997) and investigated in the near-infrared ( $3\text{--}5\ \mu\text{m}$ ) with the VLT (Pontoppidan et al. 2003), in the mid-infrared ( $5\text{--}8\ \mu\text{m}$ ) using the Infrared Space Observatory (Keane et al. 2001) and as part of the *Spitzer* Space Telescope c2d Legacy Program ( $5\text{--}8\ \mu\text{m}$  by Boogert et al. 2008;  $8\text{--}10\ \mu\text{m}$  by Bottinelli et al. 2010, and  $14.5\text{--}16.5\ \mu\text{m}$  by Pontoppidan et al. 2008). Ice features attributed to  $\text{H}_2\text{O}$ , CO,  $\text{CO}_2$ ,  $\text{NH}_3$ ,  $\text{CH}_3\text{OH}$  were securely identified in the astronomical spectra of the targeted sources.

This paper explores the variations of  $\text{CH}_3\text{OH}$  gas and ice abundances, investigating the effects of external irradiation on the gas-phase and solid-state chemistries of the youngest objects in the Coronet cluster. We present Submillimeter Array (SMA) observations towards five Coronet cluster members, thus complementing previous observations at 1.3 mm (Lindberg & Jørgensen 2012) by increasing the number of detected chemical species in this region. To probe the  $\text{CH}_3\text{OH}$  large-scale emission, an APEX  $150'' \times 150''$  map of the  $\text{CH}_3\text{OH}$   $J_K = 5_K\text{--}4_K$  Q-branch is also presented. The column densities of  $\text{H}_2\text{O}$  and  $\text{CH}_3\text{OH}$  ice are taken from Boogert et al. (2008). Furthermore, the dependence of gas-to-ice ratios on the physical conditions are addressed, by determining gas-to-ice ratios of  $\text{CH}_3\text{OH}$  in this strongly irradiated cluster and comparing them with ratios obtained towards less irradiated star-forming regions (Perotti et al. 2020, 2021).

The paper is laid out as follows. Section 2 describes both the SMA and APEX observational setups, and the data reduction strategy. In Sect. 3, we present our key observational results, while the variations on the ice and gas abundance are analysed



**Fig. 2.** SMA continuum at 1.3 mm (contours) overlaid with SCUBA 850  $\mu\text{m}$  density flux from Nutter et al. (2005). The contours start at  $5\sigma$  and continue in steps of  $5\sigma$  ( $\sigma = 18\ \text{mJy beam}^{-1}$ ). The empty circle indicates the size of the SMA primary beam, whereas the SMA synthesised beam is shown with a white ellipse in the bottom left corner. The empty rectangle shows the map covered by APEX observations. The stars mark the position of the objects located in the Coronet cluster; refer to Table 1 for their identification.

in Sect. 4. Section 5 discusses the determined  $\text{CH}_3\text{OH}$  gas-to-ice ratios and Sect. 6 summarises our conclusions and lists avenues for future studies.

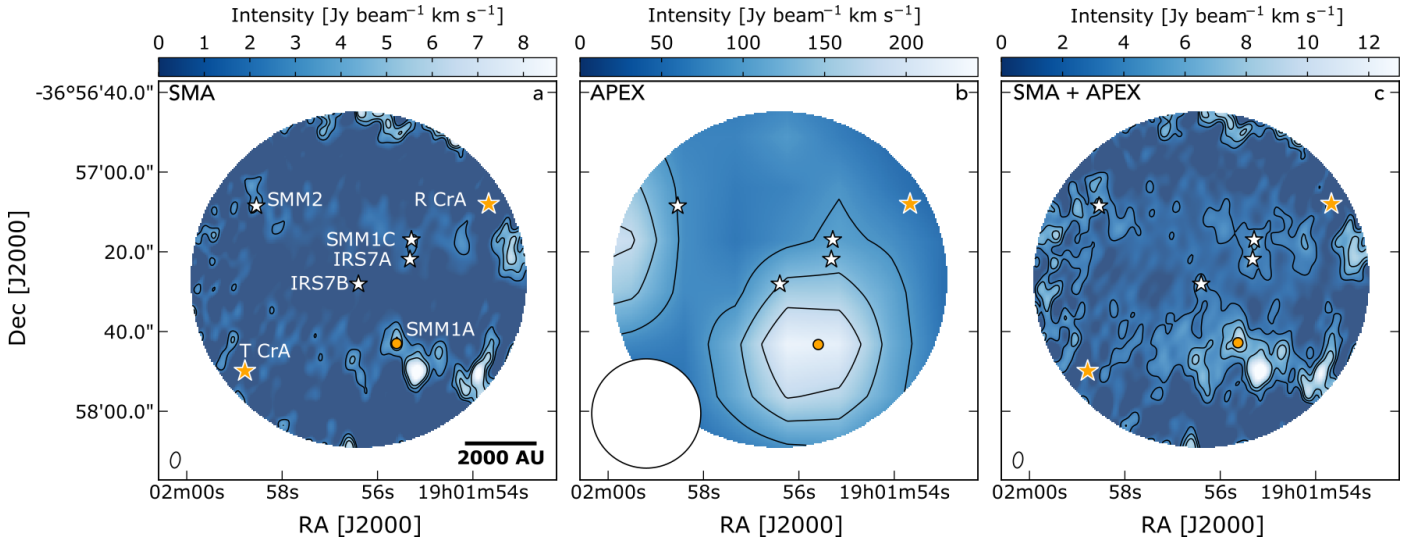
## 2. Observations

The Coronet cluster was observed with the SMA (Ho et al. 2004) on February 25, 2020, April 28, 30, and May 19, 31, 2021 in the compact configuration, resulting in projected baselines between 17–72 m. The region was covered by one pointing centred on IRS7B with coordinates  $\alpha_{\text{J2000}} = 19^{\text{h}}01^{\text{m}}56^{\text{s}}.40$ ,  $\delta_{\text{J2000}} = -36^{\circ}57'27''.00$  (Fig. 2). The SWARM correlator provided a total frequency coverage of 36 GHz, with the lower sideband covering frequencies from 210.3 to 228.3 GHz and upper sideband from 230.3 to 248.3 GHz. The spectral resolution was 0.6 MHz corresponding to  $0.7\ \text{km s}^{-1}$ .

Data calibration and imaging were performed with the CASA v. 5.7.2 package<sup>1</sup> (McMullin et al. 2007; Bean et al. 2022). The complex gains were calibrated through observations of the quasars 1924-292 and 1957-387 and the bandpass through observations of the bright quasar 3C279. The overall flux calibration was carried out through observations of Vesta and Callisto. Imaging was done using the `tclean` algorithm and a Briggs weighting with a robust parameter of 0.5. The typical synthesised beam size of the SMA dataset was  $4''.9 \times 2''.6$  ( $400\text{--}700\ \text{AU}$ ). The position angle (PA) for  $\text{CH}_3\text{OH}$ , our main molecule of interest, at 241.791 GHz was  $-8.2^\circ$ . The PA varies of  $\approx 4.5^\circ$  across the covered frequencies.

Additional information on the  $\text{CH}_3\text{OH}$  large-scale emission towards the Coronet cluster was provided by observations with APEX (Güsten et al. 2006) during the nights of October 25–November 4, 2020. The pointing was the same as reported for the SMA observations. The achieved map size is  $150'' \times 150''$ , fully covering the region observed by the SMA. The frequencies probed by the APEX observations are in the range 235.1–243 GHz, corresponding to the SMA upper side

<sup>1</sup> <http://casa.nrao.edu/>



**Fig. 3.** Primary beam corrected integrated intensity maps for  $\text{CH}_3\text{OH } J = 5_0-4_0 \text{ A}^+$  transition ( $E_u = 34.8 \text{ K}$ ) at 241.791 GHz detected by the SMA (a), by the APEX telescope (b), and in the combined interferometric SMA and single-dish APEX data (c). All lines are integrated between 6 and 12.5  $\text{km s}^{-1}$ . Contours start at  $5\sigma$  ( $\sigma_{\text{SMA}} = 0.16 \text{ Jy beam}^{-1} \text{ km s}^{-1}$ ,  $\sigma_{\text{APEX}} = 2 \text{ Jy beam}^{-1} \text{ km s}^{-1}$ ,  $\sigma_{\text{SMA+APEX}} = 0.17 \text{ Jy beam}^{-1} \text{ km s}^{-1}$ ) and follow in steps of  $5\sigma$ . The circular field-of-view corresponds to the primary beam of the SMA observations. The synthesised beams are displayed in white in the bottom left corner of each panel.

band observations, with a frequency resolution of 0.061 MHz (0.076  $\text{km s}^{-1}$ ). The resulting beam-size of the observations is 27.4'' for  $\text{CH}_3\text{OH}$ . Data reduction was carried out with the GILDAS package CLASS<sup>2</sup>. Subsequently, the reduced APEX and SMA datasets were combined in CASA v. 5.7.2 using the feathering technique, following the procedure presented in Appendix B.1. of Perotti et al. (2020). The combination of SMA and APEX short-spacing data resulted in mapping the outer regions of protostellar envelopes in the Coronet cluster on scales from  $\sim 400$ –10 000 AU.

### 3. Results

Figure 2 displays the region targeted by the SMA and APEX observations. The positions of the SMA 1.3 mm continuum peaks coincides with the location of the young stellar objects in the Coronet cluster. The strongest emission is observed towards the two Class 0 sources SMM1C and IRS7B and the Class I objects SMM2 and IRS7A. Fainter emission north of the pre-stellar core candidate SMM1A (Chen & Arce 2010) is also seen. Apart from the continuum, line emission of seven species was identified (Table A.1). These results are presented and discussed in Appendix A (Figs. A.1–A.3). In the remainder of this section we uniquely focus on the observed methanol ( $\text{CH}_3\text{OH}$ ) emission, our molecule of interest.

Amongst the molecular species detected in the Coronet cluster, methanol is the most suited for the main aim of this work: assessing the effect of external irradiation on the gas-phase and solid-state (ice) chemistries in the outer regions of protostellar envelopes. These are the regions probed in the presented SMA and APEX observations, at radii  $r > 400$  AU from the protostar. Here, gas-phase methanol is not thermally desorbed and it is a tracer of energetic input such as outflows releasing frozen  $\text{CH}_3\text{OH}$  to the gas phase (e.g. Tychoniec et al. 2021). The latter is due to the fact that the most efficient methanol formation pathway in star-forming regions ( $A_V > 9$  mag) occurs

on ice-coated dust grains (Watanabe & Kouchi 2002; Fuchs et al. 2009; Simons et al. 2020; Santos et al. 2022). Consequently, the observed non-thermal  $\text{CH}_3\text{OH}$  emission in the outer regions of protostellar envelopes does not trace the bulk of the ices, and instead, it probes a fraction of  $\text{CH}_3\text{OH}$  ice sputtered in outflows or photodesorbed (see Kristensen et al. 2010 for an observational overview of  $\text{CH}_3\text{OH}$  non-thermal desorption processes).

