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Probing sulphur chemistry in oxygen-rich asymptotic giant branch stars with ALMA

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ABSTRACT

Context. Sulphur and its isotopic ratios play a crucial role in our understanding of the physical properties of astrophysical environments; in particular, providing key insights into nucleosynthesis, interstellar medium processes, star formation, planetary system evolution, and galactic chemical evolution.

Aims. We aim to investigate the distribution of sulphur species – SO_2 , $^{34}\text{SO}_2$, SO , and ^{34}SO – towards a sample of five oxygen-rich asymptotic giant branch (AGB) stars, along with measurements of excitation temperature, column density, and isotopic ratios.

Methods. We used ALMA Band 6, 7, and 8 data of α Ceti, R Dor, W Hya, R Leo, and EP Aqr. SO_2 , $^{34}\text{SO}_2$, SO , and ^{34}SO were detected towards AGB stars using the CASSIS software. To estimate the gas temperature and column density of these species, we applied the rotational diagram method (when applicable) and the Markov chain Monte Carlo method, assuming local thermodynamic equilibrium (LTE). Finally, line imaging of different transitions was performed to infer the distributions of the detected sulphur-bearing species in our sample.

Results. The measured excitation temperatures of SO_2 for our sample sources range from ~ 200 to 600 K, with estimated column densities in the range of $1\text{--}7 \times 10^{16} \text{ cm}^{-2}$. The excitation temperatures estimated using $^{34}\text{SO}_2$ are comparable or slightly lower, while the column densities are about an order of magnitude lower than those of SO_2 . Our measured $^{32}\text{S}/^{34}\text{S}$ ratios for R Dor and W Hya are close to the solar value; however, the measured value for α Ceti is slightly higher, and the measured values for EP Aqr and R Leo are lower. Finally, spatial analysis shows that most detected lines appear as centralized emissions. Moreover, the high excitation transitions of SO_2 show compact emission and probe hot gas of the inner region circumstellar envelopes (CSEs), whereas low-excitation transitions trace slightly extended structures. However, we find some differences in the emission of detected species across our sample.

Conclusions. The excitation temperature of the observed regions of the CSE can be probed using the SO_2 molecule. The morphological correlation between SO and SO_2 emissions suggests that they are chemically linked. Differences in the emission distributions of the detected species across our sample of low mass-loss rate AGB stars such as (i) centralized emission towards α Ceti with irregular emission shapes, (ii) centralized emission with ordered circular features towards R Leo and W Hya, (iii) clumpy emission features in R Dor, and (iv) unresolved emission in EP Aqr may arise from several factors, i.e. the physical conditions of the sources (e.g. density and temperature structures of the CSEs), source multiplicity, outflows, rotation, or other associated physical processes such as thermal and nonthermal desorption, the effects of UV photons and cosmic rays, and finally the resolution of our observations. Nonetheless, the predominantly centralized distributions of SO and SO_2 in our sample support previous findings for low mass-loss rate AGB stars. Our measured $^{32}\text{S}/^{34}\text{S}$ ratios for the two stars R Dor and W Hya agree well with solar values within uncertainties, indicating that these ratios likely reflect the isotopic composition of the stars' natal clouds and deviate for three stars (α Ceti, R Leo, and EP Aqr), which could be due to the metallicity and/or excitation conditions within various sources.

Key words. astrochemistry – line: identification – instrumentation: interferometers – stars: abundances – stars: AGB and post-AGB – circumstellar matter

1. Introduction

When a star with an initial mass of $1\text{--}8 M_{\odot}$ approaches the end of its life and enters the asymptotic giant branch (AGB) phase, during which it rapidly loses mass (in the range $10^{-8}\text{--}10^{-4} M_{\odot} \text{ yr}^{-1}$) and consequently creates an expanding circumstellar envelope (CSE). In the AGB phase, the star undergoes periodic events known as thermal pulses, which trigger dredge-up episodes that bring carbon from the interior to the surface. If the carbon abundance at the surface exceeds that of oxygen, the star transforms into a carbon star. Depending on the C/O ratio, an AGB star can

be classified as oxygen-rich when $\text{C/O} < 1$, carbon-rich when $\text{C/O} > 1$, and S-type where $\text{C/O} \sim 1$.

Sulphur is the tenth most abundant element by mass in the Universe, and it plays a crucial role in biological systems (Mifsud et al. 2021; Krijt et al. 2023). Sulphur has four stable isotopes with fractional abundances of ^{32}S (94.85%), ^{33}S (0.76%), ^{34}S (4.36%), and ^{36}S (0.02%), reflecting their relative abundances at the time of the birth of the Sun (Asplund et al. 2021). The sulphur isotope ratio provides crucial and complementary information on stellar synthesis that is not traced by carbon (Yan et al. 2023). It is believed that ^{32}S and ^{34}S are mainly synthesized during the oxygen-burning process of Type II and Type Ia supernovae (SNe) and ^{33}S synthesized in explosive

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oxygen and neon-burning (Woosley & Weaver 1995; Yan et al. 2023).

Asplund et al. (2021) report a ratio of $^{32}\text{S}/^{34}\text{S} \sim 21.7$ for the sun. Danilovich et al. (2020) estimate $^{32}\text{S}/^{34}\text{S}$ ratios for two oxygen-rich AGB stars, 18.5 ± 5.8 for R Dor, in agreement with the solar value, and 42 for IK Tau, which is considerably higher than the solar value. Recently, Wallström et al. (2024) measured the $^{32}\text{S}/^{34}\text{S}$ ratio for the S-type star W Aql as 16.7 ± 5.6 , and two high-mass loss rate oxygen-rich AGB stars, IRC+10011 and IRC-10529, as 25 ± 12.5 , and 9.10 ± 5 , respectively. Their measured $^{32}\text{S}/^{34}\text{S}$ ratios are consistent with the solar value within uncertainties for W Aql and IRC+10011, but differ for IRC-10529 star. Unnikrishnan et al. (2024) report this ratio for three carbon-rich AGB stars, 15194-5115, 15082-4808, and 07454-7112, as 19 ± 4 , 28 ± 8 , and 21 ± 4 , respectively. These measured values are in good agreement with the solar value, considering their reported uncertainties, except for the high-mass oxygen-rich AGB star, IRC-10529. The reported lower value for IRC-10529 compared to the solar value of 21.7 is not discussed. We speculate that comparing the measured ratio values with the solar value may highlight variations in elemental abundances between the natal clouds of AGB stars and those of the Sun.

Several sulphur-bearing species, such as CS, SO, SO₂, and H₂S, are commonly detected in the circumstellar environments of AGB stars and star-forming regions (e.g. Omont et al. 1993; Danilovich et al. 2017b; Massalkhi et al. 2020; Fontani et al. 2023; Ghosh et al. 2024). Several studies of sulphur towards oxygen-rich AGB stars have been carried out (e.g. Lindqvist et al. 1988; Danilovich et al. 2016, 2017b, 2018; Danilovich et al. 2020). Danilovich et al. (2016) analysed the SO and SO₂ spatial distributions using single-dish observations towards five oxygen-rich AGB stars (IK Tau, R Dor, TX Cam, W Hya, R Cas) with different mass-loss rates.

They found that higher-mass-loss rate stars show shell-like distributions of SO and lower peak relative abundances. The lower-mass-loss rate stars show centralized SO distributions, with higher peak abundances close to the stars. The locations of the SO peaks (for the higher mass-loss rate stars) and the e-folding radii (for the lower mass-loss rate stars) are correlated with the OH peak abundance and the photodissociation of H₂O. Later, Danilovich et al. (2020) presented spatially resolved observations of SO and SO₂ towards two AGB stars: one with a high-mass loss rate, IK Tau, and another with a low-mass loss rate, R Dor. In the case of R Dor, the emission of two sulphur species coincides with peaks around the central star, which trace out the same density structures in the circumstellar environment. In the case of IK Tau, SO shows a shell-like structure; its peak does not appear at the star's centre, and most of the flux is resolved for the low-excitation SO₂ transition.

Massalkhi et al. (2020) observed SO and SO₂ in oxygen-rich CSE of AGB stars. They further compared their distribution with SiO to determine whether SO and SO₂ can be used as a dust precursor. They looked into the SiO abundance profile and found a decreasing trend with the envelope density, similar to C-rich AGB stars, indicating SiO adsorption onto the dust grains. Interestingly, they obtained a similar trend in the case of SO, but not as prominent as for SiO. On the other hand, SO₂ does not show such a trend, which indicates that it is a less important candidate for the precursor of the dust. Recently, Wallström et al. (2024) presented high-resolution ALMA observations of SO and SO₂ in a large sample of 17 oxygen-rich AGB stars with a variation of mass-loss rates. They show that the spatial distributions of SO and SO₂ are generally consistent with previous results, with a centralized distribution for low

mass-loss rate sources and a shell-like distribution for high-mass-loss rate sources.

In this paper, we report several transitions of SO, SO₂, and their isotopologues ^{34}SO and $^{34}\text{SO}_2$ in a wider energy range observed with ALMA Bands 6, 7, and 8 towards a sample of five oxygen-rich AGB stars, all of which have mass-loss rates of the order of $\sim 10^{-7} \text{ M}_{\odot} \text{ yr}^{-1}$. The paper is organized as follows. Section 2 describes the observations and data analysis, and the results are presented in Section 3. Discussions are summarized in Section 4, and finally the conclusion is given in Section 5.

2. Observations and data analysis

In this paper, we used the Atacama Large Millimeter/submillimeter Array (ALMA) Band 8 observations with the Atacama Compact Array (ACA) (2018.1.01440.S, PI: M. Saberi) configuration for four sources (R Dor, R Leo, W Hya, and EP Aqr) and the 12m array (2018.1.00649.S, PI: M. Saberi) for one source (*o* Ceti). Additionally, we used 12m array observations of *o* Ceti with ALMA Band 7 (2018.1.00749.S, PI: T. Khouri). For R Dor, we used ALMA 12m array Band 7 observations (2017.A.00012.S, PI: L. Decin). We also used 12m array observations of R Leo with the ALMA Band 7 set-up (2019.1.00801.S, PI: J. Champion) and 12m array observations of W Hya with ALMA Band 6 (2016.1.00374.S, PI: K. Ohnaka). Notably, we used a more or less similar angular resolution for all sources: approximately $\sim 2.5''$ with ACA and $\sim 0.2''$ with the 12m array. The targets, along with their co-ordinates, distances, LSR velocity (V_{LSR}), expansion velocities, masses, temperatures, radii, and mass-loss rates, are listed in Table 1. Table 2 provides a summary of the observations, which includes information about ALMA Band, configuration, number of antennas, observation date, angular resolution, maximum recoverable scale (MRS), field of view (FOV), sensitivity, and spectral resolution.

Data calibration and cleaning were performed following standard ALMA procedures using the Common Astronomy Software Application (CASA) (McMullin et al. 2007). The uvcontsub task was used to separate the continuum part from the line emissions. Calibration uncertainty depends on the flux calibrator used and typically ranges from 5% to 20% (Francis et al. 2020). We used CASSIS¹ software for line identification and further analysis (Vastel et al. 2015).

2.1. Line identification

The line identification of all the observed species was carried out using CASSIS software together with the Cologne Database for Molecular Spectroscopy (CDMS, Müller et al. 2001, 2005)² database. First, we applied a cut based on the rms noise for all data, then we consider the spectral profiles that are above 3σ , where σ is the rms noise. To firmly identify a molecular transition corresponding to the observed spectra, we checked line blending, systematic velocity of the source (V_{LSR}), and the Einstein coefficient of different nearby transitions. We chose a few specific transitions (if multiple potential lines are available) corresponding to a particular spectral peak. Finally, to firmly assign a molecular species to the observed spectral feature, we used the local thermodynamic equilibrium (LTE) model using Markov chain Monte Carlo (MCMC) fitting within CASSIS (see Sect. 2.2) and checked whether the synthetic spectrum of a particular species can reproduce the observed profile or not.

¹ <http://cassis.irap.omp.eu>

² <https://www.astro.uni-koeln.de/cdms>

Table 1. Our sample: M-type AGB stars.

Star	RA (J2000)	Dec (J2000)	\dot{M} ($10^{-7} M_{\odot} \text{ yr}^{-1}$)	d (pc)	V_{LSR} (km s^{-1})	V_{exp} (km s^{-1})	R_{\star} (mas)	T_{\star} (K)	P (days)	M_{\star} (M_{\odot})
<i>o</i> Ceti	02 ^h 19 ^m 20 ^s .795	−02°58′47.313″	2.0	102	47	3	15	2800	332	1.0
R Leo	09 ^h 47 ^m 33 ^s .480	+11°25′42.883″	1.0	130	0.5	6	13.5	2800	310	1.5
W Hya	13 ^h 49 ^m 01 ^s .925	−28°22′04.648″	1.3	104	41	7	20	2950	388	1.0
R Dor	04 ^h 36 ^m 45 ^s .400	−62°04′39.259″	1.6	45	6.9	6	27.5	2400	175	1.0–1.3
EP Aqr	21 ^h 46 ^m 31 ^s .880	−02°12′45.547″	1.2	114	−34	11	—	—	—	—

Notes. Columns 4 to 11 (Same as reported in Saberi et al. 2022), for the first four sources, are taken from Vlemmings et al. (2019). The stellar masses for W Hya and R Dor are from Danilovich et al. (2017a), while those for *o* Ceti and R Leo are estimated based on the $^{17}\text{O}/^{18}\text{O}$ ratio presented by De Nutte et al. (2017), and the O isotopic ratios are taken from Hinkle et al. (2016). For EP Aqr, all parameters are taken from Nhung et al. (2015, and references therein). We note that R_{\star} listed here are the measured radii in the near-infrared. †The period of R Dor varies between 175 and 362 days.

Table 2. Observation summary.

Source	ALMA Band	Configuration	No. of Antennas	Obs. Date	Ang. res. (arcsec)	MRS (arcsec)	FOV (arcsec)	Sensitivity (mJy/beam)	Spec. Resolution (kHz) (km s^{-1})
<i>o</i> Ceti	8	12m	46	28-11-18	0.172	3.212	11.96	0.16	1938.48 (~1.2)
	7	12m	44	19-11-18	0.212	3.142	20.08	0.05	1128.91 (~1.1)
	7	12m	46	26-11-18	0.213	3.033	17.21	0.06	1128.91 (~1.0)
R Leo	8	ACA	9	24-03-19	2.570	15.118	20.50	2.39	1938.48 (~1.2)
	7	12m	44	06-10-19	0.218	3.288	16.76	0.05	1128.91 (~1.0)
W Hya	8	ACA	9	08-06-19	2.039	13.634	20.50	2.17	1938.48 (~1.2)
	6	12m	44	08-07-17	0.156	2.117	22.43	0.02	976.56 (~1.1)
R Dor	8	ACA	9	08-06-19	2.135	13.812	20.50	1.79	1938.48 (~1.2)
	7	12m	44	15-09-18	0.194	2.993	16.05	0.07	1128.91 (~1.1)
Ep Aqr	8	ACA	9	24-06-19	2.452	13.939	20.50	1.35	1938.48 (~1.2)

2.2. Markov chain Monte Carlo (MCMC) fitting

MCMC fitting was employed to fit the observed line profiles of different transitions towards a sample of five AGB stars across three different ALMA bands. For the fitting, we assumed that the excitation temperature is equivalent to the gas temperature under LTE conditions. We used the Python scripting interface available in CASSIS for our model calculations to determine the best-fit physical parameters for the SO_2 and $^{34}\text{SO}_2$ transitions. We varied input parameters such as the excitation temperature, line width, and column density, while keeping the emission regions of the observed transitions, and V_{LSR} , fixed (see Table 1). After applying the LTE approximation using the MCMC approach, we extracted the best-fit physical parameters such as column density, excitation temperature, and full width at half maximum (FWHM). We applied the χ^2 minimization process to determine the best-fit model that matches the observed line profiles.

2.3. Rotation diagram analysis

ALMA observations of our AGB sample contain multiple transitions of SO_2 and $^{34}\text{SO}_2$. Hence, a rotation diagram analysis was carried out to determine the excitation temperatures and column densities. Assuming that the observed transitions of these species are optically thin and are in LTE, we performed rotational diagram analysis. To check the optically thin assumption, we estimated the optical depth (τ) of each molecular line following the equation

$$\tau = -\ln\left(1 - \frac{T_{\text{peak}}}{\phi(J(T_{\text{ex}}) - J(T_{\text{bg}}))}\right), \quad (1)$$

where T_{peak} is the peak emission (peak flux in K from Table A.1, ϕ is the beam filling factor, assumed to be 1, and

$$J_{\nu}(\nu, T_{\text{ex}}) = \left(\frac{h\nu}{k}\right) \frac{1}{\exp\left(\frac{h\nu}{kT_{\text{ex}}}\right) - 1}, \quad (2)$$

where ν is the frequency of the molecular transition, and T_{ex} was obtained from the rotational diagram of each individual source, $T_{\text{bg}} = 2.73$ K. The calculated values of τ are provided in the last column of Table A.1. All values of τ are $\ll 1$, which indicates that the lines are optically thin. For optically thin lines, column density can be expressed as (Goldsmith & Langer 1999)

$$\frac{N_u^{\text{thin}}}{g_u} = \frac{3k_B \int T_{\text{mb}} dV}{8\pi^3 \nu S \mu^2}, \quad (3)$$

where g_u is the degeneracy of the upper state, k_B is the Boltzmann constant, $\int T_{\text{mb}} dV$ is the integrated intensity, ν is the rest frequency, μ is the electric dipole moment, and S is the transition line strength. Under LTE conditions, the total column density can be written as

$$\frac{N_u^{\text{thin}}}{g_u} = \frac{N_{\text{total}}}{Q(T_{\text{rot}})} \exp(-E_u/k_B T_{\text{rot}}), \quad (4)$$

where T_{rot} is the rotational temperature, E_u is the upper state energy, and $Q(T_{\text{rot}})$ is the partition function at rotational temperature. Equation (2) can be rearranged as

$$\ln\left(\frac{N_u^{\text{thin}}}{g_u}\right) = -\left(\frac{E_u}{k_B T_{\text{rot}}}\right) + \ln\left(\frac{N_{\text{total}}}{Q(T_{\text{rot}})}\right). \quad (5)$$

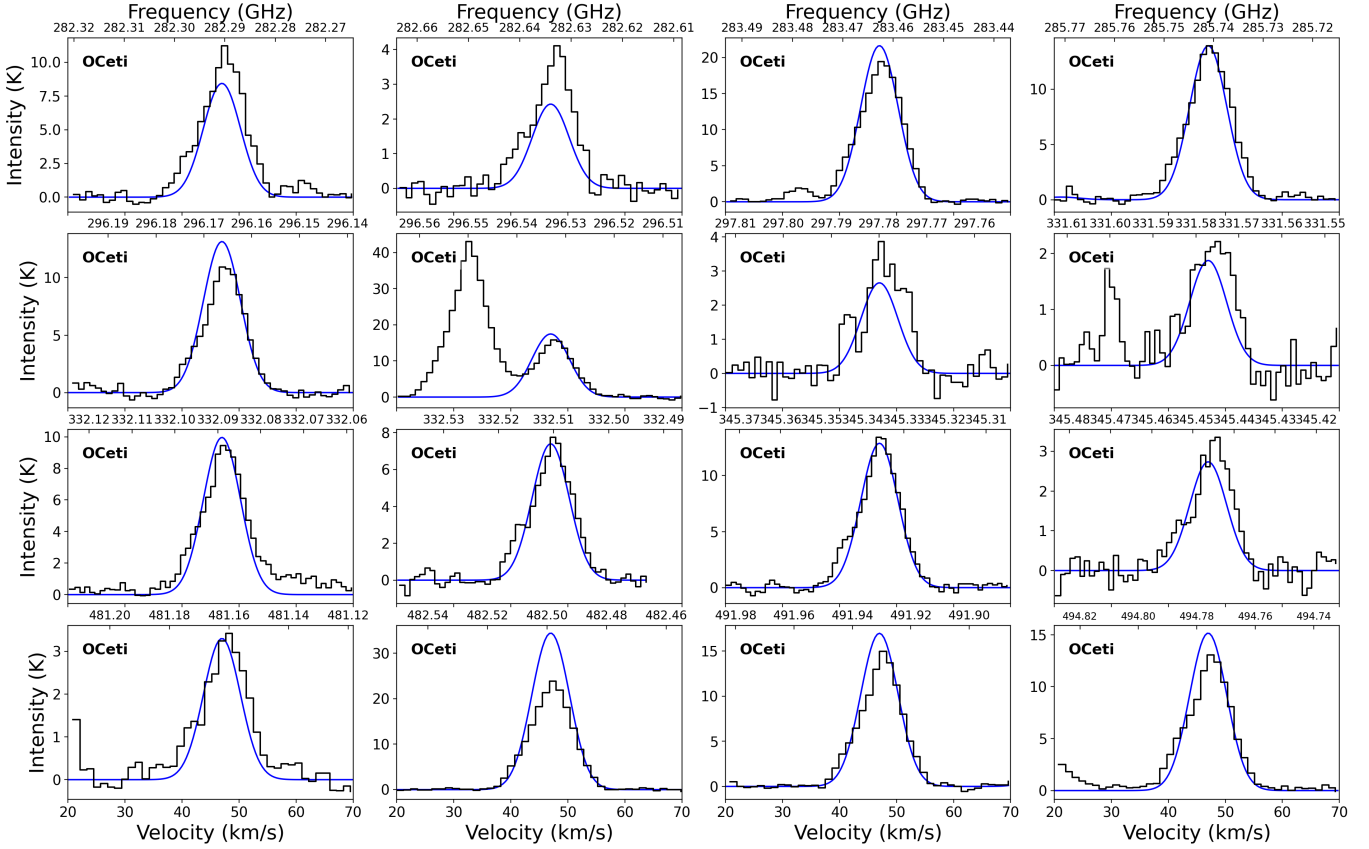


Fig. 1. Observed and modelled spectra of SO₂ transitions towards *o* Ceti. The black line represents the observed spectrum, while the blue line represents the modelled spectrum. The strong line at 296.5500 GHz, shown in the second subplot of the third row from the bottom, represents the SO line. The narrow line in the last subplot of the same row corresponds to TiO₂ (331.5996 GHz).

For the rotational diagram analysis, we combined the observed transitions from ACA and 12m array observations whenever they were available. Since the resolution of the ACA differs from that of the 12 m array, we applied a beam dilution factor. We examined the emitting region of both SO₂ transitions in ACA and 12m observations and found that the emission is unresolved in our ACA observations. On average, we extracted the spectra with a diameter of 1'' for 12m array and 3'' for ACA. Hence, we applied a scaling factor of 9 to estimate the upper-level column density of transitions observed with ACA. Finally, we used all data points in the rotational diagram analysis to obtain the rotational temperatures and column densities. We applied a rotational diagram (for species with more than two lines observed) and MCMC fitting to estimate the excitation temperature and column density. The error bars in the rotational diagram come from the Gaussian fitting error of the line profiles.

To investigate the distribution of SO₂ and ³⁴SO₂ in the oxygen-rich AGB stars, we created moment 0 maps of all transitions from both ACA and 12m array observations. Notably, the resolution of our ACA data cannot resolve the emission of SO₂, whereas 12m array observations have sufficient resolution to resolve the sulphur species spatially.

3. Results

Figure 1 depicts the observed and modelled spectra towards *o* Ceti. The spectra of other sources are provided in the appendix (see Figs. B.1 and B.2). Table A.1 summarizes all the detected transitions. The line width (FWHM), ΔV , V_{LSR} , and the

integrated intensity ($\int T_{\text{mb}} dv$) of each transition were measured by applying a single Gaussian to fit the observed spectrum. The line parameters of all the observed transitions such as the rest frequency (ν_0), quantum numbers ($J'_{K'_a K'_c} - J''_{K''_a K''_c}$), FWHM, $\int T_{\text{mb}} dv$, upper state energy (E_{up}), and V_{LSR} are noted in Table A.1.

The rotational diagrams of SO₂ and ³⁴SO₂ (when available) are shown in Figs. 2, 3, 4, 5, and 6 for *o* Ceti, R Leo, R Dor, W Hya, and EP Aqr, respectively. Table 3 compares the derived excitation temperatures and column densities from both RD and MCMC methods. The results of individual sources are discussed in the following subsections.

3.1. *o* Ceti

Many transitions of SO₂ (16 lines), ³⁴SO₂ (14 lines), SO (3 lines), and ³⁴SO (2 lines) are detected with a wide range of upper state energy ($E_{\text{up}} \sim 40\text{--}700$ K; see Table A.1). Additionally, we detected one transition of ³³SO (343.0861 GHz). We also find a line close to the frequency of a ³³SO₂ line (345.5852 GHz), but do not associate it with ³³SO₂ emission because other lines potentially expected to be seen (at 283.338 GHz, 296.270 GHz, 334.030 GHz, and 479.960 GHz) are not observed. We have multiple transitions of SO₂ and ³⁴SO₂, and hence applied a rotation diagram analysis (see Fig. 2). The rotational diagram analysis yields excitation temperatures and column densities of $T_{\text{ex}} = 542 \pm 168$ K and $N_{\text{SO}_2} = (7.06 \pm 3.55) \times 10^{16} \text{ cm}^{-2}$ for SO₂ and $T_{\text{ex}} = 214 \pm 71$ K and $N_{^{34}\text{SO}_2} = (2.16 \pm 0.87) \times 10^{15} \text{ cm}^{-2}$ for ³⁴SO₂. Our estimated temperature of SO₂ is comparable with the

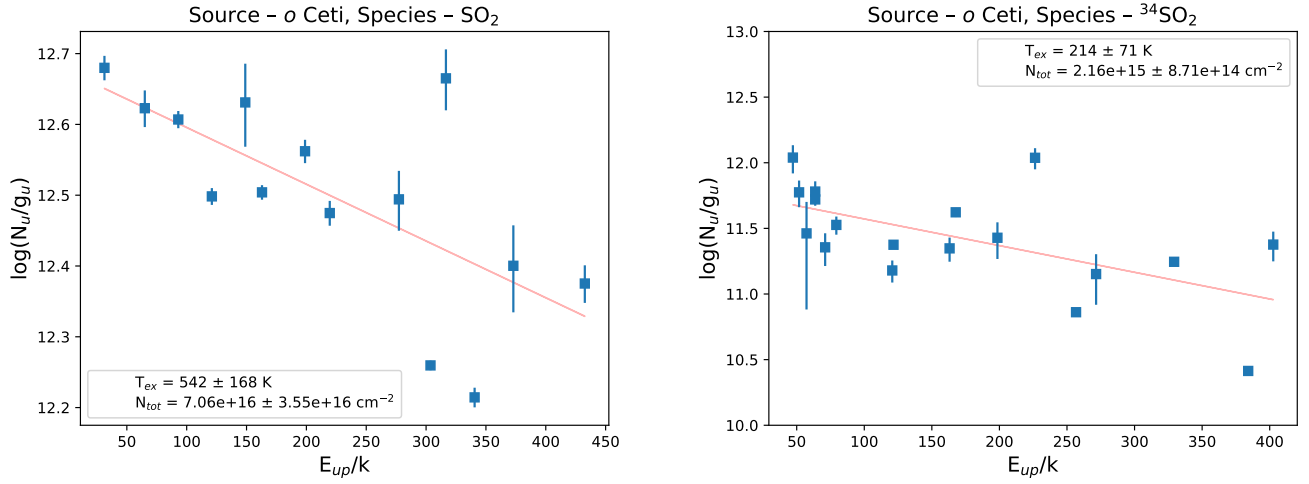


Fig. 2. Rotational diagram of SO₂ and ³⁴SO₂ for o Ceti. The filled blue squares represent the data points, and the vertical lines on each data point indicate the error bars. The solid red lines represent the fitted line.

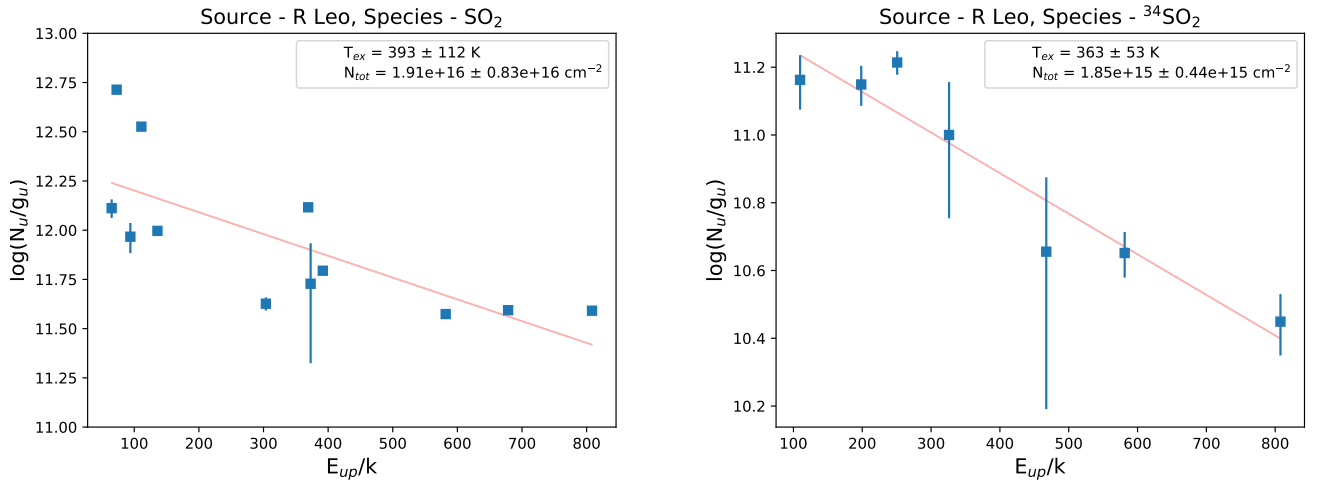


Fig. 3. Rotational diagram of SO₂ and ³⁴SO₂ for R Leo. The filled blue squares represent the data points, and the vertical lines on each data point indicate the error bars. The solid red lines represent the fitted straight line.