Figure 3 presents primary beam corrected integrated intensity (moment 0) maps of the  $\text{CH}_3\text{OH } J = 5_0-4_0 \text{ A}^+$  line for the SMA, APEX, and the combination of interferometric and single-dish data (SMA+APEX). The combined moment 0 maps of the other  $\text{CH}_3\text{OH}$  transitions belonging to the  $J = 5_K-4_K$  Q-branch are displayed in Fig. A.2. We specifically targeted this branch because its multiple transitions are conveniently observable in a narrow spectral range (0.2 GHz; Table 2) and they are associated to  $E_u$  from 34.8 to 60.7 K, well below the typical  $\text{CH}_3\text{OH}$  sublimation temperature as inferred from both observations and laboratory experiments (70–130 K; e.g. Kristensen et al. 2010; Penteado et al. 2017). The latter condition has to be satisfied in order to investigate methanol non-thermal desorption processes. Additionally, observations of  $\text{CH}_3\text{OH } J = 5_K-4_K$  Q-branch make it possible to directly compare our results with previous work on  $\text{CH}_3\text{OH}$  gas-to-ice ratios in nearby star-forming regions using the same facilities, observational strategy, and spectral setups (Perotti et al. 2020, 2021).

By looking at the SMA moment 0 map (Fig. 3; panel a), it is clear that the interferometric observations do not sufficiently recover the spatially extended emission originating from the Coronet cluster. In contrast, they filter out  $\sim 90\%$  of the extended emission detected in the APEX data. Unlike for the other species detected in the SMA data, we here want to perform a quantitative analysis of the  $\text{CH}_3\text{OH}$  emission therefore, we combine the SMA data with the APEX map to avoid underestimating the  $\text{CH}_3\text{OH}$  column densities by up to one order of magnitude.

Similarly to the  $\text{CH}_3\text{OH } J = 4_2-3_1 \text{ E}$  transition (Fig. A.2; panel g), the  $\text{CH}_3\text{OH } J = 5_0-4_0 \text{ A}^+$  emission is tracing the condensations associated to the prestellar core candidate SMM1A.

<sup>2</sup> <http://www.iram.fr/IRAMFR/GILDAS>

**Table 2.** Identified methanol transitions in the SMA and APEX datasets.

Species	Transition	Frequency <sup>(a)</sup> (GHz)	$A_{ul}^{(a)}$ (s <sup>-1</sup> )	$g_u^{(a)}$	$E_u^{(a)}$ (K)	$n_{cr}^{(b)}$ (cm <sup>-3</sup> )
CH <sub>3</sub> OH	5 <sub>+0</sub> – 4 <sub>+0</sub> E	241.700	$6.04 \times 10^{-5}$	11	47.9	$1.7 \times 10^6$
	5 <sub>-1</sub> – 4 <sub>-1</sub> E	241.767	$5.81 \times 10^{-5}$	11	40.4	$4.8 \times 10^7$
	5 <sub>0</sub> – 4 <sub>0</sub> A <sup>+</sup>	241.791	$6.05 \times 10^{-5}$	11	34.8	$4.7 \times 10^5$
	5 <sub>+1</sub> – 4 <sub>+1</sub> E	241.879	$5.96 \times 10^{-5}$	11	55.9	$1.5 \times 10^7$
	5 <sub>-2</sub> – 4 <sub>-2</sub> E	241.904	$5.09 \times 10^{-5}$	11	60.7	$2.3 \times 10^8$

**Notes.** <sup>(a)</sup>The spectroscopic data are from Xu et al. (2008). They were taken from the Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001, 2005; Endres et al. 2016) and the Jet Propulsion Laboratory catalogue (Pickett et al. 1998). <sup>(b)</sup>Calculated using a collisional temperature of 30 K and collisional rates from Rabli & Flower (2010). The latter were taken from the Leiden Atomic and Molecular Database (LAMDA; Schöier et al. 2005).

**Table 3.** Total ice and gas column densities ( $N$ ) and fractional abundances ( $X$ ) relative to H<sub>2</sub> towards the Coronet cluster members.

Object	$N_{H_2O}^{ice (a)}$ (10 <sup>18</sup> cm <sup>-2</sup> )	$N_{CH_3OH}^{ice}$ (10 <sup>17</sup> cm <sup>-2</sup> )	$N_{CH_3OH}^{gas (b)}$ (10 <sup>15</sup> cm <sup>-2</sup> )	$N_{H_2}^{SCUBA (d)}$ (10 <sup>23</sup> cm <sup>-2</sup> )	$X_{H_2O}^{ice}$ (10 <sup>-5</sup> )	$X_{CH_3OH}^{ice}$ (10 <sup>-6</sup> )	$X_{CH_3OH}^{gas}$ (10 <sup>-8</sup> )
IRS5A	$3.58 \pm 0.26$	$2.36 \pm 0.57$	$0.05 \pm 0.02^{(c)}$	$0.55 \pm 0.06$	$6.51 \pm 0.80$	$4.29 \pm 1.12$	$0.09 \pm 0.04$
HH100 IRS1	$2.45 \pm 0.24$	<2.38	$0.07 \pm 0.03^{(c)}$	$0.67 \pm 0.07$	$3.66 \pm 0.51$	<3.55	$0.10 \pm 0.05$
SMM1C	–	–	$1.95 \pm 0.23$	$1.58 \pm 0.16$	–	–	$1.23 \pm 0.19$
IRS7A	$10.89 \pm 1.92$	<4.14	$1.28 \pm 0.15$	$1.48 \pm 0.15$	$7.36 \pm 1.49$	<2.98	$0.86 \pm 0.13$
SMM1A	–	–	$3.54 \pm 0.43$	$1.90 \pm 0.19$	–	–	$1.86 \pm 0.29$
IRS7B	$11.01 \pm 1.97$	$7.49 \pm 0.33$	$1.91 \pm 0.23$	$1.92 \pm 0.19$	$5.73 \pm 1.18$	$3.90 \pm 0.43$	$0.99 \pm 0.16$
SMM2	–	–	$1.64 \pm 0.20$	$1.09 \pm 0.11$	–	–	$1.50 \pm 0.24$

**Notes.** <sup>(a)</sup>Boogert et al. (2008). <sup>(b)</sup>CH<sub>3</sub>OH column densities calculated using  $T_{rot} = 30$  K. <sup>(c)</sup>Lindberg et al. (2015). <sup>(d)</sup>H<sub>2</sub> column densities assuming  $T_{dust} = 30$  K. The errors on the H<sub>2</sub> column densities are estimated from the 10% flux calibration uncertainty at 850  $\mu$ m. The ‘–’ symbol indicates the sources for which ice column densities and abundances are not available.

The emission is extended in the APEX map (Fig. 3; panel b), and confined to two main regions, around SMM1A and to the east of SMM2. The peak intensity is observed also in this case in the vicinity of SMM1A. The same pattern is visible in the SMA + APEX moment 0 map (Fig. 3; panel c), but some CH<sub>3</sub>OH emission is also present at the SMM2, IRS7A, and IRS7B source positions although approximately a factor of 4 fainter compared to the peak position.

The CH<sub>3</sub>OH column densities towards the Coronet cluster members were determined from the integrated line intensities of the feathered SMA + APEX maps for all five CH<sub>3</sub>OH transitions belonging to the  $J = 5_K - 4_K$  branch (Fig. A.2 and Table B.1). The spectra are shown in Fig. B.1. For column density calculations, we followed two different treatments: a local thermodynamic equilibrium (LTE) analysis, assuming optically thin emission (Goldsmith & Langer 1999), and a non-LTE study employing the RADEX code (van der Tak et al. 2007). For both methods, a temperature of 30 K is assumed (Lindberg & Jørgensen 2012). Although the non-LTE analysis indicates that the CH<sub>3</sub>OH lines are not thermalised (see Fig. B.3) for densities below  $10^7$  cm<sup>-3</sup>, and four out of five transitions are optically thick (see Fig. B.4), the estimated CH<sub>3</sub>OH column densities derived with both methods agree with each other (for typical envelope densities  $n_{H_2}$  between  $10^5$  and  $10^6$  cm<sup>-3</sup>; see Table B.2). More details about the LTE and non-LTE analysis are presented in Appendix B. Table 3 lists the resulting column densities calculated from the LTE analysis and using  $T_{rot} = 30$  K. The uncertainties were calculated from the spectral rms noise and 20% calibration uncertainty.

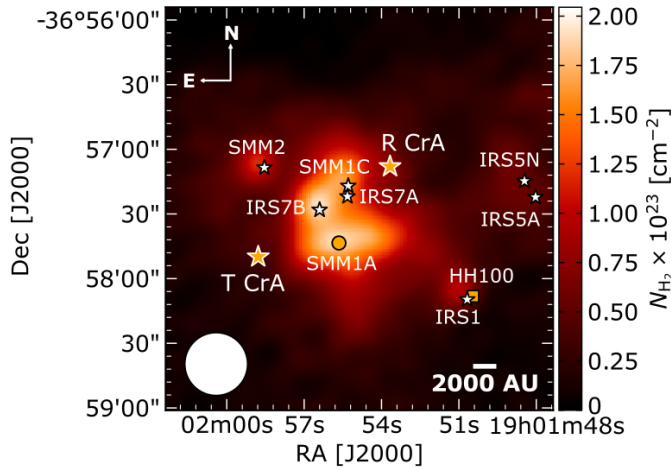
## 4. Linking methanol ice and gas

In this section, we analyse the gas-ice variations in the Coronet cluster. The H<sub>2</sub>O and CH<sub>3</sub>OH ice column densities were determined by Boogert et al. (2008) from *Spitzer* IRS mid-infrared spectra as part of the *Spitzer* Legacy Program ‘From Molecular Cores to Planet-Forming Disks’ (c2d). For the gas-phase counterpart, we made use of CH<sub>3</sub>OH gas column densities calculated in Sect. 3 and reported in Table 3.

### 4.1. H<sub>2</sub> column densities

Prior to searching for gas-ice correlations, the physical structure of the Coronet cluster is investigated by producing an H<sub>2</sub> column density map of the targeted region. This step is required as fractional abundances instead of absolute column densities need to be compared when studying gas-ice variations. This is due to the fact that gas and ice observations, by their very nature, are tracing different spatial scales, and therefore, a comparison between absolute values can lead to misinterpretations (e.g. Noble et al. 2017). At the same time we note that the H<sub>2</sub> column density calculated below provides the total H<sub>2</sub> column density – and hence it traces regions where methanol is in the gas phase as well in the ices. As a result, the CH<sub>3</sub>OH abundances presented in this work have to be interpreted as average abundances.

Figure 4 shows the H<sub>2</sub> column density map of the Coronet cluster obtained by making use of SCUBA sub-millimetre continuum maps at 850  $\mu$ m from Nutter et al. (2005). A detailed description of the formalism adopted to obtain the H<sub>2</sub> column



**Fig. 4.**  $\text{H}_2$  column density map of the R CrA region calculated from SCUBA dust emission maps at  $850\ \mu\text{m}$  by [Nutter et al. \(2005\)](#). Refer to Fig. 1 for a guide to the symbols of the Coronet cluster objects.

density from SCUBA maps is provided in Appendix C of [Perotti et al. \(2020\)](#). Briefly, in an optically thin thermal dust emission regime, the strength of the sub-millimetre radiation depends on the column density ( $N$ ), the dust temperature ( $T$ ) and the opacity ( $\kappa_\nu$ ), ([Kauffmann et al. 2008](#)). The inserted value for the opacity per unit dust+gas mass is  $0.0182\ \text{cm}^2\ \text{g}^{-1}$  (‘OH5 dust’; [Ossenkopf & Henning 1994](#)). The adopted value for the dust temperature is 30 K, based on previous observational studies of the Coronet cluster and radiative transfer models by [Lindberg & Jørgensen \(2012\)](#).