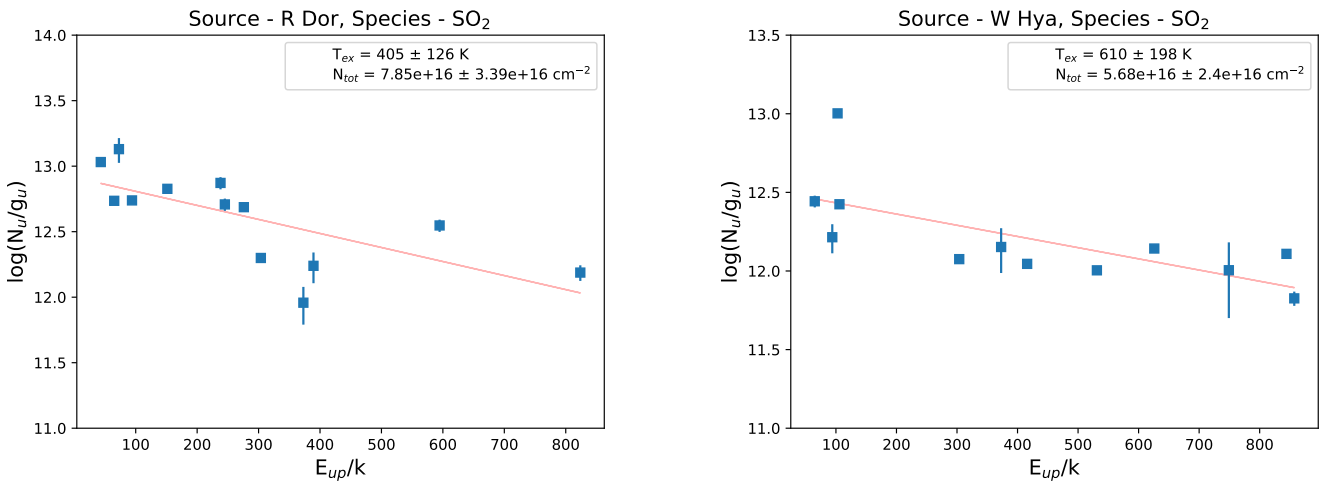


Fig. 4. Rotational diagram of SO₂ for R Dor. The filled blue squares represent the data points, and the vertical lines on each data point indicate the error bars. The solid red line represents the fitted straight line.

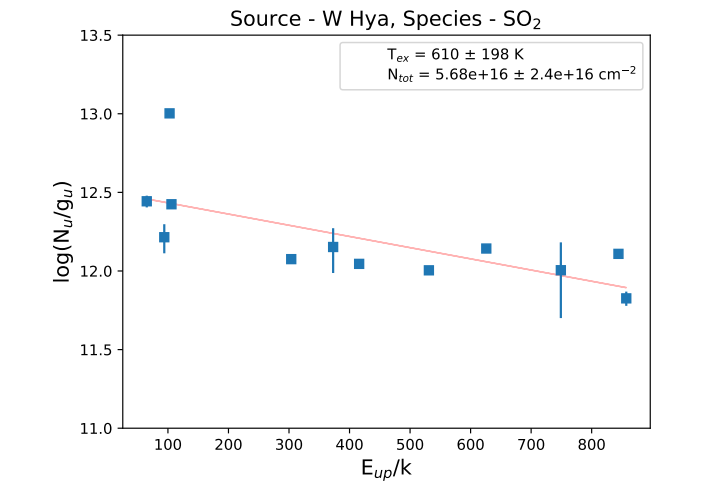
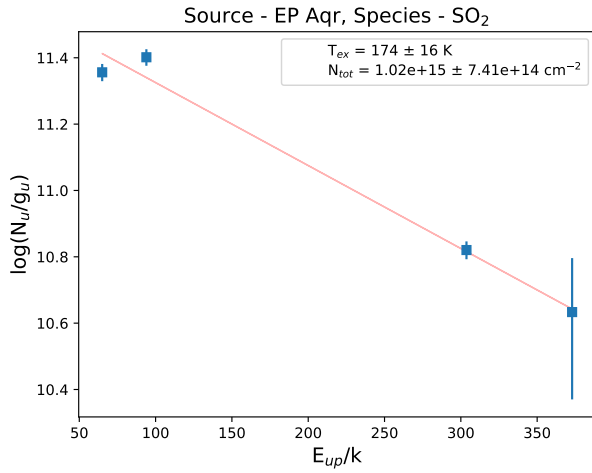


Fig. 5. Rotational diagram of SO₂ for W Hya. The filled blue squares represent the data points, and the vertical lines on each data point indicate the error bars. The solid red line represents the fitted straight line.

Table 3. Rotational temperature and column density of SO₂ and ³⁴SO₂ from two methods.

Source	$T_{\text{ex}}(\text{SO}_2)$ (K)		$N(\text{SO}_2)$		$T_{\text{ex}}(^{34}\text{SO}_2)$ (K)		$N(^{34}\text{SO}_2)$		$N_{\text{SO}_2}/N_{^{34}\text{SO}_2}$
	RD	MCMC	RD	MCMC	RD	MCMC	RD	MCMC	
<i>o</i> Ceti	542±168	410±36	7.06e+16	6.76e+16	214±71	296±4	2.16e+15	2.21e+15	32±6
R Leo	393±112	489±29	1.91e+16	2.49e+16	363±53	436±48	1.85e+15	1.89e+15	10±3
R Dor	405±126	364±26	7.85e+16	4.05e+16	–	258±54	–	1.55e+15	26±8
W Hya	610±198	419±4	5.68e+16	5.54e+16	–	573±22	–	3.17e+15	18±4
EP Aqr	174±16	209±4	1.02e+15	1.97e+15	–	176±37	–	1.61e+14	12±5

Notes. RD – rotational diagram method; MCMC – Monte Carlo Markov chain method. For the MCMC fitting, we used the following emission regions: *o* Ceti – 1'' (~102 au), R Dor – 1.5'' (~68 au), R Leo 1'' (~130 au), W Hya – 1'' (~104 au), and EP Aqr – 3'' (~330 au), which were obtained from the moment 0 map analysis.

**Fig. 6.** Rotational diagram of SO₂ for EP Aqr. The filled blue squares represent the data points, and the vertical lines on each data point indicate the error bars. The solid red line represents the fitted straight line.

previous measurement based on TiO (474 ± 69) K (Kamiński et al. 2017) and AIF observations within uncertainties (Saberi et al. 2022). For *o* Ceti, a very faint signal at 493.3622 ($E_{\text{up}} \approx 564$ K) GHz transition of H₂S is observed towards the central source, but it is below the detection limit. Assuming LTE, we created synthetic spectra to match the observed profile. For this, we adopted the same temperature of 542 K, obtained from the RD of *o* Ceti, and estimated the upper limit for the column density of H₂S to be $6.53 \times 10^{14} \text{ cm}^{-2}$. However, we did not detect any signal of the high-excitation transition ($E_{\text{up}} \approx 564$ K) towards the arc-like structure of *o* Ceti. Danilovich et al. (2017b) found that H₂S likely plays an important role in oxygen-rich AGB stars with high mass-loss rates, but is unlikely to be significant in stars of other chemical types or in oxygen-rich stars with lower mass-loss rates.

Figures C.1–C.3 show the moment 0 maps of all detected transitions (also see Fig. 1 in the supplementary document). A spatial analysis with high-resolution (~0.2'') observations in different ALMA bands shows centralized emission for all transitions with emitting regions around ~1'' towards the main continuum source. The shape of the emission shows irregular appearances, which could be due to the binary companion. The average separation between binary sources is 0.472'' (Vlemmings et al. 2015), corresponding to ~48 au, considering a distance of 102 pc. Nonetheless, we cannot distinguish the distinct features of SO and SO₂ around Mira A and B with the present angular resolution.

Furthermore, for several transitions ($E_{\text{up}} < 200$ K) we find extended emission in the south-west position, which is completely offset and unconnected to the emission that is close to the continuum peak. For example, the low-excitation transitions of SO and SO₂ (see Figs. C.4 and C.5) show an extended emission, about 3'' (~300 au) offset from the continuum position. Because of the lower intensity of the offset emission, it does not appear as a prominent feature in the integrated intensity maps. Hence, we plotted channel maps for these transitions, in which the extended emission is more clearly visible. In these figures, we see the extended feature in a few channels around $\pm 3 \text{ km s}^{-1}$ of the $V_{\text{LSR}} = 47 \text{ km s}^{-1}$. Previously, Wong et al. (2016) reported the extended spatially resolved emission in SiO ($J = 5-4$) ALMA observation around the same position. Our observations show SO and SO₂ emissions around the same region where SiO appears. Since SiO is a good tracer of shocks (Fontani et al. 2019), the formation of gas-phase SO and SO₂ is probably caused by shock-induced processes in the same environment. It could also be a region with higher gas density. The morphological correlation between SO, SO₂, and SiO emissions implies that the shock processes might enhance the gas-phase abundance of these molecules. Shocks can significantly influence chemistry by increasing temperatures and densities and disrupting dust grains. In particular, shock waves can sputter dust grains and ice mantles (Suutarinen et al. 2014), releasing sulphur-bearing species such as atomic sulphur (S), hydrogen sulphide (H₂S), or sulphur monoxide (SO) into the gas phase. These species can then undergo rapid gas-phase reactions to form SO₂. Furthermore, we estimated that the temperature of this region is $\sim 133 \pm 20$ K, based on an LTE analysis of all detected transitions of SO and SO₂ towards the offset position (see Table A.2).

3.2. R Dor

A list of detected transitions of SO₂ (14 lines), ³⁴SO₂ (2 lines), SO (1 line), and ³⁴SO (1 line) towards R Dor is provided in Table A.1. We applied rotational diagram analysis for SO₂ to obtain the excitation temperature (see Fig. 4). Our derived excitation temperature is $T_{\text{ex}} = 405 \pm 126$ K and $N_{\text{SO}_2} = (7.85 \pm 3.39) \times 10^{16} \text{ cm}^{-2}$. We do not have multiple transitions of ³⁴SO₂ and SO. Hence, we measured the excitation temperature and the column density of these species using the MCMC method. Our estimated excitation temperature using SO₂ is consistent with previous measurements: ~505 K (Khouri et al. 2024). Moment 0 maps of different transitions of SO₂ are shown in Fig. C.6 (also see Fig. 2 in the supplementary document). It is clear from Fig. 2 in the supplementary document that SO₂ is unresolved with ACA observation, but with high-resolution

data it is well resolved and shows clumpy emission structures (see Fig. C.6) where the continuum peak is at the centre. Furthermore, emission shows more compact features for transitions with high upper-state energy than the lower E_{up} transition. Our results, obtained at an angular resolution of $0.2''$, are consistent with previous ALMA observations at a comparable resolution ($\sim 0.15''$) (Danilovich et al. 2020). The co-spatial emission of SO and SO_2 (see Fig. C.6) suggests that both molecules trace the same wind structure within the circumstellar environment. Furthermore, since the MRS of our observations is approximately $3''$, we compared the observed intensities of transitions with similar E_{up} in our data and previous results obtained using the APEX 12m single-dish telescope (Danilovich et al. 2016). This comparison was performed to determine whether any flux has been resolved out in the interferometric data. However, we found good agreement between the single-dish and interferometric results for transitions with comparable E_{up} .

3.3. R Leo

All detected transitions of SO_2 (10 lines), $^{34}\text{SO}_2$ (7 lines), SO (1 line), and ^{34}SO (1 line) towards R Leo are summarized in Table A.1. We have multiple transitions of SO_2 and $^{34}\text{SO}_2$. Hence, we applied a rotational diagram to estimate the excitation temperature and column density (see Figs. 3). The rotational diagram analysis yields excitation temperatures and column densities of $T_{\text{ex}} = 393 \pm 112$ K and $N_{\text{SO}_2} = (1.91 \pm 0.83) \times 10^{16} \text{ cm}^{-2}$ for SO_2 and $T_{\text{ex}} = 363 \pm 53$ K and $N_{^{34}\text{SO}_2} = (1.85 \pm 0.44) \times 10^{15} \text{ cm}^{-2}$ for $^{34}\text{SO}_2$. For SO, we do not have multiple transitions. Consequently, we employed the MCMC method to derive a column density of SO of $6.3 \times 10^{15} \text{ cm}^{-2}$ and an excitation temperature of $T_{\text{ex}} = 300$ K. Figures C.7 and C.8 (also see Fig. 3 in the supplementary document) show moment 0 maps of all detected transitions of SO_2 , $^{34}\text{SO}_2$, SO, and ^{34}SO . It is clear from Fig. 3 in the supplementary document that emission of SO_2 is not resolved with a low-resolution ($\sim 2.6''$) observation. In contrast, SO_2 can be resolved well with high-resolution ($\sim 0.22''$) data, as is shown in Figs. C.7 and C.8. We find a morphological correlation in the emission between SO_2 , $^{34}\text{SO}_2$, and SO. The emission from the low-excitation transitions of ($E_{\text{up}} = 81$ K) and $^{34}\text{SO}_2$ ($E_{\text{up}} = 77$ K) appears to be slightly more extended compared to higher-excitation transitions. Overall, the emission from both SO and SO_2 is centrally concentrated and co-spatial, suggesting that these species likely trace the same wind structures in the CSE of R Leo.

3.4. W Hya

The observed transitions of SO_2 (13 lines) and $^{34}\text{SO}_2$ (3 lines) towards W Hya are noted in Table A.1. The estimated excitation temperature using the SO_2 rotational diagram is $T_{\text{ex}} = 610 \pm 198$ K and the column density $N_{\text{SO}_2} = (5.68 \pm 2.4) \times 10^{16} \text{ cm}^{-2}$. Moment 0 maps for different transitions of SO_2 with high- ($\sim 0.16''$) and low-resolution ($\sim 2.0''$) observations are depicted in Fig. C.9 and Fig. 4 in the supplementary document, respectively. In the case of W Hya, emission from SO_2 is not resolved in the ACA observations, similar to R Dor and R Leo. However, we find an additional structure of strong SO_2 emission located approximately $10''$ offset from the central continuum peak. It possibly traces the structures at the outer envelope of W Hya. At high resolution, the SO_2 emission is resolved. It exhibits an ordered, centrally concentrated structure, with a slight offset between the continuum and molecular emission peak, similar

to what is observed in R Leo. However, we do not detect the additional features at the offset position, which we found from ACA observations. The structure at the offset position is most likely resolved out due to the limited MRS of $2.1''$ in the high-resolution observations. We do not have an SO transition in the frequency range of our observation. Nonetheless, modelling results from previous studies suggest that the SO_2 distribution is similar in size and abundance to the circumstellar SO distribution (Danilovich et al. 2016).