The youngest cluster members are all situated in the densest areas of the targeted region (Fig. 4). The apex of the estimated  $\text{H}_2$  column density lies to the south-east of R CrA, and in particular at the IRS7B, SMM1A, SMM1C and IRS7A positions ( $N_{\text{H}_2} > 1.45 \times 10^{23}\ \text{cm}^{-2}$ ). SMM2, IRS1, IRS5N and IRS5A are located in slightly less dense regions ( $0.5 \times 10^{23}\ \text{cm}^{-2} < N_{\text{H}_2} < 1.25 \times 10^{23}\ \text{cm}^{-2}$ ). The measured  $\text{H}_2$  column densities and their uncertainties for individual sources are listed in Table 3.

#### 4.2. Gas-ice variations

The analysis of gas-ice variations in the Coronet cluster is addressed by comparing fractional abundances ( $X$ ) of gas and ice species relative to  $\text{H}_2$  (Table 3). In addition, this section also provides an overview of gas-ice variations in Coronet with respect to two other nearby low-mass star-forming clusters: Serpens SVS 4 and Orion B35A, which have been observed with the same facilities, angular resolution, and receiver settings adopted for observations of the Coronet cluster to assure a meaningful comparison and reduce observational bias (see [Perotti et al. 2020, 2021](#)).

Figure 5 displays  $\text{CH}_3\text{OH}$  gas abundances as function of  $\text{H}_2\text{O}$  (a) and  $\text{CH}_3\text{OH}$  ice (b). A correlation between  $\text{CH}_3\text{OH}$  gas and  $\text{H}_2\text{O}$  ice for the three star-forming regions is seen in panel a with the Coronet cluster showing gas-ice abundances as low as the Orion B35A cloud. The lower methanol gas abundances in the Coronet cluster and in Orion B35A may be due to reduced ice mantle formation and enhanced photodissociation of methanol molecules upon desorption in these two regions with stronger UV field (see Sect. 5 for more details). The spread among the  $\text{H}_2\text{O}$  abundances for the Coronet cluster and Orion B35A is

limited, spanning less than a factor of 2.6 including the uncertainties. The same applies to the  $\text{CH}_3\text{OH}$  gas abundances. In contrast, Serpens SVS 4 exhibits the highest values compared to the other two regions and the largest spread, with  $\text{CH}_3\text{OH}$  gas abundances ranging up to  $\sim 20 \times 10^{-8}$ . One outlier is identified, Serpens SVS 4–12, which is characterised by the lowest  $\text{CH}_3\text{OH}$  gas and highest  $\text{H}_2\text{O}$  ice abundances relative to the other Serpens SVS 4 protostars. This result is not surprising since this is the most extinguished object of all targeted sources ( $A_V \sim 95\ \text{mag}$ ; [Pontoppidan et al. 2004](#)) with the lowest S/N (see [Perotti et al. 2020](#) for a detailed discussion on Serpens SVS 4–12).

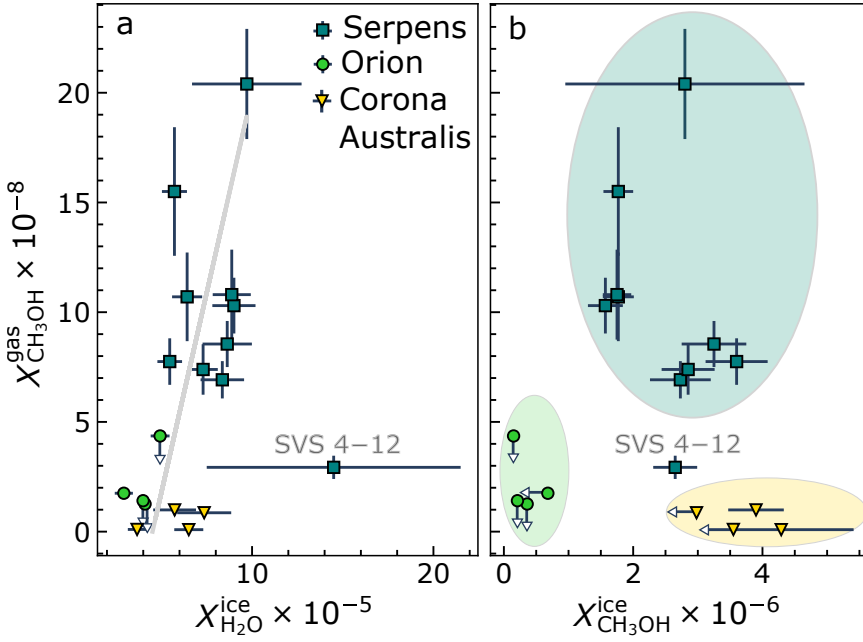
Figure 5b compares  $\text{CH}_3\text{OH}$  gas and ice abundances. In contrast to panel a, no clear trend is seen. However, when Serpens SVS 4–12 is excluded from the analysis it is possible to identify three different groups corresponding to the three regions: (i) Orion B35A with the lowest measurements of both  $\text{CH}_3\text{OH}$  in the gas and in the solid state; (ii) Serpens SVS 4 with intermediate to high values of  $\text{CH}_3\text{OH}$  ice and the highest gas abundances, and finally (iii) the Coronet cluster, with some of the highest  $\text{CH}_3\text{OH}$  ice but surprisingly low  $\text{CH}_3\text{OH}$  gas abundances. The observed behaviour for the Coronet cluster is likely due to  $\text{CH}_3\text{OH}$  destruction in the gas phase as a consequence of strong external irradiation ([Lindberg & Jørgensen 2012](#)).

Figure 6 shows  $\text{H}_2\text{O}$  (a) and  $\text{CH}_3\text{OH}$  (b) ice column densities ( $N$ ) as function of cloud visual extinction ( $A_V$ ) for lines of sights in the three nearby star-forming clusters. The cloud extinction estimates were taken from [Pontoppidan et al. \(2004\)](#) for Serpens SVS 4, from [Perotti et al. \(2021\)](#) for Orion B35A, and from [Alves et al. \(2014\)](#) for the Coronet cluster, then converted to visual extinction using the conversion factors (3.6 for  $A_J$ , 5.55 for  $A_H$  and 8.33 for  $A_K$ ) from [Weingartner & Draine \(2001\)](#) for  $RV = 5.5$  for the dense ISM.

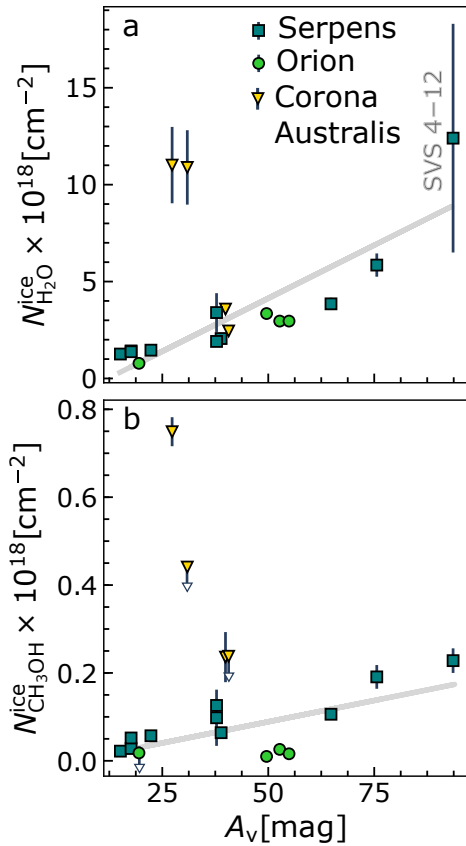
The distribution of data points in panels a and b shows a positive trend for Serpens SVS 4 and Orion B35A, indicating that to higher visual extinctions correspond higher columns of  $\text{H}_2\text{O}$  (a) and  $\text{CH}_3\text{OH}$  (b). This observation is particularly valid for  $\text{H}_2\text{O}$  (a) but it also applies to  $\text{CH}_3\text{OH}$  when an evaluation is done per star-forming region. The opposite behaviour is found for the Coronet cluster members, for both ice species. This negative trend could simply reflect inaccurate extinction values due to the lack of background stars in the densest area of the Corona Australis complex, where the Coronet cluster members are located (see Sect. 2 of [Alves et al. 2014](#) for more details).

## 5. Discussion

In this study, the  $\text{CH}_3\text{OH}$  chemical behaviour has been investigated with millimetric and infrared facilities. This provides direct observational constraints on the  $\text{CH}_3\text{OH}$  gas-to-ice ratio in protostellar envelopes and on its dependency on the physical conditions of star-forming regions. Laboratory experiments predict that  $\text{CH}_3\text{OH}$  ice in cold dark clouds is non-thermally desorbed with an efficiency spanning over approximately three orders of magnitude ( $10^{-6}$ – $10^{-3}$  molecules/photon; [Öberg et al. 2009b](#); [Bertin et al. 2016](#); [Cruz-Diaz et al. 2016](#); [Martín-Doménech et al. 2016](#)). This range is expected to increase up to  $\sim 10^{-2}$  molecules/photon ([Basalgète et al. 2021a,b](#)), as X-rays emitted from the central YSO strongly impact the abundances of  $\text{CH}_3\text{OH}$  in protostellar envelopes ([Notsu et al. 2021](#)). Generally, the majority of non-thermally desorbed  $\text{CH}_3\text{OH}$  molecules fragment during the desorption and a fraction of those recombine (e.g. [Bertin et al. 2016](#)). The  $\text{CH}_3\text{OH}$  non-thermal desorption efficiency and by inference the gas-to-ice ratio is highly dependent on the ice structure and composition



**Fig. 5.** Relationship between  $\text{CH}_3\text{OH}$  and  $\text{H}_2\text{O}$  fractional abundances ( $X$ ) towards three nearby low-mass star-forming regions. The triangles represent the gas and ice fractional abundances relative to  $\text{H}_2$  for the Coronet cluster members listed in Table 3 for which both gas and ice measurements are available. The squares and circles symbolise the values obtained for Serpens SVS 4 and Orion B35A from Perotti et al. (2020, 2021), respectively. The  $\text{CH}_3\text{OH}$  gas abundances are compared to the  $\text{H}_2\text{O}$  (panel a) and  $\text{CH}_3\text{OH}$  (panel b) ice abundances. Upper limits are marked as empty arrows. The grey line in panel a represents the linear fit to the detections. In panel b three different groups are identified, excluding Serpens SVS 4–12; they correspond to the three star-forming regions and they are marked with coloured ovals.



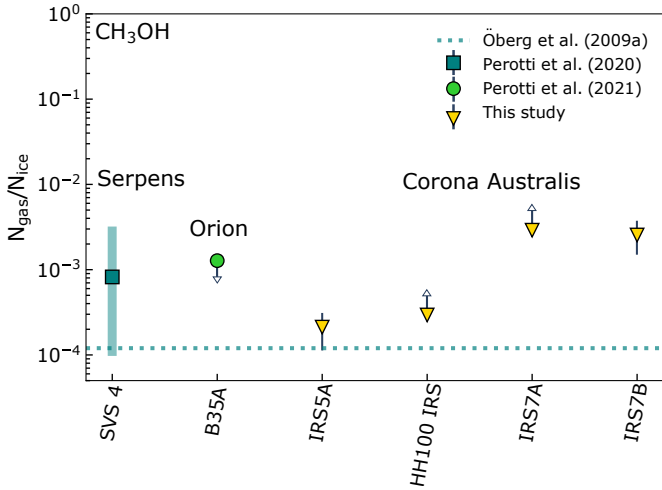
**Fig. 6.** Relationship between  $\text{H}_2\text{O}$  (a),  $\text{CH}_3\text{OH}$  (b) ice column densities ( $N$ ) and visual extinction ( $A_V$ ) for lines of sight in nearby star-forming clusters: Corona Australis Coronet (triangles), Serpens SVS 4 (squares), and Orion B35A (circles). Upper limits are marked as arrows. The grey lines represent linear fits to the detections. We note that the visual extinction values obtained for CrA have to be taken with care as their determination is hampered by the lack of background stars towards the densest region of the Coronet cluster (see Sect. 4.2). The data for Serpens SVS 4 and Orion B35A were taken from Perotti et al. (2020, 2021), respectively.