3.5. EP Aqr

For EP Aqr, we have several transitions of SO_2 (five lines) and one transition of $^{34}\text{SO}_2$ (one line) but only with ALMA Band 8 ACA observations (see Table A.1). The rotational diagram yields $T_{\text{ex}} = 174 \pm 16$ K and $N_{\text{SO}_2} = (1.02 \pm 0.74) \times 10^{15} \text{ cm}^{-2}$. Figure 5 in the supplementary document shows the moment 0 maps of SO_2 transitions. The slightly lower temperature obtained compared to other sources (see Table 3) could be due to the large emission regions and the lack of high-resolution data with multiple transitions for this source. All transitions are spatially unresolved because of the low angular resolution of $\sim 2.5''$. Tuan-Anh et al. (2019) reported one transition of SO_2 and found that the emission is confined towards the centre of the source, with an emitting diameter of $\sim 0.25''$ (~ 30 au). They suggest that SO_2 could be a useful tracer of the mass-loss mechanism in its early phase. Their analysis also provided clear evidence of rotation, possibly combined with moderate expansion. Furthermore, by applying the temperature radial dependence from Hoai et al. (2019) and considering the radial extent of the SO_2 emission, they estimated a gas temperature of a few hundred Kelvin, consistent with values typical of oxygen-rich AGB stars (Yamamura et al. 1999). However, a multi-transition (with a wide range of E_{up}) with high-angular resolution observations could provide better constraints on both the spatial distribution and the actual gas temperature traced by SO_2 .

4. Discussion

4.1. SO and SO_2 distributions

SO and SO_2 have previously been found by Danilovich et al. (2016) and Wallström et al. (2024) to have centralized distributions for low-mass-loss rate AGB stars and shell-like distributions for high-mass-loss sources. They also observed that sources with shell-like emission are typically detected in lower-energy SO_2 lines, while those with centralized emission are more often detected in higher-energy lines. All five stars in our sample have relatively low mass-loss rates. We investigated the distribution of SO_2 and SO towards these sources. For *o* Ceti, we observe slightly irregular emission contours for SO_2 , $^{34}\text{SO}_2$, SO, and ^{34}SO , and emission peaks are offset from the continuum peak. The emissions of SO_2 and $^{34}\text{SO}_2$ are spatially correlated, as are the emissions of SO and ^{34}SO . However, with the current angular resolution, we cannot distinguish the distributions of our detected species between the Mira A and B components. Furthermore, we detected extended emission for several transitions of SO and SO_2 and some other species (see Table A.2) at the south-west position, which is $\sim 3''$ offset from the main continuum and molecular emission peaks (e.g. see Figs. C.4 and C.5). The extended emission region is also seen with the SiO (5–4) transition (Wong et al. 2016), suggesting that shocks may play a significant role in the origin of SO and SO_2 . In addition, the measured temperature of this arc-like structure is lower than that

of the central region. Distributions of ^{34}SO and SO_2 show a morphological correlation in R Dor. The emissions of these species are clumpy, but the emission peak appears at the centre of the continuum peak. For R Leo, we find ordered emission structures of SO_2 and $^{34}\text{SO}_2$ centred at the continuum peak. We find similar features for SO and ^{34}SO but they are slightly extended compared to the SO_2 and $^{34}\text{SO}_2$ emission. SO and ^{34}SO probably trace the cold gas of the outer layer of the CSE. In W Hya, we find SO_2 and $^{34}\text{SO}_2$ contours with an approximate circular pattern without smooth arcs, and emission peaks are slightly offset from the continuum peak. Apart from the central emission around the continuum peak, additional strong emission peaks appear to the north-west, which is $10''$ offset from the continuum peak. SO_2 and $^{34}\text{SO}_2$ are not resolved with the ACA observations of EP Aqr. In summary, our results support the previous conclusion that SO and SO_2 exhibit a centralized distribution in low-mass-loss rate AGB stars.

Recently, Danilovich et al. (2025) observed low-excitation ($E_{\text{low}} < 100$ K) transitions of SO_2 towards the intermediate-mass AGB star OH 30.1–0.7. The source is thought to have a very high mass loss rate ($>10^{-4} M_{\odot} \text{ yr}^{-1}$). They found shell-like distributions of SO_2 , consistent with the morphology discussed above for sources with high-mass loss rates. Their results suggest a low excitation temperature for SO_2 around the source. They also conclude that their result aligns with the observed trend that AGB stars with high-mass loss rates tend to show lower-energy SO_2 transitions (Wallström et al. 2024). In contrast, we have detected both low- and high-excitation transitions in a sample of AGB stars with low mass-loss rates. We obtained high excitation temperatures of SO_2 for all sources (see Table 3), of the order of a few hundred Kelvin, indicating that SO_2 originates mainly from the intermediate region of the CSE in all sources. The average high angular resolution of our observations is $\sim 0.2''$ and the MRS $\sim 3''$. Since the emitting regions of the detected species are confined to within $\sim 2.0''$ (see Figs. C.1 C.2, C.3, C.6, C.7, C.8, and C.9), and the MRS is consistently higher than this, we can assume that we have recovered all of the flux. Therefore, we do not expect significant flux loss due to spatial filtering.

The centralized emission of SO and SO_2 towards α Ceti, W Hya, and R Leo could be due to several physical properties of their surrounding environments. Khouri et al. (2018) found dust surrounding both Mira A and B, as well as a dust tail connecting the two stars. They found a region around Mira A that exhibits high gas density close to the star, which is characterized by a steep decline in density at its outer edge. Mira is a binary system, and within this context, the observed SO and SO_2 emission is centrally concentrated but displays an irregular morphology, suggesting that their formation in the inner envelope might be influenced by both local density enhancements and dynamical effects related to binarity. R Leo has an episodic and patchy mass ejection circumstellar environment (Hoai et al. 2023), leading to localized shocks in the inner envelope, which in turn may enhance SO and SO_2 formation, producing the observed centralized emission. The CSE of the stable component of W Hya shows an approximately spherical morphology, with the gas and dust expanding radially (Hoai et al. 2022). The SO and SO_2 emissions are centrally concentrated, indicating an origin in the warm, dense inner envelope, where possibly shocks and dust-formation processes predominate in sulphur chemistry.

Ionization effects are likely to play only a minor role in these sources, as the species reside in the inner envelope, where they are expected to be shielded from the interstellar radiation field. Furthermore, the centralized SO and SO_2 emissions indicate that these molecules are not primarily produced by outer-shell

photochemistry. However, the present data cannot resolve the molecular emission features between the individual components in sources with binary companions. Furthermore, the modelling results of Danilovich et al. (2016) found that, in their sample of low-mass-loss-rate stars, both SO and SO_2 exhibit centrally peaked abundance distributions. In contrast, in higher-mass-loss rate stars, the SO abundance peaks further out in the wind. The SO_2 abundance profile might follow a similar profile (Danilovich et al. 2020). However, dedicated studies of individual sources, combining observations and modelling, are needed to better understand the origin of emission features of sulphur-bearing molecules.

4.2. Sulphur isotopic ratio in various astrophysical environments

All sources in our sample are oxygen-rich AGB stars with a similar metallicity and mass-loss rates of approximately $10^{-7} M_{\odot} \text{ yr}^{-1}$ (see Table 1). We estimated the $^{32}\text{S}/^{34}\text{S}$ ratio for α Ceti and R Leo using column density values that were obtained using RD analysis, and for W Hya, R Dor, and EP Aqr values obtained using MCMC analysis. Our measured isotopic ratios with their uncertainties are listed in Table 3. We estimated a $^{32}\text{S}/^{34}\text{S}$ ratio of 32 ± 6 for α Ceti, 10 ± 3 for R Leo, 26 ± 8 for R Dor, 18 ± 4 for W Hya, and 12 ± 5 for EP Aqr. Comparing these ratios with the solar value of 21.7, our estimated values for R Dor and W Hya are in good agreement with the solar value. The ratios for EP Aqr and R Leo are slightly lower, and for α Ceti, it is slightly higher than the expected value. We also measured $^{33}\text{S}/^{34}\text{S}$ ratio ~ 0.35 for α Ceti, which is higher than the solar value (~ 0.17 , Asplund et al. 2021) by a factor of 2. However, this is only in one source and based on one line of ^{33}SO . Hence, we need multi-transition observations to constrain this ratio and to compare this with different astrophysical environments.

Along with SO_2 and $^{34}\text{SO}_2$ transitions, we also have one pair of SO and ^{34}SO detection towards α Ceti, R Leo, and R Dor. In the case of α Ceti, an approximate $^{32}\text{SO}/^{34}\text{SO}$ ratio of 7 ± 1 was obtained using the integrated intensities of transitions with similar E_{up} and transition probabilities. This is considerably lower than the ratio of 32.4 obtained from SO_2 isotopologue ratio. This could be partly due to opacity effects, which we did not take into account in our calculations. The emitting regions of two molecules may also differ, leading to the over- or under-estimation of the column densities and derived isotopic ratios. For R Dor, we have one transition of SO with high $E_{\text{up}} = 143$ K and ^{34}SO with low $E_{\text{up}} = 80$ K. Also, their transitions are very different. Hence, we did not estimate the $^{32}\text{SO}/^{34}\text{SO}$. Our measured value of $^{32}\text{SO}/^{34}\text{SO}$ is 6.0 ± 0.5 for R Leo, which was obtained using integrated intensities of two transitions with a similar E_{up} and transition probability. This value is also considerably lower than the value of 26 obtained from the SO_2 isotopologue ratio, likely for the same reason as was discussed for α Ceti.

Our measured value of $^{32}\text{SO}_2/^{34}\text{SO}_2$ towards two oxygen-rich AGB stars, R Dor and W Hya, agrees well with the solar value within the uncertainties, indicating that ^{32}S and ^{34}S are not significantly processed during the AGB phase. However, the measured ratio for α Ceti is slightly higher and for R Leo and EP Aqr the measured ratios are lower. The obtained value of the $^{32}\text{S}/^{34}\text{S}$ ratio in R Leo and IRC-10529 is significantly lower than the solar value, which can be attributed to the fact that these sources might have a lower metallicity or have formed in an environment with lower $^{32}\text{S}/^{34}\text{S}$. Another possibility could be the age effect of the different environments. Conversely,

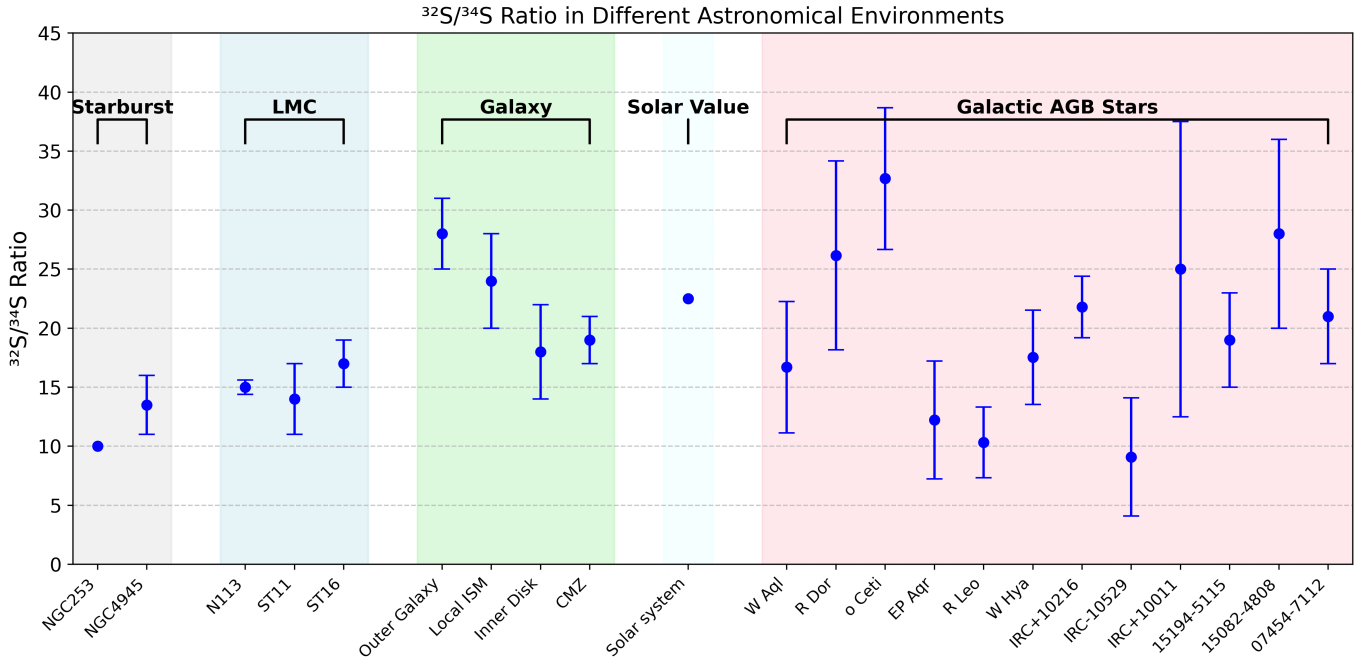


Fig. 7. $^{32}\text{S}/^{34}\text{S}$ ratios for different environments. N113 – Gong et al. (2023); ST11, ST16 – Shimonishi et al. (2020); NGC 253 – Martín et al. (2021); NGC 4945 – Wang et al. (2004); Solar System – Anders & Grevesse (1989); CMZ, inner disc, local ISM, outer Galaxy – Yan et al. (2023); IRC+10216 – Mauersberger et al. (2004); W Aql, IRC+10011, IRC-10529 – Wallström et al. (2024); 15194-5115, 15082-4808, 07454-7112 – Unnikrishnan et al. (2024); o Ceti, R Dor, W Hya, R Leo, EP Aqr – This work. Additionally, Danilovich et al. (2020) reported a ratio of 42 for the Galactic AGB star, IK Tau, which is not included in the plot due to the absence of measured uncertainties.

our estimated $^{32}\text{SO}/^{34}\text{SO}$ ratio deviates from the solar value, which could be attributed to the SO line being optically thick. In such cases, the observed line intensity underestimates the true SO column density, leading to an artificially low isotopic ratio. To verify the accuracy of the $^{32}\text{S}/^{34}\text{S}$ ratio obtained here, it would be useful to derive the sulphur isotopic ratio using other sulphur-bearing species abundant in the CSEs of oxygen-rich AGB stars, such as CS and SiS, in addition to SO and SO_2 , as was previously performed by Velilla Prieto et al. (2017) for IK Tau.