(pure  $\text{CH}_3\text{OH}$  versus  $\text{CH}_3\text{OH}$  mixed with  $\text{CO}$  molecules), as well as on the temperature, photon energy and flux (e.g. Öberg 2016; Carrascosa et al. 2023). To support laboratory experiments and computations, in this paper we provide observational measurements of the  $\text{CH}_3\text{OH}$  gas-to-ice ratio in the outer regions of protostellar envelopes.

Our targeted regions (Serpens SVS 4, Orion B35A and Corona Australis Coronet) share a number of similarities: they are low-mass star-forming regions, they are clustered and finally, they are affected by the presence of outflows and/or Herbig-Haro objects. However, they show distinct physical conditions for instance, the Serpens SVS 4 cluster is influenced by the presence of the Class 0 binary SMM4 (Pontoppidan et al. 2004), whereas B35A is affected by the nearby high-mass star  $\lambda$  Orionis (Reipurth & Friberg 2021) and the Coronet is strongly irradiated by the Herbig Ae/Be star R CrA (Lindberg & Jørgensen 2012). B35A is exposed to an interstellar radiation field  $\chi_{\text{ISRF}}$  of 34 (Wolfire et al. 1989)<sup>3</sup> enhanced by the neighbouring  $\lambda$  Orionis (Dolan & Mathieu 2002), whereas the radiation field in the Coronet cluster has been estimated to approximately  $\chi_{\text{ISRF}} \sim 750$  to account for the high fluxes at millimetre wavelengths (Lindberg & Jørgensen 2012).

In addition, the molecular clouds in which our targeted regions are located do not share a common formation history. The ongoing low-mass star formation in Serpens Main might have been triggered by cloud-cloud collisions (Duarte-Cabral et al. 2010) or compression from a shock wave from a supernova. However, there is no clear evidence that a nearby supernova ever occurred (Herczeg et al. 2019). In contrast to Serpens, the formation of the ring constituting the  $\lambda$  Orionis region is attributed to a supernova explosion that occurred roughly 1–6 Myr ago (Dolan & Mathieu 1999, 2002; Kounkel 2020). The low-mass star formation in B35A was likely generated by the presence of neighbouring massive stars and their stellar winds (Barrado et al. 2018). Finally, star formation in CrA has been supposedly promoted by a high-velocity cloud impact onto the Galactic

<sup>3</sup> We note that Wolfire et al. (1989) report a  $\log(G_0)$  of 1.3 which is here converted to  $\chi_{\text{ISRF}}$  by applying the formalism described by Draine & Bertoldi (1996).



**Fig. 7.**  $\text{CH}_3\text{OH}$  gas-to-ice ratios ( $N_{\text{gas}}/N_{\text{ice}}$ ) towards low-mass protostars in the Coronet cluster in Corona Australis (triangles). The square and the circle symbolise the average ratios for ten sources in Serpens SVS 4 and four sources in Orion B35A, respectively. The green shaded area corresponds to the interval of gas-to-ice ratios observed for the ten sources in the Serpens SVS 4 cluster. The dotted line indicates the average value for four sources from Öberg et al. (2009a). Upper and lower limits are marked as arrows.

plane (Neuhäuser & Forbrich 2008) or by the expansion of the UpperCenLupus (UCL) superbubble (Mamajek et al. 2002). The fact that the three regions likely did not form in the same way implies different initial physical conditions for the production of  $\text{CH}_3\text{OH}$ .

Figure 7 illustrates the distribution of the  $\text{CH}_3\text{OH}$  gas-to-ice ratios ( $N_{\text{gas}}/N_{\text{ice}}$ ) towards low-mass protostars located in the different molecular clouds described above. The average  $\text{CH}_3\text{OH}$  gas-to-ice ratio ( $1.2 \times 10^{-4}$ ) determined from millimetric (single-dish) and infrared measurements by Öberg et al. (2009a) for four Class 0/I sources in Perseus, Taurus and Serpens is over-plotted. A detailed comparison between the gas-to-ice ratio determined by Öberg et al. (2009a) and the ratios calculated for the sources in the Serpens SVS 4 and the Orion B35A cloud is provided in Perotti et al. (2021). In this section we focus on the ratios obtained for CrA and their comparison with the values determined for the two other star-forming regions.

The distribution of  $\text{CH}_3\text{OH}$  gas-to-ice ratios towards the Coronet cluster covers approximately one order of magnitude. The uncertainty on the determination of the  $\text{CH}_3\text{OH}$  ice column densities for HH100 IRS1 and IRS7A (Boogert et al. 2008) results in lower limits for their gas-to-ice ratios, consequently we cannot provide solid comparisons for these two sources. In contrast, the values obtained for IRS5A and IRS7B fall in the range of gas-to-ice ratios previously obtained for the protostars located in the Serpens SVS 4 cluster.

The  $\text{CH}_3\text{OH}$  gas-to-ice ratios constrained from gas and ice observations of a sample of protostellar envelopes (Fig. 7) do not point to one value but, instead, to a distribution (from  $1.2 \times 10^{-4}$  to  $3.1 \times 10^{-3}$ ) which validates previous observational studies (e.g. Öberg et al. 2009a; Perotti et al. 2020). This distribution agrees with a more complex scenario for the  $\text{CH}_3\text{OH}$  desorption elucidated in the laboratory in which a large fraction of  $\text{CH}_3\text{OH}$  molecules do not desorb intact, but instead fragment and eventually recombine during and after the desorption (e.g. Bertin et al. 2016) leading to  $\text{CH}_3\text{OH}$  gas-to-ice ratios lower than  $10^{-3}$ .

A second interesting observation is that the  $\text{CH}_3\text{OH}$  gas-to-ice ratios fall in a fairly narrow range; this suggests that the  $\text{CH}_3\text{OH}$  chemistry at play in the cold outer regions of protostellar envelopes belonging to different low-mass star-forming regions is relatively independent of variations in the physical conditions. We note that a similar conclusion could not be tested for the innermost envelope regions where thermal desorption dominates and young discs are present; our observations are simply not sensitive to these regions. Our analysis implicitly indicates that the  $\text{CH}_3\text{OH}$  chemistry in the outer envelope regions is set to a large degree already during the prestellar stage. A related result came to light when comparing abundances of complex organic molecules observed in distinct star-forming environments Scibelli & Shirley (2020), (e.g. Galactic disc versus Galactic centre, high-mass versus low-mass star-forming region; Jørgensen et al. 2020; Coletta et al. 2020; Yang et al. 2021; Nazari et al. 2022a). A larger sample is however necessary to prove this statement, and additional observations, laboratory experiments and modelling efforts are required to elucidate further the impact of the physical evolution of protostars on their chemical signatures and vice versa.

## 6. Conclusions

We present 1.3 mm SMA and APEX observations towards the Coronet cluster in Corona Australis, a unique astrochemical laboratory to study the effect of external irradiation on the chemical and physical evolution of young protostars. In addition, we determined  $\text{CH}_3\text{OH}$  gas-phase abundances and compared those with  $\text{CH}_3\text{OH}$  ice abundances from Boogert et al. (2008) to directly constrain the  $\text{CH}_3\text{OH}$  gas-to-ice ratio in cold protostellar envelopes. Our key findings are:

- We find a positive trend between  $\text{H}_2\text{O}$  ice and  $\text{CH}_3\text{OH}$  gas abundances in Serpens SVS 4, Coronet, and Orion B35A with the latter two regions showing the lowest gas abundances. This result is attributed to reduced ice mantle formation and substantial destruction of methanol gas molecules in these two regions characterised by stronger external UV field;
- The distribution of  $\text{CH}_3\text{OH}$  gas-to-ice ratios in CrA constrained from millimetric and infrared observations spans approximately one order of magnitude (from  $1.2 \times 10^{-4}$  to  $3.1 \times 10^{-3}$ ) and reinforces previous observational constraints on the  $\text{CH}_3\text{OH}$  gas-to-ice ratio;
- Similarities are found between  $\text{CH}_3\text{OH}$  gas-to-ice ratios determined in different low-mass star forming regions (Serpens SVS 4, Orion B35A and Corona Australis Coronet), characterised by distinct physical conditions and formation histories. This result suggests that the  $\text{CH}_3\text{OH}$  chemistry – in the outer regions of low-mass envelopes – is relatively independent of variations in the physical conditions and thus it is set to a large degree during the prestellar stage.

Multiple avenues for future studies can be pursued. Comparisons of gas and ice abundances of key species are currently limited by the low number of infrared surveys. Current infrared facilities such as the *James Webb* Space Telescope are now routinely probing interstellar ices towards background stars (McClure et al. 2023), pre-stellar cores, proto-stellar envelopes (Yang et al. 2022) and near edge-on discs (Sturm et al. 2023), thus significantly increasing the number of studied regions. Therefore, it will be feasible to search for gas-ice trends found in other regions, improve our constraints on the  $\text{CH}_3\text{OH}$  gas-to-ice ratios, and ultimately determine ratios for a larger set of major interstellar molecules. This approach will provide important feedback on

the interactions between ice and gas material during its journey from the molecular cloud to the disc and on the impact of the physical conditions on the physico-chemical evolution of protostars. In this context, the Corona Australis complex still has much to teach us, hosting one of the youngest population of protostars observed so far.

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## Appendix A: Overview of molecular line detections

This section reports the primary beam corrected integrated intensity (moment zero) maps of the species detected in the SMA dataset (Figures A.1–A.3). For our molecule of interest, CH<sub>3</sub>OH, we also include in A.2 the combined SMA+APEX integrated intensity maps. Please refer to Appendix B.1 of Perotti et al. (2020) for a detailed description of the feathering procedure used to combine interferometric and single-dish data.

The eighteen molecular transitions of seven species identified in the SMA datasets are listed in Table A.1. The detected species are: carbon monoxide (<sup>12</sup>CO, <sup>13</sup>CO and C<sup>18</sup>O), formaldehyde (H<sub>2</sub>CO), methanol (CH<sub>3</sub>OH), sulphur oxide (SO), sulphur dioxide (SO<sub>2</sub>), cyclopropenylidene (c-C<sub>3</sub>H<sub>2</sub>), and cyanoacetylene (HC<sub>3</sub>N). The line emission is predominantly extended and it does not necessarily coincide with the position of the compact continuum sources in the Coronet cluster (Figures A.1–A.3). In contrast, it is mostly localised to the east of the Herbig Ae/Be star R CrA and to the south of the young stellar objects IRS7A and IRS7B, at the pre-stellar core SMM1A position. This statement applies especially to H<sub>2</sub>CO and CH<sub>3</sub>OH, and is consistent with previous 1.3 mm SMA observations of the region by Lindberg & Jørgensen (2012) covering the 216.849–218.831 GHz and 226.849–228.831 GHz frequencies ranges. Interestingly, no molecular line emission was detected towards SMM1A in the SMA dataset presented by Chen & Arce (2010), although the same number of antennas and configuration were used and the SMA correlator covered spectral ranges similar to the ones presented in this work. However, the SMA observations by Chen & Arce (2010) were designed to study the dust continuum and they were characterised by a shorter bandwidth and likely lower line sensitivity. Unfortunately, a comparison of the line sensitivity of the two datasets could not be made as this information is not reported in Chen & Arce (2010) and a reduced, continuum-subtracted spectrum is not provided either.

The detected species in the SMA data can be roughly divided into three groups described based on their emission pattern (Table A.1): (I) CO species, (II) S-species, and (III) carbon-chain species. Overall, the emission reported in the SMA dataset is dominated by molecular tracers such as H<sub>2</sub>CO, CH<sub>3</sub>OH, SO, SO<sub>2</sub>. The emission of H<sub>2</sub>CO and CH<sub>3</sub>OH is consistent with temperature variations in the field of view of the observations, and particularly with the external irradiation from the variable stars in the region (R CrA and T CrA). In contrast, the shock tracers (SO, SO<sub>2</sub>) probe the numerous outflows and Herbig-Haro objects in the region, which are indicators of a young stellar population in the earliest evolutionary stages.

### Appendix A.1. CO species

Figure A.1 displays the SMA primary beam corrected integrated intensity maps for <sup>12</sup>CO, <sup>13</sup>CO, C<sup>18</sup>O and H<sub>2</sub>CO. The distribution of the CO-species emission is extended, especially for <sup>12</sup>CO and its isotopologues. The three transitions of H<sub>2</sub>CO show the same emission pattern (Figure A.1, panels d–f). The emission is localised in two ridges: the northern ridge between R CrA and SMM1C, and the southern ridge to the south of IRS7B, in the vicinity of SMM1A. No bright line emission is observed towards the young stellar objects SMM1C, IRS7A and IRS7B. The emission offset from the YSO positions and the elongated shape of the ridges have been previously identified in this region and interpreted as indicators of external heating (Lindberg & Jørgensen 2012). The emission localised in the southern ridge traces the three condensations within the pre-stellar core

SMM1A (SMM1A-a, SMM1A-b, and SMM1A-c), analysed in dust continuum observations by Chen & Arce (2010).

The strength of the H<sub>2</sub>CO emission, for the three detected H<sub>2</sub>CO lines is comparable in both ridges. A similar morphology pattern is observed for CH<sub>3</sub>OH where the emission peak is located within the SMM1A core. However, neither the CH<sub>3</sub>OH  $J=4_2 - 3_1$  line nor the CH<sub>3</sub>OH transitions belonging to the  $J = 5_K - 4_K$  branch show a northern ridge (Figures 3 and A.2).

### Appendix A.2. S species

The SO and SO<sub>2</sub> line emissions appear relatively bright and compact (Figure A.3, panels m–o). The distribution of the emission is identical for the two SO transitions, and it peaks towards IRS7A and SMM1C. The strength of the emission among the SO transitions is very similar (7–8 Jy beam<sup>-1</sup> km s<sup>-1</sup>; Figure A.3, panels m–n). The SO<sub>2</sub> emission is concentrated towards IRS7A and is approximately a factor of 2.5 fainter compared to SO (Figure A.3, panels o). In contrast to the majority of the CO-species, the S-species emission have their strongest emission associated with the YSO positions. SO and SO<sub>2</sub> are molecular tracers of energetic inputs in the form of outflows and jets. The luminous SO and SO<sub>2</sub> emissions appear to indicate that higher velocity (~5–10 km s<sup>-1</sup>) flows of matter are present in the field-of-view of the observations, associated with IRS7A and SMM1C. This result is in agreement with studies of the Herbig-Haro objects in the Coronet cluster attributed to SMM1C (Wang et al. 2004; Peterson et al. 2011).

### Appendix A.3. Carbon-chain species

The third group of lines corresponds to carbon-chain species such as c-C<sub>3</sub>H<sub>2</sub> and HC<sub>3</sub>N which show a fairly weak emission (1.6–3.0 Jy beam<sup>-1</sup> km s<sup>-1</sup>) compared to the previous two groups (CO- and S-species). Their emission is diffuse and characterised by a fairly low S/N (Figure A.3, panels p–r). Similarly to <sup>13</sup>CO and C<sup>18</sup>O, the carbon-chain species have their brightest emission associated with IRS7B.

Table A.1: Identified molecular transitions in the SMA dataset.

Species	Transition	Frequency <sup>(a)</sup> [GHz]	$A_{ul}^{(a)}$ [s <sup>-1</sup> ]	$E_u^{(a)}$ [K]
CO-species				
<sup>12</sup> CO	2 – 1	230.538	$6.92 \times 10^{-7}$	17
C <sup>18</sup> O	2 – 1	219.560	$6.01 \times 10^{-7}$	16
<sup>13</sup> CO	2 – 1	220.398	$6.04 \times 10^{-7}$	16
H <sub>2</sub> CO	3 <sub>0,3</sub> – 2 <sub>0,2</sub>	218.222	$2.82 \times 10^{-4}$	21
H <sub>2</sub> CO	3 <sub>2,2</sub> – 2 <sub>2,1</sub>	218.476	$1.58 \times 10^{-4}$	68
H <sub>2</sub> CO	3 <sub>2,1</sub> – 2 <sub>2,0</sub>	218.760	$1.58 \times 10^{-4}$	68
CH <sub>3</sub> OH	4 <sub>2</sub> – 3 <sub>1</sub> , E	218.440	$4.68 \times 10^{-5}$	46
CH <sub>3</sub> OH <sup>b</sup>	5 <sub>+0</sub> – 4 <sub>+0</sub> E	241.700	$6.04 \times 10^{-5}$	47.9
CH <sub>3</sub> OH <sup>b</sup>	5 <sub>-1</sub> – 4 <sub>-1</sub> E	241.767	$5.81 \times 10^{-5}$	40.4
CH <sub>3</sub> OH <sup>b</sup>	5 <sub>0</sub> – 4 <sub>0</sub> A <sup>+</sup>	241.791	$6.05 \times 10^{-5}$	34.8
CH <sub>3</sub> OH <sup>b</sup>	5 <sub>+1</sub> – 4 <sub>+1</sub> E	241.879	$5.96 \times 10^{-5}$	55.9
CH <sub>3</sub> OH <sup>b</sup>	5 <sub>-2</sub> – 4 <sub>-2</sub> E	241.904	$5.09 \times 10^{-5}$	60.7
S-species				
SO	5 <sub>5</sub> – 4 <sub>4</sub>	215.220	$1.20 \times 10^{-4}$	44
SO	6 <sub>5</sub> – 5 <sub>4</sub>	219.949	$1.35 \times 10^{-4}$	35
SO <sub>2</sub>	4 <sub>2,2</sub> – 3 <sub>1,3</sub>	235.151	$7.69 \times 10^{-5}$	19
Carbon-chain species				
<i>c</i> -C <sub>3</sub> H <sub>2</sub>	6 <sub>0,6</sub> – 5 <sub>1,5</sub>	217.822	$5.37 \times 10^{-4}$	39
<i>c</i> -C <sub>3</sub> H <sub>2</sub>	5 <sub>1,4</sub> – 4 <sub>2,3</sub>	217.940	$3.98 \times 10^{-4}$	35
HC <sub>3</sub> N	24 – 23	218.324	$8.32 \times 10^{-4}$	131

**Notes.** <sup>(a)</sup> From the Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001, 2005; Endres et al. 2016) and the Jet Propulsion Laboratory catalogue (Pickett et al. 1998). The CH<sub>3</sub>OH transitions belonging to the CH<sub>3</sub>OH  $J_K = 5_K - 4_K$  Q-branch are also detected in the APEX data.

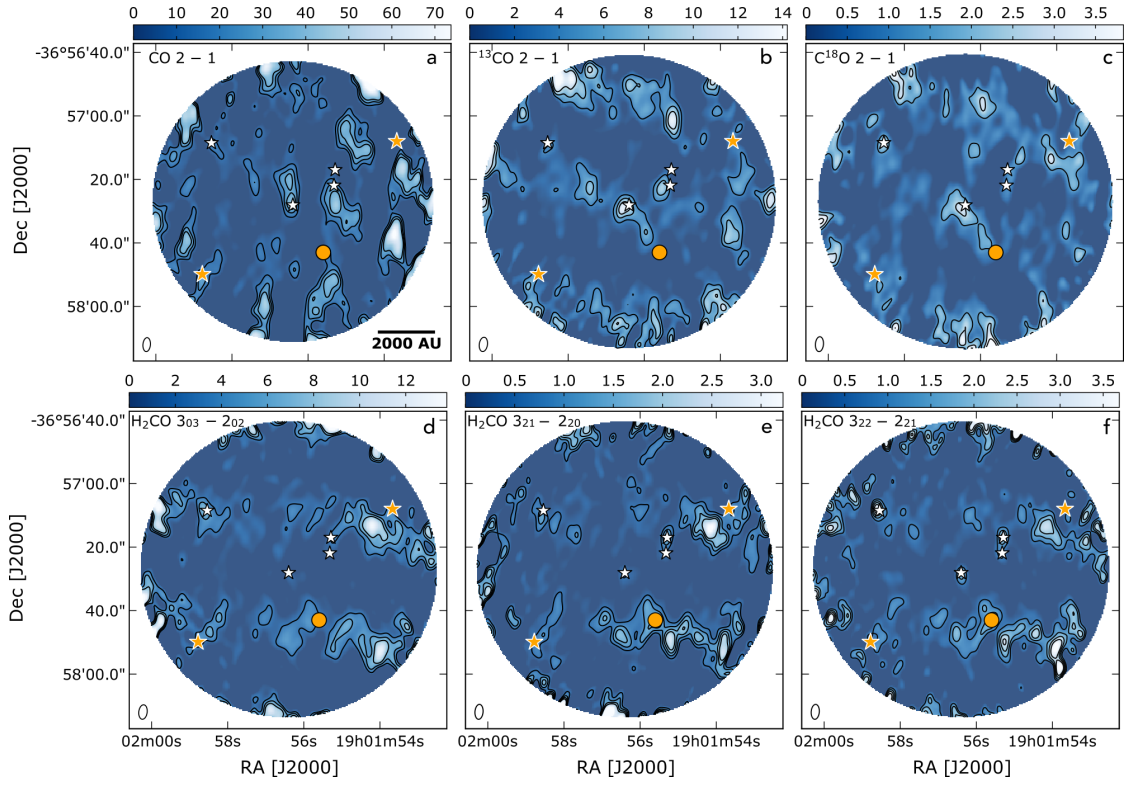


Fig. A.1: Primary beam corrected integrated intensity maps in  $\text{Jy beam}^{-1}\text{km s}^{-1}$  for the different species detected with the SMA. In each plot contours start at  $5\sigma$  and continue in intervals of  $5\sigma$ , except for panel a-c for which contours start at  $10\sigma$  and continue in intervals of  $10\sigma$ . Shown are (a) CO  $J = 2 - 1$  ( $\sigma = 0.96 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (b)  $^{13}\text{CO}$   $J = 2 - 1$  ( $\sigma = 0.32 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (c)  $\text{C}^{18}\text{O}$   $J = 2 - 1$  ( $\sigma = 0.13 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (d)  $\text{H}_2\text{CO}$   $3_{03} - 2_{02}$  ( $\sigma = 0.23 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ), (e)  $\text{H}_2\text{CO}$   $3_{21} - 2_{20}$  ( $\sigma = 88 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ), (f)  $\text{H}_2\text{CO}$   $3_{22} - 2_{21}$  ( $\sigma = 0.14 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ). The circular field-of-view corresponds to the SMA primary beam. The size of the synthesised beam is shown in the lower left corner of each image. The white and orange stars and the orange circle mark the position of the Coronet cluster members as in Fig. 1.