Furthermore, we compared our derived $^{32}\text{S}/^{34}\text{S}$ ratios with different astrophysical environments (see Fig. 7), such as the Large Magellanic Cloud (LMC), representing a low-metallicity environment with a metallicity of $\sim 0.3\text{--}0.5\ Z_{\odot}$ (Westerlund 1997); starburst galaxies, NGC 253 with a metallicity of $\sim 0.5\text{--}1.5\ Z_{\odot}$ (Beck et al. 2022); NGC 4945 with a metallicity of $\sim 0.5\text{--}1.2\ Z_{\odot}$ (Mouhcine et al. 2005; Stanghellini et al. 2015); the Milky Way (the central part has solar metallicity and decreasing metallicity with distance), considering different environments; the outer Galaxy has a metallicity similar to the LMC (Shimonishi et al. 2021), the local interstellar medium (ISM); the inner disc; the CMZ; the Solar System; and finally, a list of AGB stars that includes our work and some literature values. It is clear from Figure 7 that the $^{32}\text{S}/^{34}\text{S}$ ratios for most of the Galactic AGB stars except R Leo and IRC-10529 are comparable to the solar value, including different environments from our Galaxy and Solar System. Since ^{32}S and ^{34}S are primarily produced through explosive nucleosynthesis during Type II and Type Ia SNe, the measured values in AGB stars reflect the abundances in their natal clouds. In contrast, $^{32}\text{S}/^{34}\text{S}$ ratios in LMC and starburst galaxies are lower compared to the solar value. Based on comprehensive calculations, Woosley & Weaver (1995) reported that ^{32}S and ^{33}S are the primary yield (in the sense that the stellar yields do not

strongly depend on the initial metallicity of the stellar model) and do not depend on the initial metallicity of the stellar yield, i.e. metal content of a star at its birth, while ^{34}S is not a clean primary isotope, and its yield decreases with metallicity. Recently, Gong et al. (2023) also reported that $^{32}\text{S}/^{33}\text{S}$ and $^{34}\text{S}/^{33}\text{S}$ in the LMC are lower than in the Milky Way and NGC 253 and concluded that this can be attributed to a combination of age (the LMC may be composed of older stellar populations that have not undergone as much recent nucleosynthesis), low metallicity (which influences the types and outcomes of nuclear reactions in stars), and the star formation history (e.g. rate of star formation).

The $^{32}\text{S}/^{34}\text{S}$ isotopic ratio shows an increasing trend from NGC 253 (youngest evolutionary stage) to the Milky Way (oldest) (see Fig. 7), which may correlate with the age of stellar populations across these galaxies. NGC 253, a starburst galaxy with the lowest $^{32}\text{S}/^{34}\text{S}$ ratio, is dominated by massive young stars due to intense recent star formation, which may indicate minimal isotopic processing. NGC 4945, with a slightly higher $^{32}\text{S}/^{34}\text{S}$ ratio, has a mix of younger and older stars, reflecting a moderately older stellar population than NGC 253. The LMC, with an even higher $^{32}\text{S}/^{34}\text{S}$ ratio, contains stars that are several billion years old, making it older than NGC 4945 but younger than the Milky Way. The Milky Way, with the highest $^{32}\text{S}/^{34}\text{S}$ ratio, hosts the oldest stellar populations, including stars that are over 13 billion years old, suggesting significant nucleosynthetic evolution. This trend indicates that lower $^{32}\text{S}/^{34}\text{S}$ ratios are associated with younger, less processed stellar populations. In contrast, higher ratios correspond to older, more evolved galactic environments where sulphur chemistry has been processed significantly. However, this analysis is based on a small sample. Therefore, we need a large sample and systematic analysis to confirm this trend with different multi-transition sulphur-bearing

species along with different isotopic ratios such as $^{33}\text{S}/^{34}\text{S}$, $^{32}\text{S}/^{33}\text{S}$, and $^{32}\text{S}/^{34}\text{S}$.

5. Conclusions

We have analysed the data of five oxygen-rich AGB stars with low resolution ($\sim 2.5''$) and complementary high resolution ($\sim 0.2''$), obtained with the ALMA ACA and 12m array. The major conclusions of this paper are summarized as follows:

- We report the detection of 16 transitions of SO_2 , 14 of $^{34}\text{SO}_2$, 3 of SO , and 2 of ^{34}SO towards α Ceti; 14 transitions of SO_2 , 2 of $^{34}\text{SO}_2$, and 1 transition each of SO and ^{34}SO towards R Dor; 10 transitions of SO_2 , 7 of $^{34}\text{SO}_2$, and 1 transition each of SO and ^{34}SO towards R Leo; 13 transitions of SO_2 and 3 of $^{34}\text{SO}_2$ towards W Hya; and 5 transitions of SO_2 and 1 of $^{34}\text{SO}_2$ towards EP Aqr;
- We estimate the excitation temperature using detected sulphur-bearing species to be around 200–600 K across different sources and report column densities of $\sim (1-7) \times 10^{16} \text{ cm}^{-2}$ for SO_2 and $\sim (0.8-3) \times 10^{15} \text{ cm}^{-2}$ for $^{34}\text{SO}_2$. Based on the spatial distribution analysis and temperature measurement, we conclude that the detected molecule traces the intermediate region ($\sim 50-500$ au away from the central star) of the CSE;
- We measure the $^{32}\text{S}/^{34}\text{S}$ ratios for five sources, and two of them, R Dor and W Hya, are found to be close to the solar value within uncertainties. For α Ceti, it is slightly higher, and for R Leo and Ep Aqr, the ratios are slightly lower. Overall, 13 Galactic AGB stars have measured $^{32}\text{S}/^{34}\text{S}$ ratios (see Fig. 7), of which 8 are consistent with the solar value within the uncertainties, while 5 show some variations. It is believed that nucleosynthesis in AGB stars does not significantly alter the sulphur isotopic ratio, and the observed values are expected to reflect the abundances of their natal clouds. Therefore, the observed discrepancies in the $^{32}\text{S}/^{34}\text{S}$ ratios among Galactic AGB stars may arise from assumptions of solar metallicity, which might not hold universally, and/or variations in the excitation conditions within their CSEs. Further analysis using different sulphur-bearing species is needed to provide deeper insights, which is beyond the scope of this work;
- We find that SO_2 , $^{34}\text{SO}_2$, SO , and ^{34}SO emission contours are slightly irregular and the emission peaks are offset from the continuum peak of α Ceti. We report an extended emission structure towards the south-west, offset by $\sim 3''$ (~ 300 au) from the continuum peak, where emissions of SO and SO_2 correlate with SiO molecular distribution (Wong et al. 2016). This region exhibits several transitions of SO and SO_2 and some other species reported in Table A.2, likely enhanced by shock chemistry and most probably a low-velocity shock. We estimate the temperature of the arc-like emission to be $\sim 133 \pm 20$ K;
- SO and $^{34}\text{SO}_2$ show clumpy structures in R Dor, with the peak of the molecular emission centred on the continuum peak. Our results are consistent with a previous study by Danilovich et al. (2020);
- The emission of SO_2 , $^{34}\text{SO}_2$, SO , and ^{34}SO transitions in R Leo appear as ordered elliptical concentric emission contours and molecular emission peak slightly offset from the continuum peak. Similar to α Ceti, low-excitation transitions in R Leo show a slightly extended feature in moment 0 maps compared to higher-energy transitions;
- The distribution of SO_2 and $^{34}\text{SO}_2$ in W Hya appears as ordered approximately circular concentric emission

contours, and where the emission peak is slightly offset from the continuum peak as α Ceti. Furthermore, we find an additional structure of strong SO_2 emission located in the north-west, which is $\sim 10''$ (~ 1040 au) offset from the central continuum peak;

- For EP Aqr, we have only low-resolution data with which SO_2 and $^{34}\text{SO}_2$ are mostly unresolved. Hence, we cannot provide a detailed insight into the distribution of detected species;

In summary, we have studied five oxygen-rich AGB stars with comparable low mass-loss rates and find some differences in the distribution and morphology of the emissions from SO_2 , $^{34}\text{SO}_2$, SO , and ^{34}SO across different sources. These differences are likely due to the chemistry being highly sensitive to the physical conditions of the sources, such as the density and temperature structures of the CSEs, source multiplicity, outflows, rotation, and other physical processes associated with the sources. Nonetheless, the mostly centralized distributions of SO and SO_2 in our sample support previous results for low-mass-loss rate AGB stars.

Data availability

The data underlying this article (as presented in the appendix) are available on Zenodo under a Creative Commons Attribution license at <https://zenodo.org/records/17302962>

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Appendix A: Summary of detected species and their transitions

In this section we present the line parameters of all detected species towards a sample of five AGB stars and summary of detected transitions towards *o* Ceti arc-like structure

Table A.1: Line parameters of observed transitions towards our sample AGB stars.

Species	QNs	Freq (MHz)	$\log_{10}(A_{ij})$ (s ⁻¹)	E_{up} (K)	$S_{ij}\mu^2$ (D ²)	Peak Flux (K)	FWHM (km s ⁻¹)	Integrated Intensity (K km s ⁻¹)	τ
<i>o</i> Ceti									
SO ₂	7(4,4) - 6(3,3)	491934.72	-3.02	65.01	10.27	13.89 ± 0.54	8.63 ± 0.39	126.97 ± 7.56	0.0265
SO ₂	12(3,9) - 11(2,10)	494779.73	-3.26	93.96	9.80	12.24 ± 0.37	9.12 ± 0.32	118.34 ± 5.53	0.0234
SO ₂	25(2,24) - 24(1,23)	482503.15	-2.98	303.69	40.61	22.89 ± 0.23	8.79 ± 0.10	213.32 ± 3.30	0.0441
SO ₂	19(9,11) - 20(8,12)	481166.24	-3.83	372.91	4.44	3.15 ± 0.29	9.64 ± 1.03	32.16 ± 4.52	0.0060
SO ₂	38(2,36) - 39(1,39)	481237.15	-5.39	697.69	0.24	1.80 ± 0.13	8.03 ± 0.67	15.34 ± 1.68	0.0034
SO ₂	11(6,6) - 12(5, 7)	331580.24	-4.36	148.95	2.36	2.25 ± 0.20	8.41 ± 0.85	20.04 ± 2.69	0.0042
SO ₂	21(2,20) - 21(1,21)	332091.43	-3.82	219.53	15.20	8.74 ± 0.23	9.74 ± 0.30	90.20 ± 3.64	0.0165
SO ₂	4(3,1) - 3(2, 2)	332505.24	-3.48	31.29	6.92	7.00 ± 0.18	8.89 ± 0.27	65.94 ± 2.62	0.0132
SO ₂	13(2,12) - 12(1,11)	345338.54	-3.62	92.98	13.41	12.50 ± 0.23	8.47 ± 0.18	112.19 ± 3.14	0.0237
SO ₂	20(1,19) - 20(0,20)	282292.81	-4.00	198.88	15.69	10.17 ± 0.25	8.97 ± 0.26	96.77 ± 3.67	0.0192
SO ₂	24(8,16) - 25(7,19)	282636.22	-4.36	432.62	8.12	3.66 ± 0.15	8.40 ± 0.39	32.61 ± 1.99	0.0069
SO ₂	16(0,16) - 15(1,15)	283464.77	-3.57	121.03	33.60	18.98 ± 0.34	8.93 ± 0.18	179.68 ± 4.91	0.0361
SO ₂	17(3,15) - 17(2,16)	285743.59	-3.74	162.93	23.41	13.45 ± 0.21	8.97 ± 0.16	127.86 ± 3.02	0.0255
SO ₂	26(2,24) - 26(1,25)	296168.67	-3.73	340.55	32.84	10.67 ± 0.22	8.44 ± 0.20	95.43 ± 3.05	0.0202
SO ₂	24(4,20) - 24(3,21)	296535.42	-3.59	316.58	41.34	38.12 ± 2.47	8.40 ± 0.63	339.51 ± 33.53	0.0739
SO ₂	18(7,11) - 19(6,14)	297782.59	-4.34	277.30	5.48	3.32 ± 0.21	8.66 ± 0.64	30.49 ± 2.96	0.0062
³⁴ SO ₂	7(4,4) - 6(3,3)	479309.99	-3.06	63.70	10.25	1.31 ± 0.08	11.14 ± 0.96	15.46 ± 1.62	0.0025
³⁴ SO ₂	27(1,27) - 26(0,26)	482025.14	-2.82	329.28	63.92	3.14 ± 0.14	9.77 ± 0.56	32.47 ± 2.34	0.0059
³⁴ SO ₂	5(4,2) - 5(3,3)	345651.29	-3.71	51.76	4.49	0.77 ± 0.12	6.80 ± 1.18	5.53 ± 1.27	0.0014
³⁴ SO ₂	4(4,0) - 4(3,1)	345678.79	-3.88	47.17	2.44	0.77 ± 0.12	6.80 ± 1.24	5.53 ± 1.34	0.0014
³⁴ SO ₂	6(4,2) - 6(3,3)	345553.09	-3.63	57.27	6.34	0.60 ± 0.29	5.98 ± 3.34	3.80 ± 2.80	0.0011
³⁴ SO ₂	7(4,4) - 7(3,5)	345519.66	-3.59	63.70	8.09	1.12 ± 0.14	8.49 ± 1.26	10.10 ± 1.98	0.0021
³⁴ SO ₂	8(4,4) - 8(3,5)	345168.66	-3.56	71.05	9.77	0.78 ± 0.14	5.57 ± 1.18	4.58 ± 1.28	0.0015
³⁴ SO ₂	9(4,6) - 9(3,7)	345285.62	-3.54	79.32	11.41	0.92 ± 0.09	8.14 ± 0.97	7.93 ± 1.25	0.0017
³⁴ SO ₂	13(4,10) - 13(3,11)	344807.91	-3.50	121.63	17.92	0.85 ± 0.67	9.70 ± 8.77	8.78 ± 10.51	0.0016
³⁴ SO ₂	16(4,12) - 16(3,13)	332836.22	-3.52	163.13	23.23	1.10 ± 0.15	8.82 ± 1.39	10.31 ± 2.15	0.0021
³⁴ SO ₂	19(1,19) - 18(0,18)	344581.04	-3.29	167.66	42.24	4.16 ± 0.22	8.28 ± 0.50	36.54 ± 2.92	0.0078
³⁴ SO ₂	17(4,14) - 17(3,15)	345929.35	-3.47	178.77	24.49	1.09 ± 0.13	8.26 ± 1.16	9.54 ± 1.77	0.0020
³⁴ SO ₂	15(7,9) - 16(6,10)	333683.92	-4.28	226.38	3.78	1.38 ± 0.17	5.64 ± 0.78	8.24 ± 1.51	0.0026
³⁴ SO ₂	28(3,25) - 27(4,24)	333364.11	-3.93	402.63	15.68	0.60 ± 0.10	11.66 ± 2.25	7.45 ± 1.90	0.0011
SO	10(9) - 9(9)	295355.70	-5.29	120.24	0.33	4.42 ± 0.17	5.92 ± 0.27	27.75 ± 1.64	0.0083
SO	7(6) - 6(5)	296550.06	-3.49	64.89	13.83	38.12 ± 11.86	8.40 ± 3.02	339.48 ± 161.25	0.0739
SO	8(8) - 7(7)	344310.61	-3.29	87.48	18.56	41.14 ± 0.49	8.19 ± 0.11	356.97 ± 6.47	0.0802
³⁴ SO	8(7) - 7(6)	333900.98	-3.32	79.86	16.24	7.00 ± 0.26	6.66 ± 0.29	46.62 ± 7.24	0.0132
³⁴ SO	7(7) - 6(6)	295396.33	-3.49	69.87	16.15	8.83 ± 0.23	6.17 ± 0.19	57.74 ± 3.72	0.0166
R Dor									
SO ₂	10(3,7) - 10(0,10)	479317.52	-4.92	72.71	0.20	0.09 ± 0.01	12.52 ± 2.01	1.18 ± 0.25	0.0002
SO ₂	19(9,11) - 20(8,12)	481166.24	-3.83	372.91	4.44	0.11 ± 0.02	8.25 ± 1.84	0.98 ± 0.29	0.0003
SO ₂	25(2,24) - 24(1,23)	482503.15	-2.98	303.69	40.61	2.42 ± 0.07	7.73 ± 0.24	19.85 ± 0.82	0.0062
SO ₂	7(4,4) - 6(3,3)	491934.72	-3.02	65.01	10.27	1.50 ± 0.04	7.75 ± 0.24	12.28 ± 0.50	0.0038
SO ₂	12(3,9) - 11(2,10)	494779.73	-3.26	93.96	9.80	1.49 ± 0.06	8.04 ± 0.36	12.68 ± 0.75	0.0038
SO ₂	18(0,18) - 17(1,17)	321330.17	-3.39	151.50	39.15	50.06 ± 1.36	9.53 ± 0.30	505.98 ± 20.99	0.1345
SO ₂	40(5,35) - 39(6,34)	321420.53	-4.08	823.69	17.34	6.45 ± 0.58	7.53 ± 0.78	51.44 ± 7.02	0.0162
SO ₂	22(8,14) - 23(7,17)	321782.62	-4.22	389.37	6.95	4.89 ± 0.85	4.49 ± 0.89	23.25 ± 6.13	0.0122
SO ₂	34(4,30) - 33(5,29)	322475.41	-4.05	594.66	15.82	10.28 ± 0.72	9.89 ± 0.80	107.77 ± 11.55	0.0261
SO ₂	11(2,10) - 10(1,9)	323026.46	-3.73	70.21	10.90	21.89 ± 1.07	9.40 ± 0.53	218.00 ± 16.28	0.0565
SO ₂	22(1,21) - 22(0,22)	323526.42	-3.87	238.04	15.57	21.47 ± 1.50	9.87 ± 0.80	224.69 ± 23.93	0.0554
SO ₂	8(2,6) - 7(1,7)	334673.35	-3.90	43.15	4.95	10.13 ± 0.38	9.93 ± 0.43	106.65 ± 6.18	0.0257
SO ₂	23(3,21) - 23(2,22)	336089.23	-3.57	276.02	28.41	30.17 ± 0.56	8.70 ± 0.19	278.10 ± 7.88	0.0789
SO ₂	16(7, 9) - 17(6,12)	336669.58	-4.23	245.11	4.34	5.57 ± 0.39	7.56 ± 0.61	44.68 ± 4.80	0.0140
³⁴ SO ₂	18(4,14) - 18(3,15)	323806.69	-3.54	195.47	27.04	4.11 ± 2.12	5.16 ± 3.09	22.48 ± 17.77	0.0130
³⁴ SO ₂	27(1,27) - 26(0,26)	482025.14	-2.82	329.28	63.92	0.23 ± 0.04	6.12 ± 1.06	1.52 ± 0.35	0.0006
SO	11(10) - 10(10)	336553.80	-5.21	142.88	4.34	0.92 ± 0.12	6.61 ± 0.97	6.44 ± 3.1	0.0022
³⁴ SO	8(7) - 7(6)	333900.98	-3.32	79.86	16.24	2.95 ± 0.2616	9.60 ± 1.12	30.02 ± 6.30	0.0073
R Leo									
SO ₂	25(2,24) - 24(1,23)	482503.15	-2.98	303.69	40.61	0.46 ± 0.02	7.98 ± 0.46	3.90 ± 0.30	0.0012
SO ₂	7(4, 4) - 6(3,3)	491934.72	-3.02	65.01	10.27	0.47 ± 0.03	6.19 ± 0.51	3.07 ± 0.33	0.0012
SO ₂	12(3, 9) - 11(2,10)	494779.73	-3.26	93.96	9.80	0.26 ± 0.03	7.58 ± 0.99	2.12 ± 0.36	0.0007
SO ₂	28(2,26) - 28(1,27)	340316.41	-3.59	391.79	32.05	4.61 ± 0.08	8.32 ± 0.16	40.67 ± 1.05	0.0119
SO ₂	21(8,14) - 22(7,15)	341275.52	-4.16	369.13	6.37	2.20 ± 0.07	7.30 ± 0.27	17.01 ± 0.83	0.0056