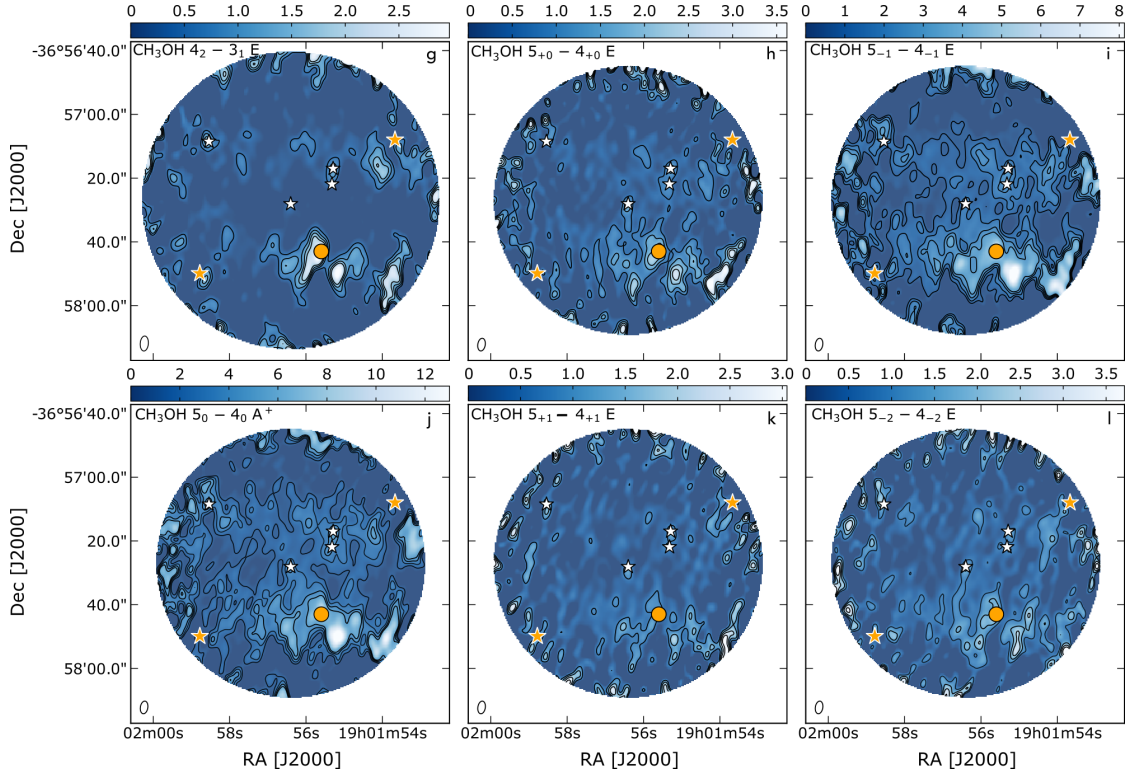


Fig. A.2: Same as Figure A.1, but for primary beam corrected integrated intensity maps in  $\text{Jy beam}^{-1}\text{km s}^{-1}$  for: (g)  $\text{CH}_3\text{OH } 4_2-3_1 \text{ E}$  ( $\sigma=88 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ), (h)  $\text{CH}_3\text{OH } 5_{+0}-4_{+0} \text{ E}$  ( $\sigma=0.12 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (i)  $\text{CH}_3\text{OH } 5_{-1}-4_{-1} \text{ E}$  ( $\sigma=0.13 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (j)  $\text{CH}_3\text{OH } 5_0-4_0 \text{ A}^+$  ( $\sigma=0.17 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (k)  $\text{CH}_3\text{OH } 5_{+1}-4_{+1} \text{ E}$  ( $\sigma=0.13 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), and (l)  $\text{CH}_3\text{OH } 5_{-2}-4_{-2} \text{ E}$  ( $\sigma=0.15 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ) detected by the SMA (panel g), and in the combined SMA and APEX data (panels h-l).

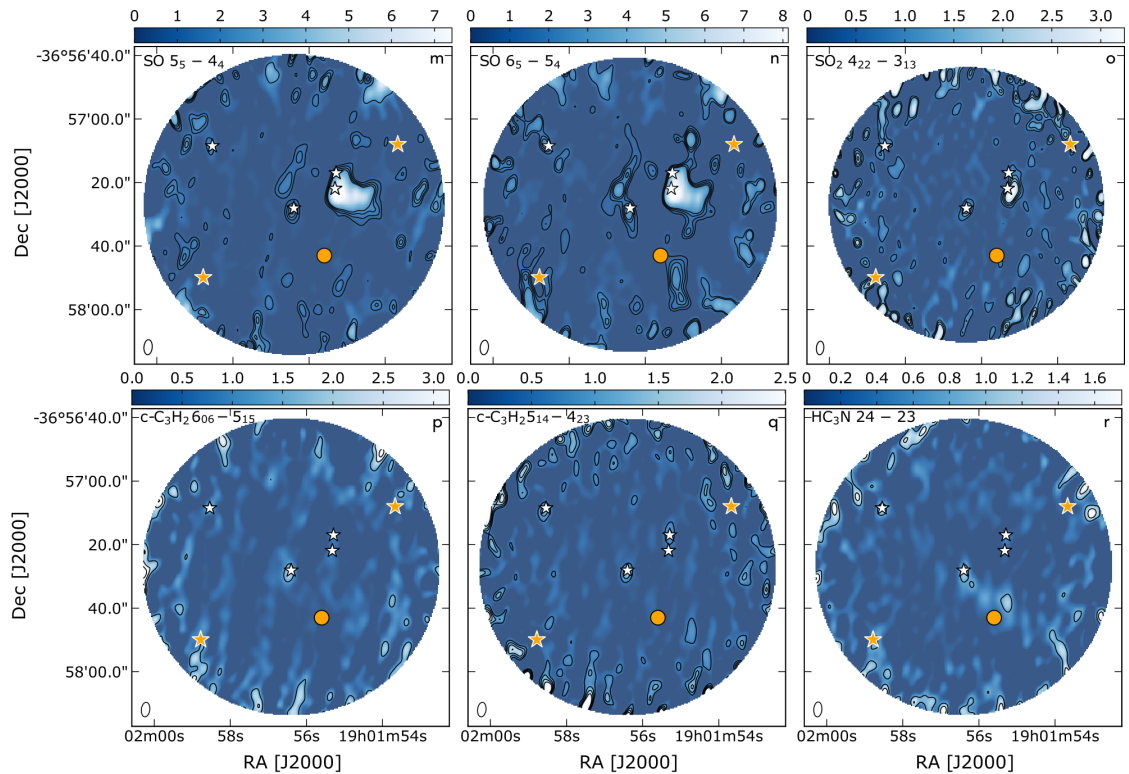


Fig. A.3: Same as Figure A.1, but for primary beam corrected integrated intensity maps in  $\text{Jy beam}^{-1}\text{km s}^{-1}$  for: (m)  $\text{SO } 5_5-4_4$  ( $\sigma=0.12 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (n)  $\text{SO } 6_5-5_4$  ( $\sigma=0.13 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (o)  $\text{SO}_2 4_{22}-3_{13}$  ( $\sigma=0.15 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (p)  $c\text{-C}_3\text{H}_2 6_{06}-5_{15}$  ( $\sigma=0.13 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), (q)  $c\text{-C}_3\text{H}_2 5_{14}-4_{23}$  ( $\sigma=0.11 \text{ Jy beam}^{-1}\text{km s}^{-1}$ ), and (r)  $\text{HC}_3\text{N } J=24-23$  ( $\sigma=71 \text{ mJy beam}^{-1}\text{km s}^{-1}$ ) detected by the SMA.

## Appendix B: Gas-phase methanol

### Appendix B.1. Spectra of methanol

Figure B.1 shows the spectra for the brightest transitions of the CH<sub>3</sub>OH  $J_K = 5_K - 4_K$  Q-branch,  $J=5_0 - 4_0$  A<sup>+</sup> and  $J=5_{-1} - 4_{-1}$  E, respectively, extracted from the combined SMA + APEX dataset. Table A.1 reports the spectral data for both CH<sub>3</sub>OH transitions. The CH<sub>3</sub>OH emission lines have a Gaussian shape and do not show blue or red-shifted components.

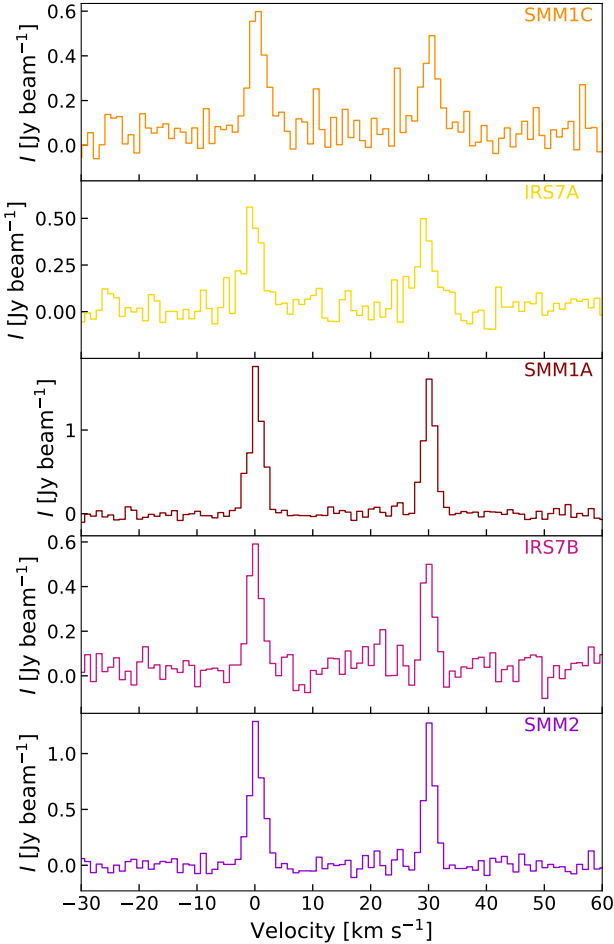


Fig. B.1: CH<sub>3</sub>OH  $J=5_0 - 4_0$  A<sup>+</sup> and CH<sub>3</sub>OH  $J=5_{-1} - 4_{-1}$  E spectra towards the Coronet cluster members detected in the combined SMA+APEX dataset.

### Appendix B.2. Derivation of CH<sub>3</sub>OH column densities

The column densities of gas-phase CH<sub>3</sub>OH for the Coronet cluster members were calculated using the integrated intensities of the combined interferometric (SMA) and single-dish (APEX) data reported in Table B.1. The SMA+APEX primary beam corrected integrated intensity (moment 0) maps for the CH<sub>3</sub>OH  $J_K = 5_K - 4_K$  Q-branch are displayed in Figure A.2.