Table A.1: continued.

Species	QNs	Freq (MHz)	$\log_{10}(A_{ij})$ (s^{-1})	E_{up} (K)	$S_{ij}\mu^2$ (D^2)	Peak Flux (K)	FWHM ($km\ s^{-1}$)	Integrated Intensity ($K\ km\ s^{-1}$)	τ
SO ₂	40(4,36) - 40(3,37)	341403.07	-3.39	808.36	71.97	8.47 ± 0.14	6.39 ± 0.13	57.39 ± 1.49	0.0221
SO ₂	36(5,31) - 36(4,32)	341673.96	-3.36	678.51	68.31	6.93 ± 0.12	7.46 ± 0.15	54.83 ± 1.47	0.0181
SO ₂	34(3,31) - 34(2,32)	342761.63	-3.46	581.92	50.74	3.77 ± 0.09	9.78 ± 0.26	39.05 ± 1.35	0.0097
SO ₂	14(4,10) - 14(3,11)	351873.87	-3.46	135.87	19.63	4.76 ± 0.10	8.14 ± 0.20	41.03 ± 1.36	0.0124
SO ₂	12(4, 8) - 12(3,9)	355045.52	-3.47	111.00	16.30	24.23 ± 0.41	4.53 ± 0.09	116.36 ± 3.01	0.0650
³⁴ SO ₂	12(4,8) - 12(3,9)	342332.01	-3.52	109.67	16.34	0.62 ± 0.07	7.46 ± 1.03	4.87 ± 0.89	0.0015
³⁴ SO ₂	14(7,7) - 15(6,10)	353002.39	-4.24	212.59	3.23	0.08 ± 0.08	3.69 ± 3.99	0.32 ± 0.45	0.0001
³⁴ SO ₂	21(4,18) - 21(3,19)	352082.92	-3.44	250.78	30.92	0.79 ± 0.04	12.81 ± 0.77	10.67 ± 0.85	0.0020
³⁴ SO ₂	19(8,12) - 20(7,13)	354397.76	-4.16	326.16	5.23	0.39 ± 0.11	2.71 ± 0.90	1.11 ± 0.48	0.0009
³⁴ SO ₂	24(9,15) - 25(8,18)	355322.20	-4.11	467.40	7.27	0.24 ± 0.11	2.71 ± 1.34	0.70 ± 0.46	0.0005
³⁴ SO ₂	34(3,31) - 34(2,32)	354277.56	-3.44	581.30	48.47	0.79 ± 0.08	5.48 ± 0.64	4.61 ± 0.71	0.0020
³⁴ SO ₂	40(4,36) - 40(3,37)	353949.96	-3.36	807.68	68.60	0.58 ± 0.08	6.68 ± 1.03	4.09 ± 0.84	0.0014
SO	7(8) - 6(7)	340714.16	-3.30	81.24	16.24	18.11 ± 0.23	7.09 ± 0.10	136.05 ± 2.62	0.0481
³⁴ SO	8(9) - 7(8)	339857.26	-3.29	77.34	21.11	6.24 ± 0.25	3.40 ± 0.11	22.48 ± 1.70	0.0162
W Hya									
SO ₂	10(3, 7) - 10(0,10)	479317.52	-4.92	72.71	0.20	0.15 ± 0.04	6.28 ± 1.92	1.02 ± 0.41	0.0002
SO ₂	19(9,11) - 20(8,12)	481166.24	-3.83	372.91	4.44	0.18 ± 0.04	7.41 ± 1.77	1.43 ± 0.45	0.0003
SO ₂	25(2,24) - 24(1,23)	482503.15	-2.98	303.69	40.61	1.14 ± 0.04	9.08 ± 0.33	10.97 ± 0.53	0.0019
SO ₂	7(4,4) - 6(3,3)	491934.72	-3.02	65.01	10.27	0.74 ± 0.04	8.35 ± 0.54	6.59 ± 0.57	0.0012
SO ₂	12(3,9) - 11(2,10)	494779.73	-3.26	93.96	9.80	0.48 ± 0.07	7.30 ± 1.16	3.74 ± 0.79	0.0008
SO ₂	36(10,26) - 37(9,29)	250816.79	-4.47	857.17	13.40	1.75 ± 0.12	7.27 ± 0.57	13.48 ± 1.40	0.0027
SO ₂	32(4,28) - 31(5,27)	252563.89	-4.38	531.10	14.32	3.04 ± 0.09	6.79 ± 0.24	21.87 ± 1.01	0.0048
SO ₂	38(5,33) - 37(6,32)	253935.88	-4.39	749.09	16.29	3.52 ± 1.28	6.72 ± 2.34	25.05 ± 12.62	0.0056
SO ₂	30(9,21) - 31(8,24)	266943.32	-4.41	625.92	10.76	3.84 ± 0.07	5.87 ± 0.12	23.90 ± 0.65	0.0062
SO ₂	13(3,11) - 13(2,12)	267537.45	-3.82	105.82	18.32	8.55 ± 0.12	8.60 ± 0.14	77.99 ± 1.63	0.0141
SO ₂	28(4,24) - 28(3,25)	267719.84	-3.67	415.88	55.04	12.96 ± 0.19	7.13 ± 0.12	97.94 ± 2.18	0.0215
SO ₂	9(5,5) - 10(4, 6)	268168.33	-4.62	102.70	2.02	4.86 ± 6.57	6.34 ± 9.89	32.63 ± 67.40	0.0079
SO ₂	11(3,9) - 11(2,10)	268169.79	-3.81	844.26	15.81	4.86 ± 6.62	6.34 ± 9.96	32.63 ± 67.86	0.0079
³⁴ SO ₂	11(3,9) - 11(2,10)	253936.32	-3.89	82.05	15.42	3.52 ± 1.29	6.72 ± 2.36	25.05 ± 12.73	0.0056
³⁴ SO ₂	27(1,27) - 26(0,26)	482025.14	-2.82	329.28	63.92	0.23 ± 0.03	6.52 ± 0.95	1.62 ± 0.31	0.0004
³⁴ SO ₂	7(4,4) - 6(3,3)	479309.99	-3.06	63.70	10.25	0.15 ± 0.04	6.28 ± 1.95	1.02 ± 0.42	0.0002
Ep Aqr									
SO ₂	10(3,7) - 10(0,10)	479317.52	-4.92	72.71	0.20	0.14 ± 0.02	10.28 ± 2.10	1.50 ± 0.41	0.0008
SO ₂	19(9,11) - 20(8,12)	481166.24	-3.83	372.91	4.44	0.15 ± 0.04	6.95 ± 2.26	1.09 ± 0.47	0.0009
SO ₂	25(2,24) - 24(1,23)	482503.15	-2.98	303.69	40.61	1.52 ± 0.06	8.63 ± 0.41	13.91 ± 0.88	0.0094
SO ₂	7(4,4) - 6(3,3)	491934.72	-3.02	65.01	10.27	2.10 ± 0.08	5.55 ± 0.25	12.37 ± 0.75	0.0130
SO ₂	12(3,9) - 11(2,10)	494779.73	-3.26	93.96	9.80	1.81 ± 0.07	6.74 ± 0.29	12.92 ± 0.73	0.0112
³⁴ SO ₂	27(1,27) - 26(0,26)	482025.14	-2.82	329.28	63.92	0.10 ± 0.02	8.02 ± 1.92	0.88 ± 0.38	0.0006

Table A.2: Species detected towards *o* Ceti arc-like structure

Species	QNs	Freq (MHz)	$\log_{10}(A_{ij})$ (s^{-1})	E_{up} (K)	$S_{ij}\mu^2$ (D^2)
SO ₂	4(3,1) - 3(2, 2)	332505.24	-3.48	31.29	6.92
SO ₂	7(4,4) - 6(3,3)	491934.72	-3.02	65.01	10.27
SO ₂	12(3,9) - 11(2,10)	494779.73	-3.26	93.96	9.80
SO ₂	16(0,16) - 15(1,15)	283464.77	-3.57	121.03	33.60
SO ₂	20(1,19) - 20(0,20)	282292.81	-4.00	198.88	15.69
SO	7(6) - 6(5)	296550.06	-3.49	64.89	13.83
SO	8(8) - 7(7)	344310.61	-3.29	87.48	18.56
³⁰ SiO	J = 7-6	296575.74	-2.86	56.94	67.18
Si ¹⁸ O	J = 7-6	282434.73	-2.93	54.22	67.18
²⁹ SiO	J = 8-7	342980.84	-2.67	74.08	76.79
CO	J = 3-2	345795.99	-5.62	55.31	0.05
¹³ CO	J = 3-2	330587.96	-5.95	31.73	0.07
PO	J = 13/2-11/2, $\Omega = 1/2$, F= 7- 6, l=e	283586.81	-3.91	50.26	6.92
PO	J = 13/2-11/2, $\Omega = 1/2$, F= 6- 5, l=f	283785.40	-3.91	50.32	5.92

Appendix B: Observed and synthetic spectra

Here we present the synthetic spectra obtained with MCMC fitting and the observed spectra for R Dor, R Leo, W Hya, and Ep Aqr.

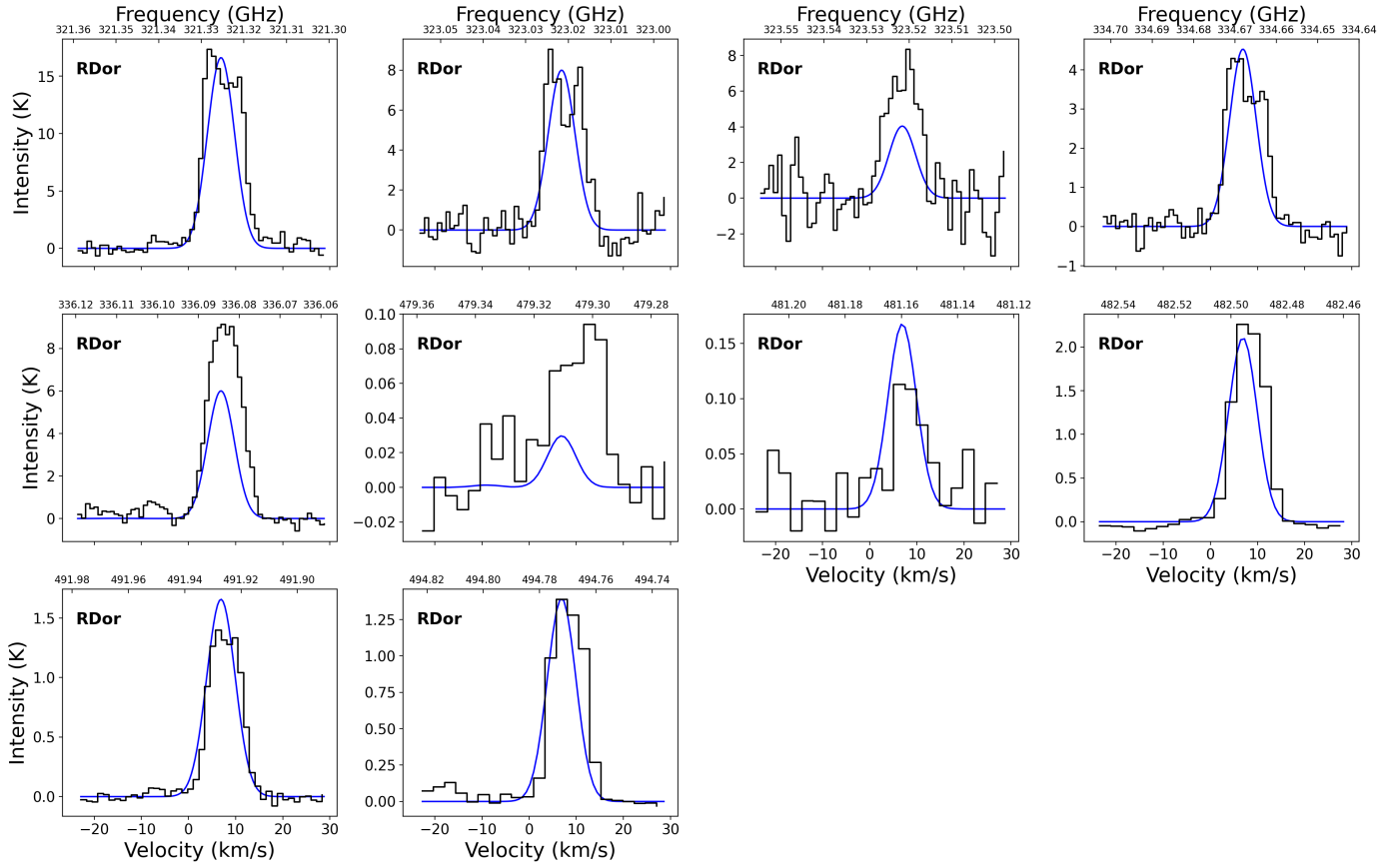


Fig. B.1: Observed and synthetic spectra of SO₂ transitions towards R Dor. The black line represents the observed spectrum, while the blue line represents the modelled spectrum.

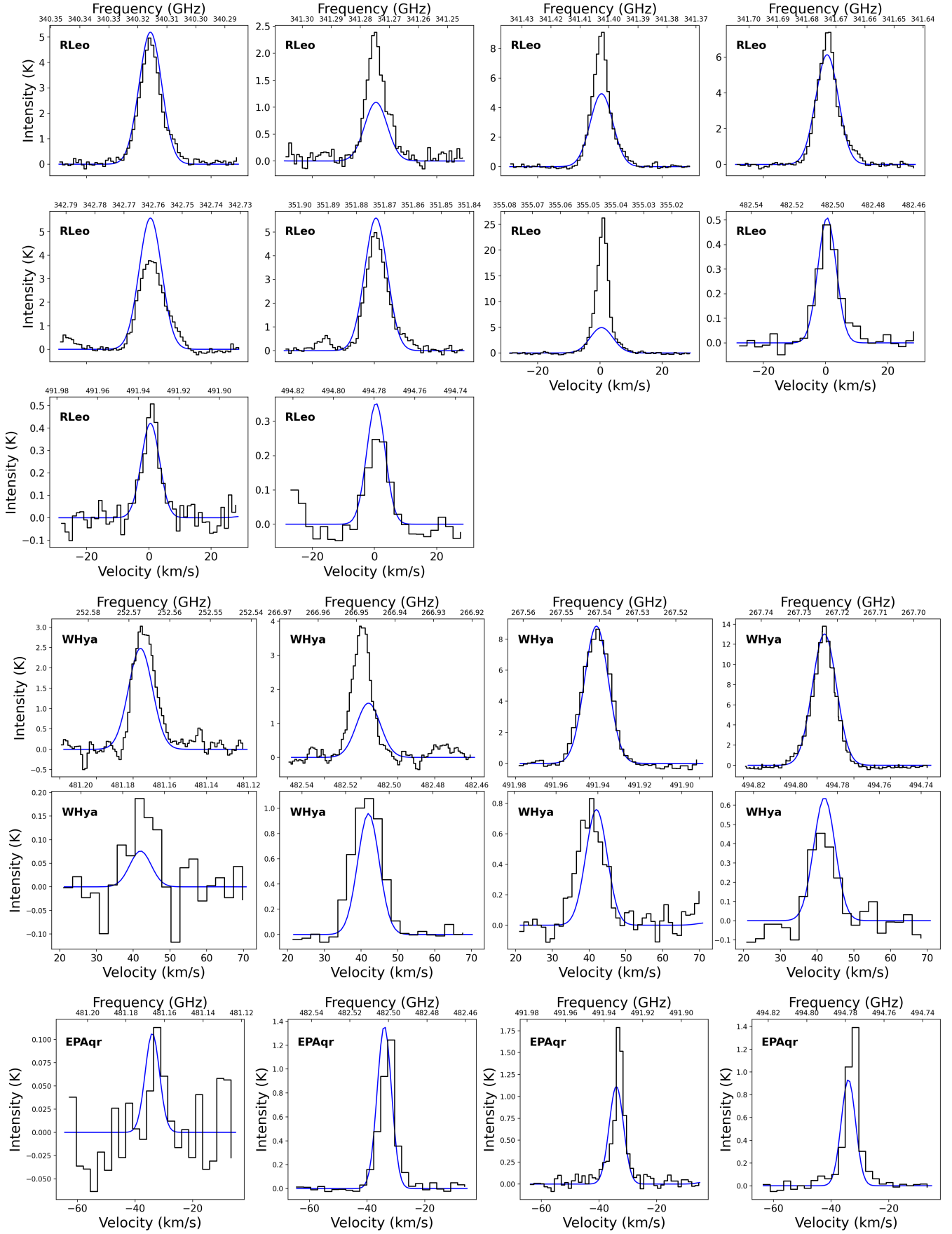


Fig. B.2: Observed and synthetic spectra of SO₂ transitions towards R Leo, W Hya, and EP Aqr. The black line represents the observed spectrum, while the blue line represents the modelled spectrum.

Appendix C: Moment 0 and channel maps

In the appendix, we present moment 0 maps of all detected strong transitions, as well as channel maps of SO and SO₂, which are provided to illustrate the extended atmosphere of *o* Ceti. Figures C.1, C.2, and C.3 show the moment 0 maps of the detected transitions towards *o* Ceti (one additional figure (Fig. 1) of moment 0 maps of SO₂ is provided in the supplementary document). The channel maps of SO and SO₂ are presented in Figures C.4 and C.5, respectively. Moment 0 maps of the detected transitions towards R Dor are shown in Figure C.6 (one additional figure (Fig. 2) of moment 0 maps of SO₂ is provided in the supplementary document). For R Leo, the moment 0 maps are shown in Figure C.7 and C.8 (one additional figure (Fig. 3) of moment 0 maps of SO₂ is provided in the supplementary document). Figure 4 in supplementary document and Fig. C.9 display the moment 0 maps of the transitions identified towards W Hya. Moment 0 maps of the detected transitions towards EP Aqr are available in the supplementary document (Fig. 5).

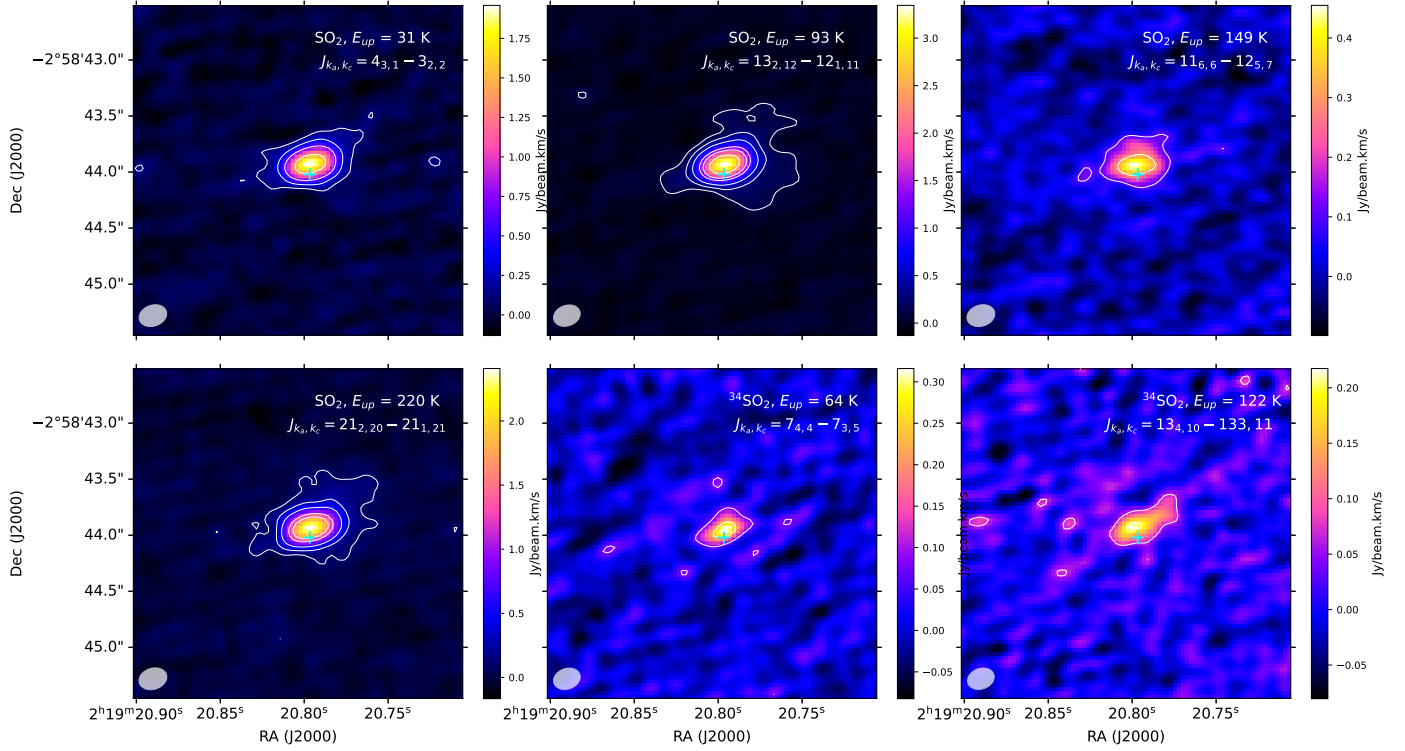


Fig. C.1: Moment 0 maps of SO₂ transitions using ALMA Band 7 observations towards source *o* Ceti. Contours are drawn with steps of 3σ , 9σ , 18σ , 36σ , 48σ , 72σ , and 96σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

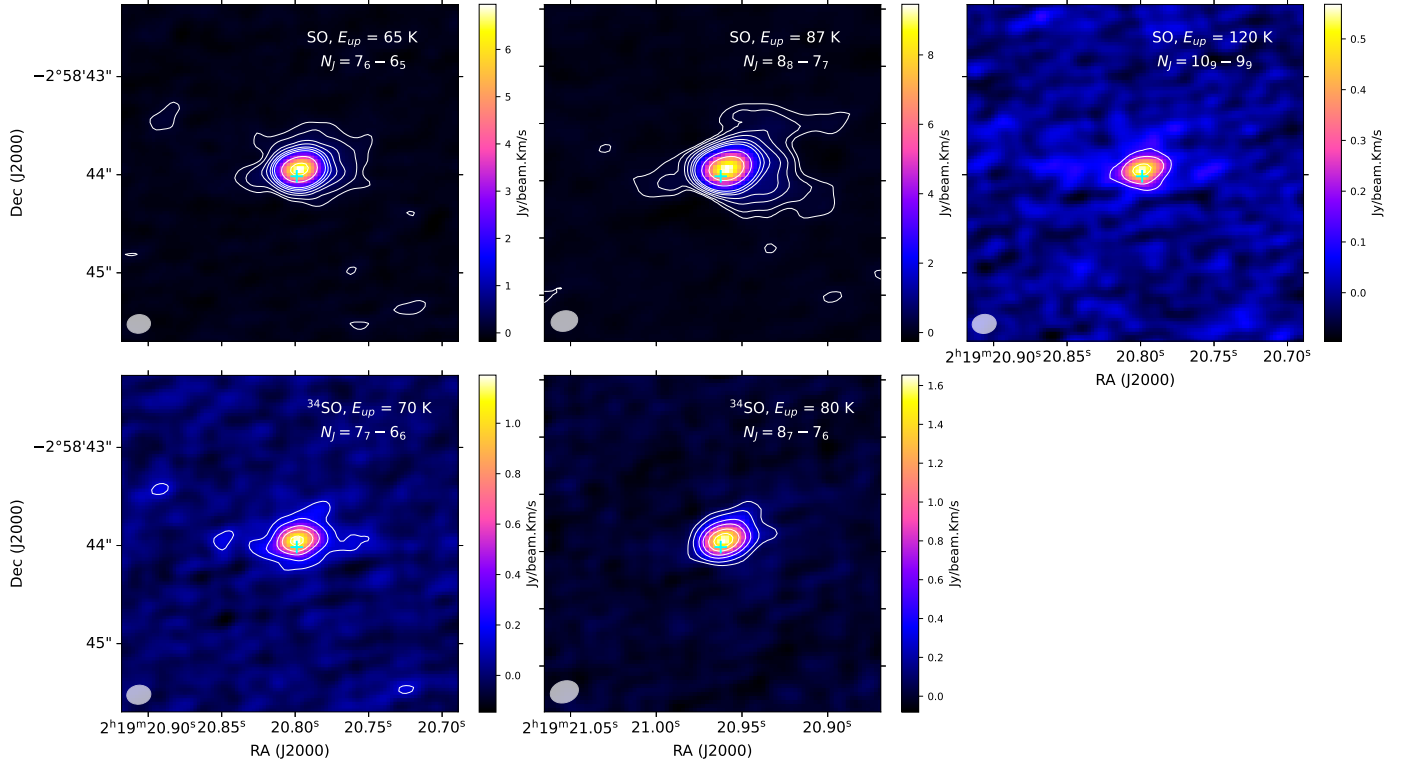


Fig. C.2: Moment 0 maps of SO and ^{34}SO transitions using ALMA Band 7 observations towards source ρ Ceti. Contours are drawn with steps of 3σ , 9σ , 18σ , 36σ , 48σ , and 72σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