For the column density calculation one of the two adopted procedures assumes LTE conditions, optically thin line emission, homogeneous source filling the telescope beam and a single excitation temperature to describe the level population. The equations used follow the prescription of Goldsmith & Langer (1999), but assuming a fixed rotational temperature equal to 30 K (Lindberg & Jørgensen 2012), are reported in Appendix A.4 of Perotti et al. (2021) and thus they will not be repeated here.

The values used for the column densities calculation are listed in Table 2 and have been compiled from the Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001, 2005; Endres et al. 2016) and the Jet Propulsion Laboratory catalogue (Pickett et al. 1998). In particular, the spectroscopic data are from Xu et al. (2008). The rotational diagrams for the Coronet cluster members are shown in Figure B.2. The CH<sub>3</sub>OH column densities and their uncertainties are reported in Table 3. Using a different rotational temperature (e.g. 20 K or 40 K) results in column densities being increased or decreased by approximately a factor of 0.5, respectively. The uncertainties on the gas column densities were calculated based on the spectral rms noise and on the 20% calibration uncertainty.

### Appendix B.3. Non-LTE analysis

Further analysis of the CH<sub>3</sub>OH emission was performed using the non-LTE code RADEX (van der Tak et al. 2007). In particular, all the five CH<sub>3</sub>OH transitions were used to compare the observed relative intensities with the RADEX results for a H<sub>2</sub> number density ( $n_{\text{H}_2}$ ) in the range  $10^3 - 10^7$  cm<sup>-3</sup> and a kinetic temperature equal to 30 K. Figure B.3 shows that all lines are in sub-thermal regime for  $n_{\text{H}_2}$  below  $10^7$  cm<sup>-3</sup>, typical for protostellar envelopes. Figure B.4 illustrates the optical depth of the five transitions for CH<sub>3</sub>OH column densities between  $10^{13} - 10^{16}$  cm<sup>-2</sup>. The three brightest transitions ( $J = 5_{+0} - 4_{+0}$ ,  $J = 5_{-1} - 4_{-1}$ ,  $J = 5_0 - 4_0$ ) are optically thick for CH<sub>3</sub>OH column densities equal to  $10^{15}$  cm<sup>-2</sup>, while the  $J = 5_{+1} - 4_{+1}$  transition is marginally optically thick and the weakest transition,  $J = 5_{-2} - 4_{-2}$ , can be considered as optically thin. Figures B.5–B.7 show a visual comparison between CH<sub>3</sub>OH column densities calculated using RADEX and obtained from the LTE analysis for each Coronet cluster member. Thin contours represent lower limits for the integrated line intensities of the optically thick lines from Table B.1, while thick contours indicate integrated line intensities for the less optically thick lines. The thickness of these contours are given by the errors of the integrated line intensities. A range of CH<sub>3</sub>OH column densities is estimated for each source from the weakest transition (CH<sub>3</sub>OH  $J = 5_{-2} - 4_{-2}$ ) and assuming  $n_{\text{H}_2}$  between  $10^5$  and  $10^6$  cm<sup>-3</sup> (see Table B.2). The vertical black line in Figures B.5–B.7 depicts the CH<sub>3</sub>OH column densities calculated with the LTE analysis listed in Table 3. We note that the column densities obtained with both methods are consistent for  $n_{\text{H}_2}$  equal to  $10^5 - 10^6$  cm<sup>-3</sup> (see Table B.2).

Table B.1: Integrated CH<sub>3</sub>OH line intensities in units of Jy beam<sup>-1</sup> km s<sup>-1</sup> towards the Coronet cluster members.

Source	$5_{+0} - 4_{+0} E^+$	$5_{-1} - 4_{-1} E$	$5_0 - 4_0 A^+$	$5_{+1} - 4_{+1} E$	$5_{-2} - 4_{-2} E$
SMM1C	1.41±0.29	2.48±0.50	3.86±0.77	1.24±0.25	1.01±0.21
IRS7A	0.79±0.17	2.26±0.45	2.94±0.59	0.61±0.13	0.62±0.14
SMM1A	2.43±0.49	6.44±1.29	7.23±1.45	1.15±0.24	2.57±0.52
IRS7B	1.74±0.35	1.93±0.39	2.97±0.60	0.96±0.20	1.57±0.32
SMM2	0.86±0.18	1.97±0.40	3.79±0.76	0.86±0.18	1.26±0.26

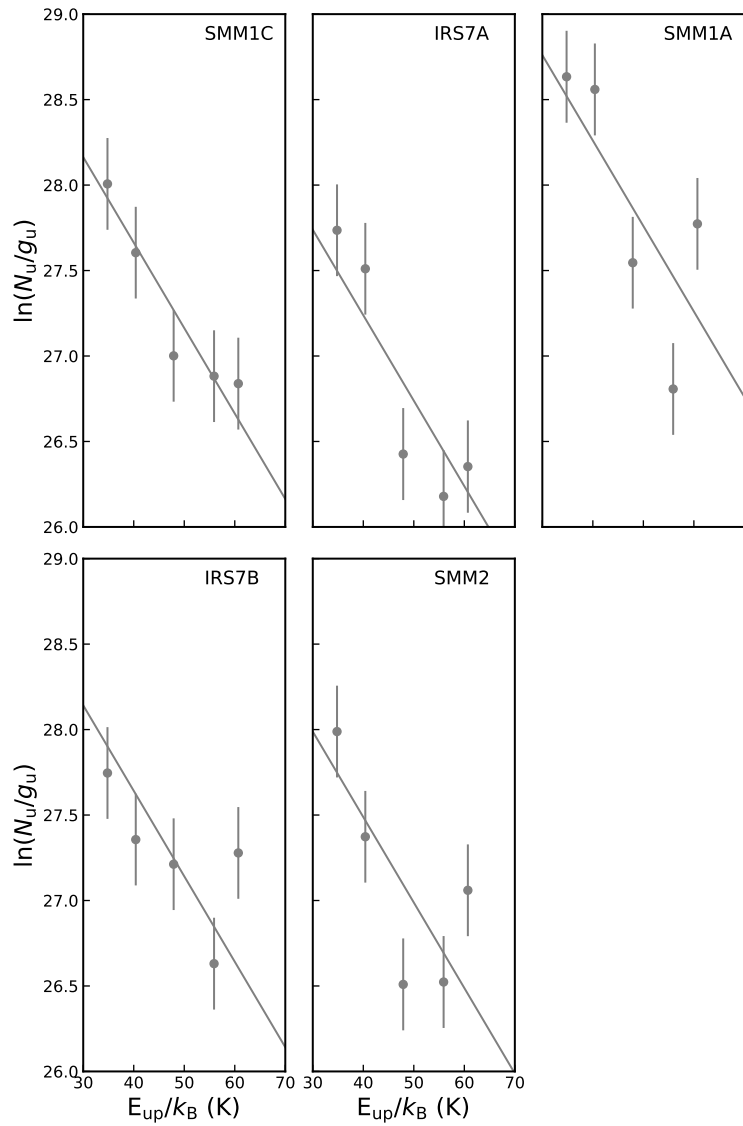
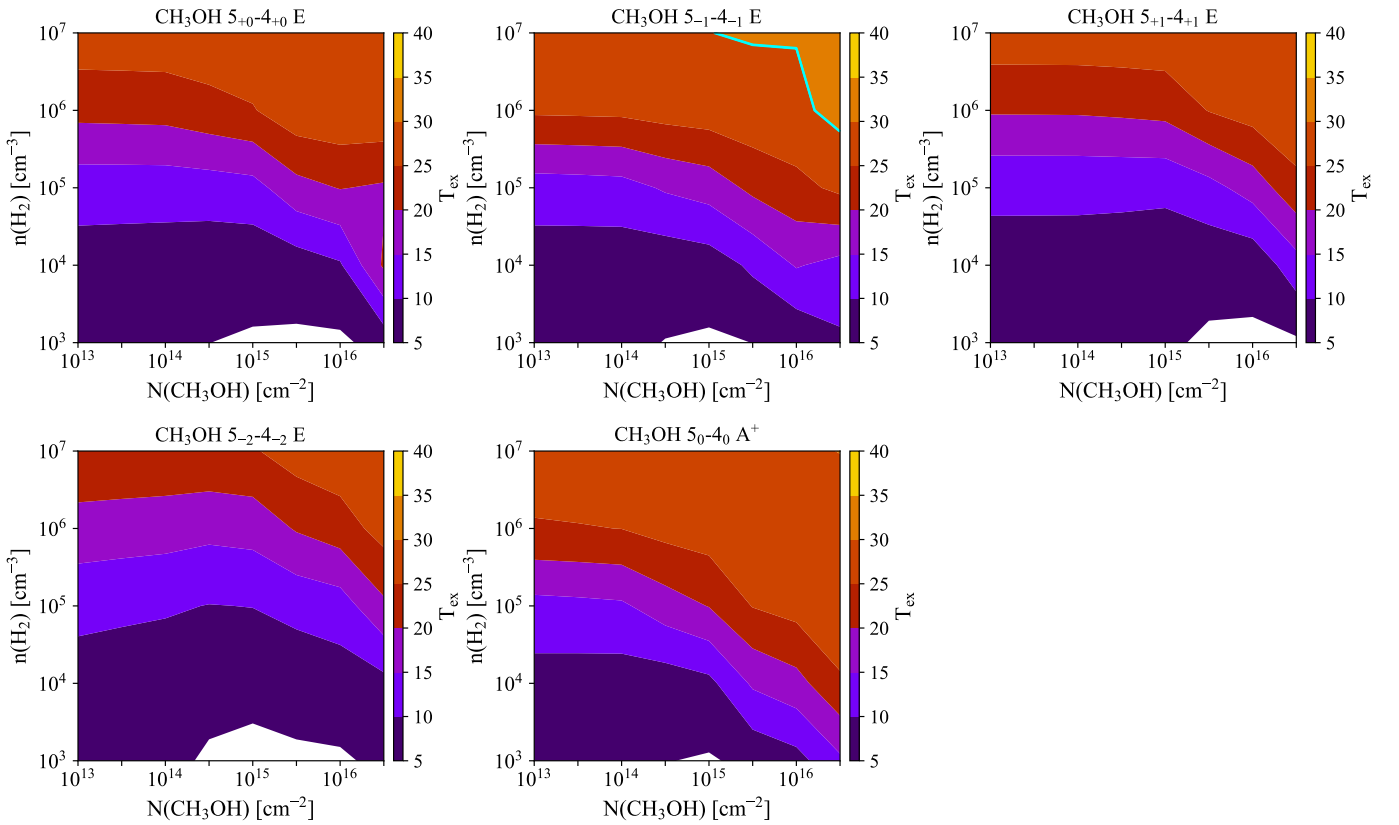


Fig. B.2: Rotational diagrams of CH<sub>3</sub>OH for the Coronet cluster members. The solid line shows the fixed slope for  $T_{\text{rot}} = 30$  K. Error bars are for  $1\sigma$  uncertainties. The column densities are reported in Table 3. Only transitions belonging to the  $J = 5_K - 4_K$  Q-branch have been considered.

Table B.2: Comparison between total CH<sub>3</sub>OH gas column densities ( $N$ ) calculated using the LTE and the non-LTE analyses towards the Coronet cluster members.

Object	$N_{\text{CH}_3\text{OH}}^{\text{gas, LTE (a)}}$ [ $10^{15} \text{ cm}^{-2}$ ]	$N_{\text{CH}_3\text{OH}}^{\text{gas, non-LTE (b)}}$ [ $10^{15} \text{ cm}^{-2}$ ]
SMM1C	$1.95 \pm 0.23$	1.03 – 5.50
IRS7A	$1.28 \pm 0.15$	0.94 – 3.50
SMM1A	$3.54 \pm 0.43$	1.50 – 10.0
IRS7B	$1.91 \pm 0.23$	1.07 – 7.30
SMM2	$1.64 \pm 0.20$	1.04 – 6.30

**Notes.** <sup>(a)</sup>CH<sub>3</sub>OH column densities calculated with the LTE analysis and fixing  $T_{\text{rot}} = 30 \text{ K}$ . <sup>(b)</sup> CH<sub>3</sub>OH column densities determined using the non-LTE code RADEX for  $T_{\text{kin}} = 30 \text{ K}$  and  $n_{\text{H}_2} = 10^5 - 10^6 \text{ cm}^{-3}$ .


 Fig. B.3: RADEX model results for the five CH<sub>3</sub>OH transitions (Table A.1) using  $T_{\text{kin}} = 30 \text{ K}$ . The cyan line indicates  $T_{\text{ex}} = 30 \text{ K}$ .

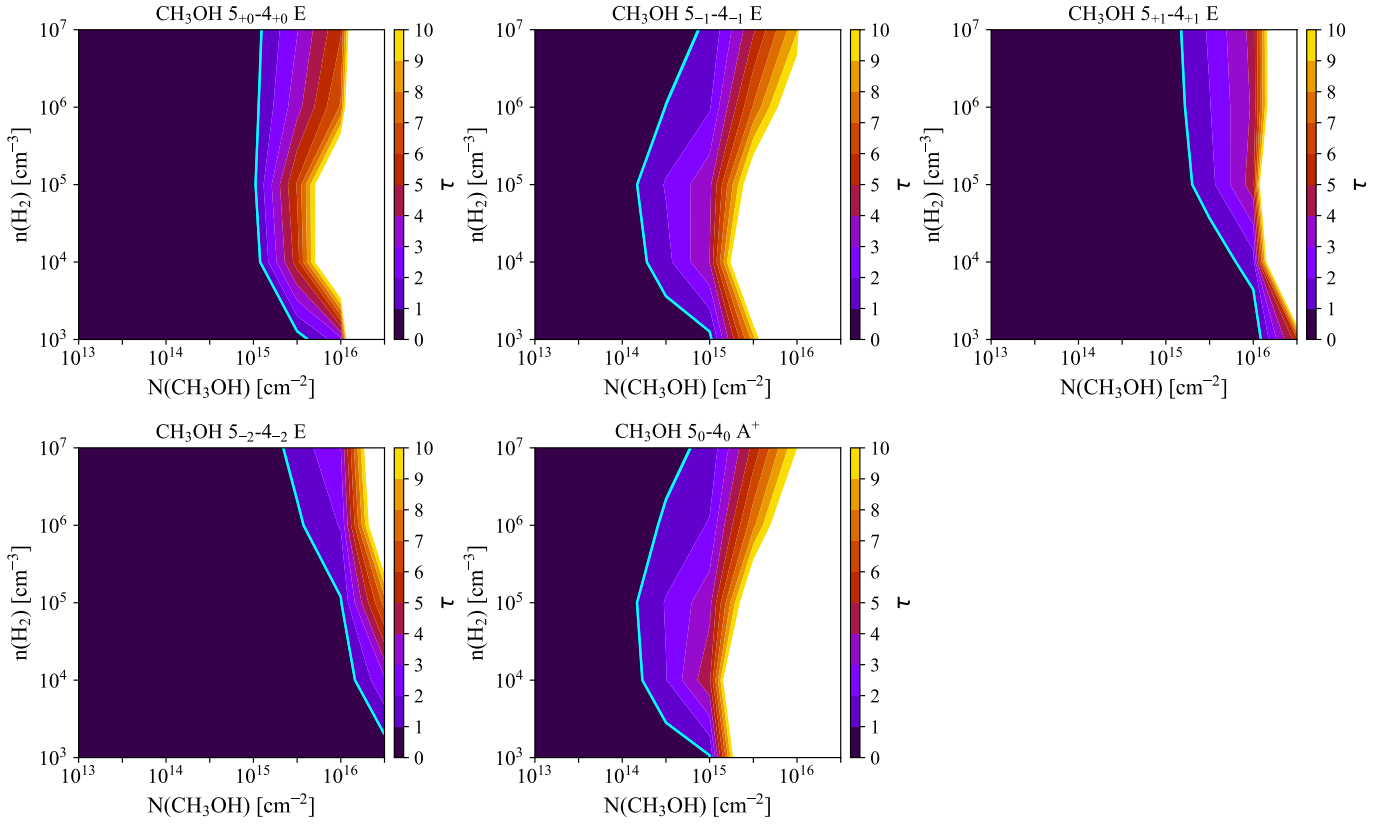


Fig. B.4: Optical depth for the five CH<sub>3</sub>OH transitions obtained with RADEX, using  $T_{\text{kin}} = 30$  K. The cyan line represents  $\tau = 1$ .

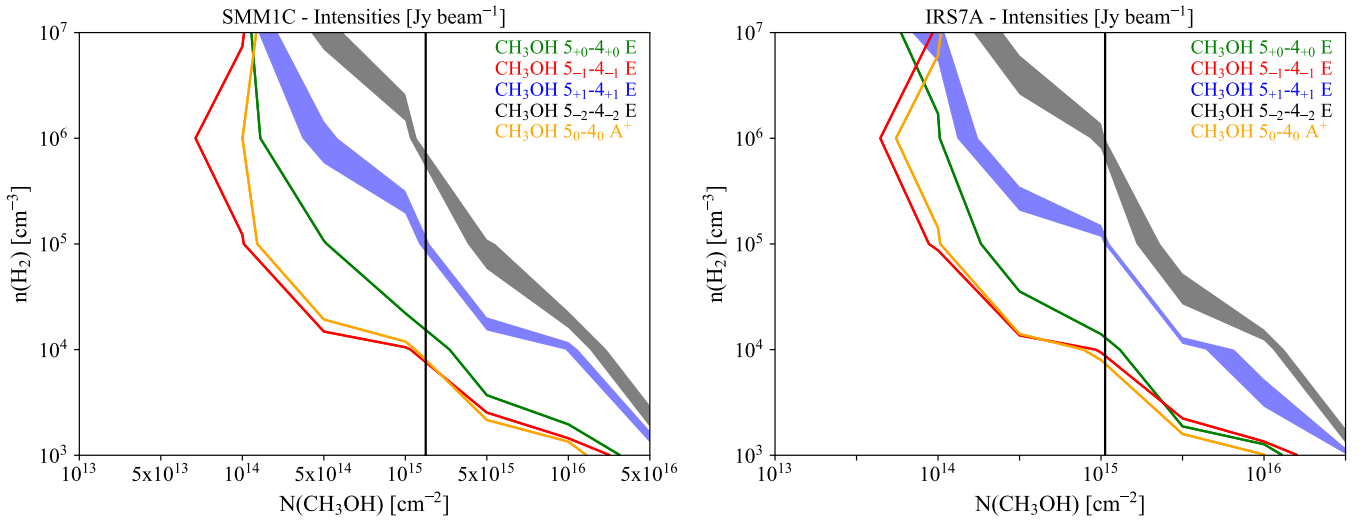


Fig. B.5: RADEX model results for the observed CH<sub>3</sub>OH lines towards SMM1C (left) and IRS7A (right), using  $T_{\text{kin}} = 30$  K. Red, orange and green contours represent lower limits for the integrated line intensities of optically thick lines whereas blue and grey contours represent the marginally optically thin lines and include the errors from Table B.1. The vertical line is the CH<sub>3</sub>OH column density calculated assuming LTE conditions and optically thin emission reported in Table 3.

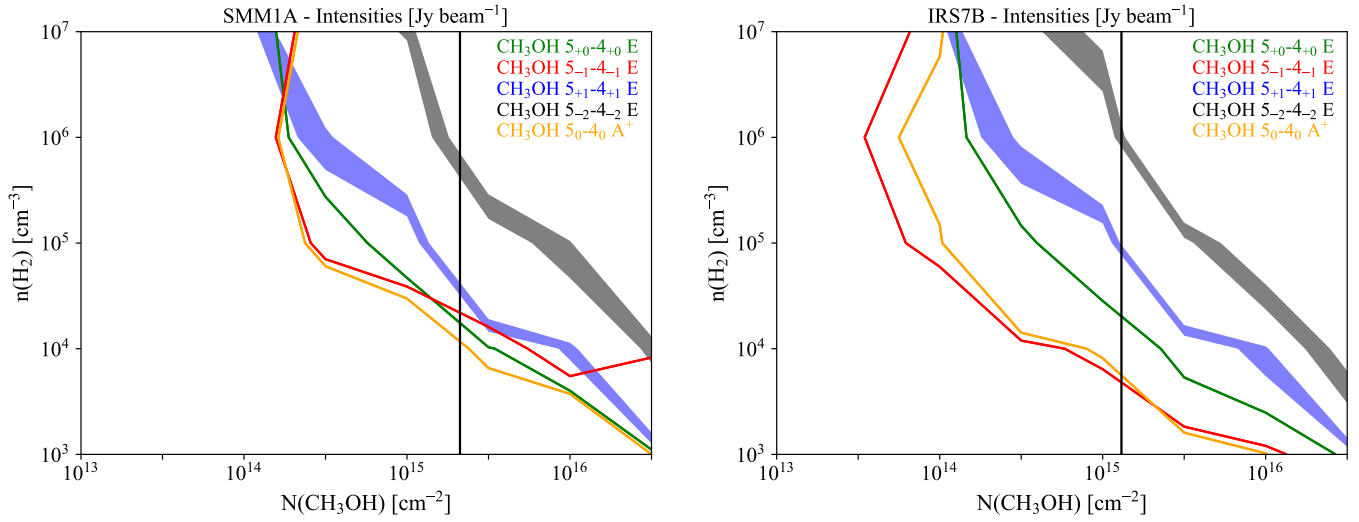


Fig. B.6: Same as Figure B.5, but for SMM1A (left) and IRS7B (right).

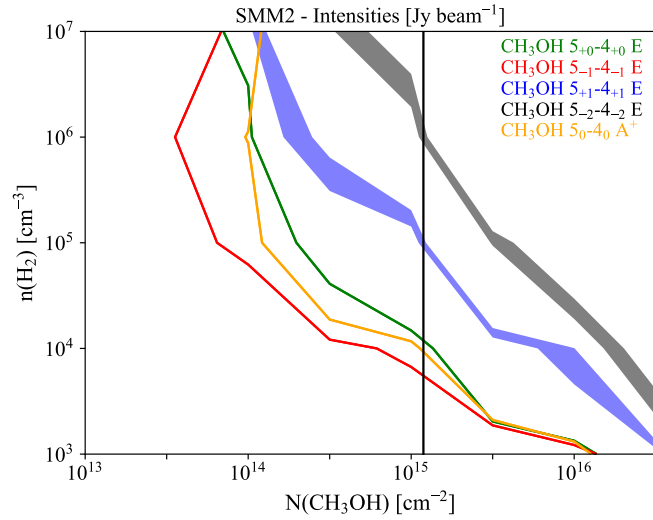


Fig. B.7: Same as Figure B.5, but for SMM2.