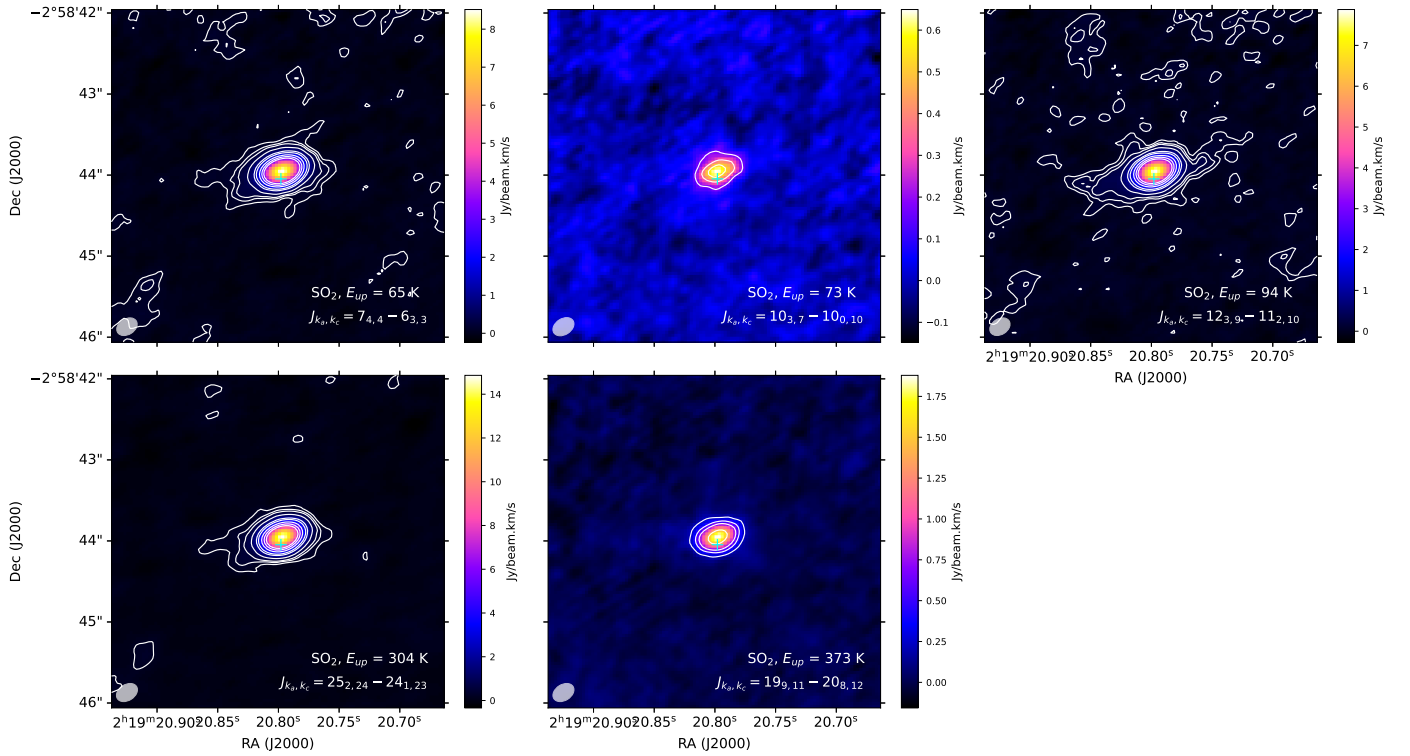


Fig. C.3: Moment 0 maps of SO_2 transitions using ALMA Band 8 observations towards source ρ Ceti. Contours are drawn with steps of 3σ , 9σ , 18σ , 36σ , 48σ , 72σ , and 96σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

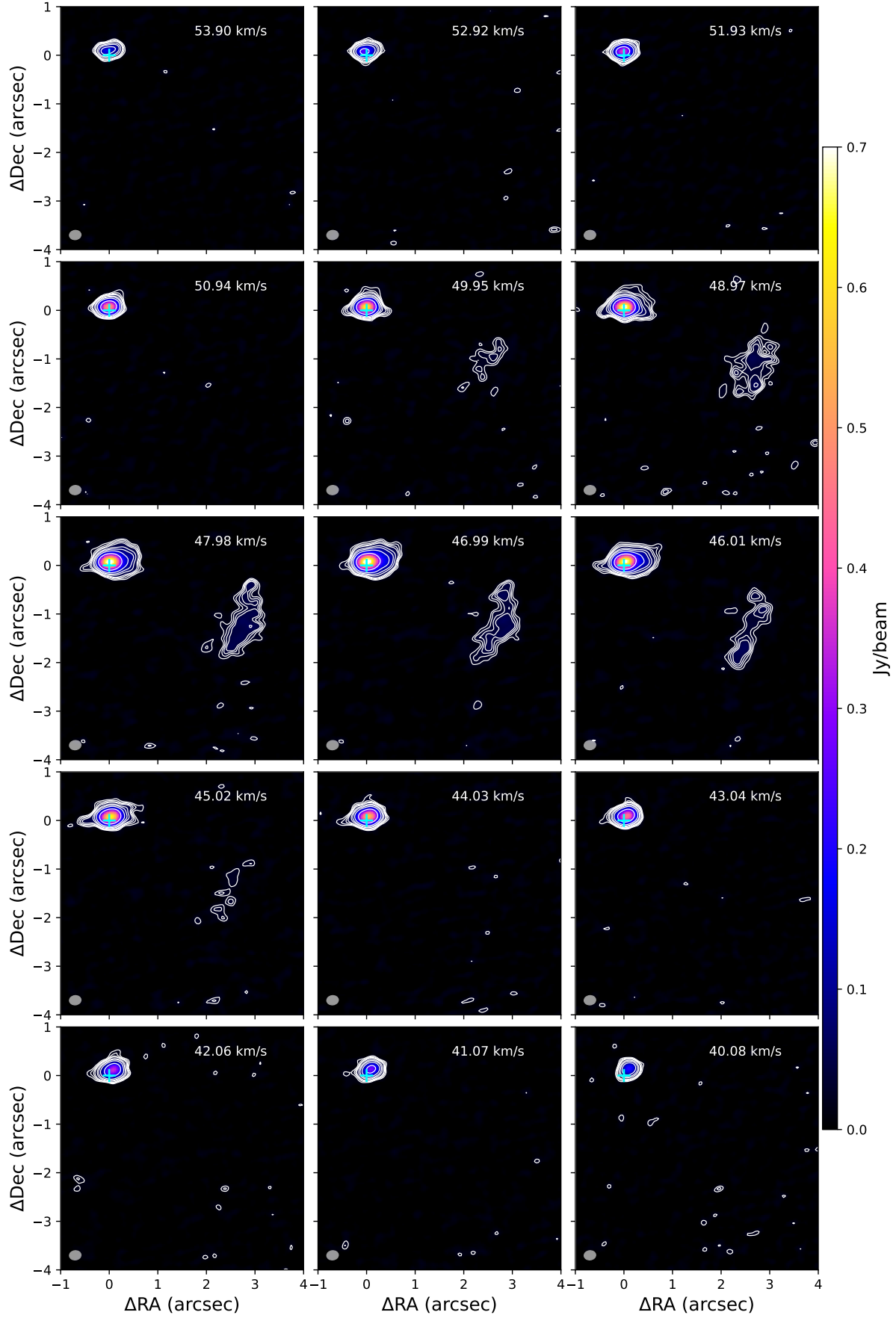


Fig. C.4: Channel maps of SO ($J_{K_a, K_c} = 67 - 56$) low excitation transition ($E_{up} = 65$ K) towards *o* Ceti. Contours are drawn with steps of 3σ , 4σ , 5σ , 6σ , 9σ , 12σ , 18σ , and 36σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

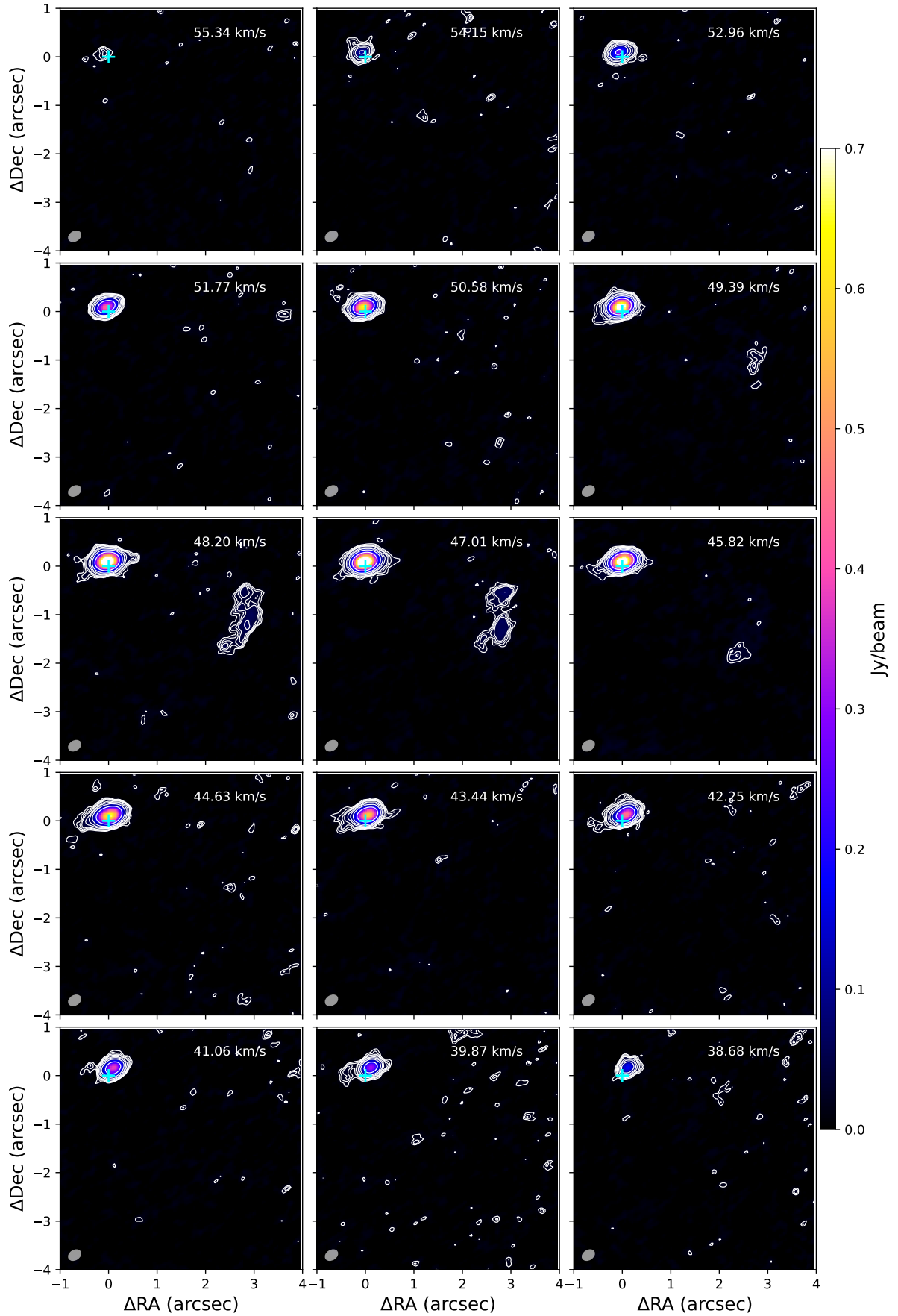


Fig. C.5: Channel maps of SO_2 ($J_{K_a, K_c} = 7_{4,4} - 6_{3,3}$) low excitation transition ($E_{up} = 65$ K) towards ρ Ceti. Contours are drawn with steps of 3σ , 4σ , 5σ , 6σ , 9σ , 12σ , 18σ , and 36σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

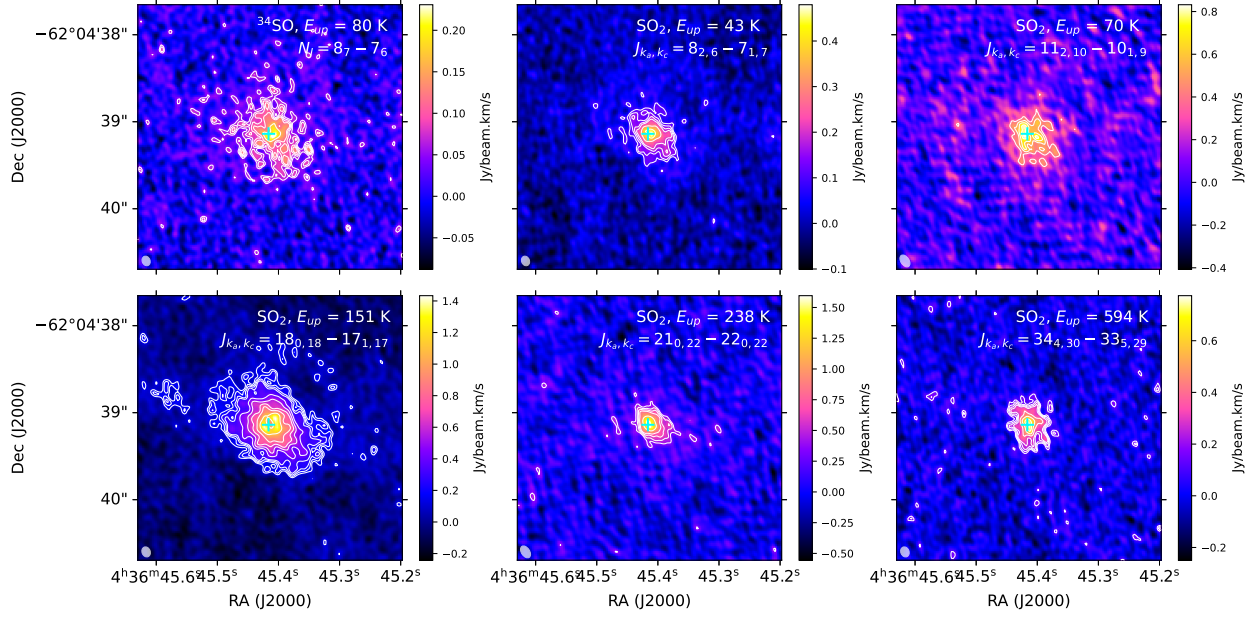


Fig. C.6: Moment 0 maps of SO₂ transitions towards R Dor using ALMA 12m array observations. Contours are drawn with steps of 3σ , 4σ , 5σ , 6σ , 9σ , 12σ , 15σ , 18σ , and 24σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

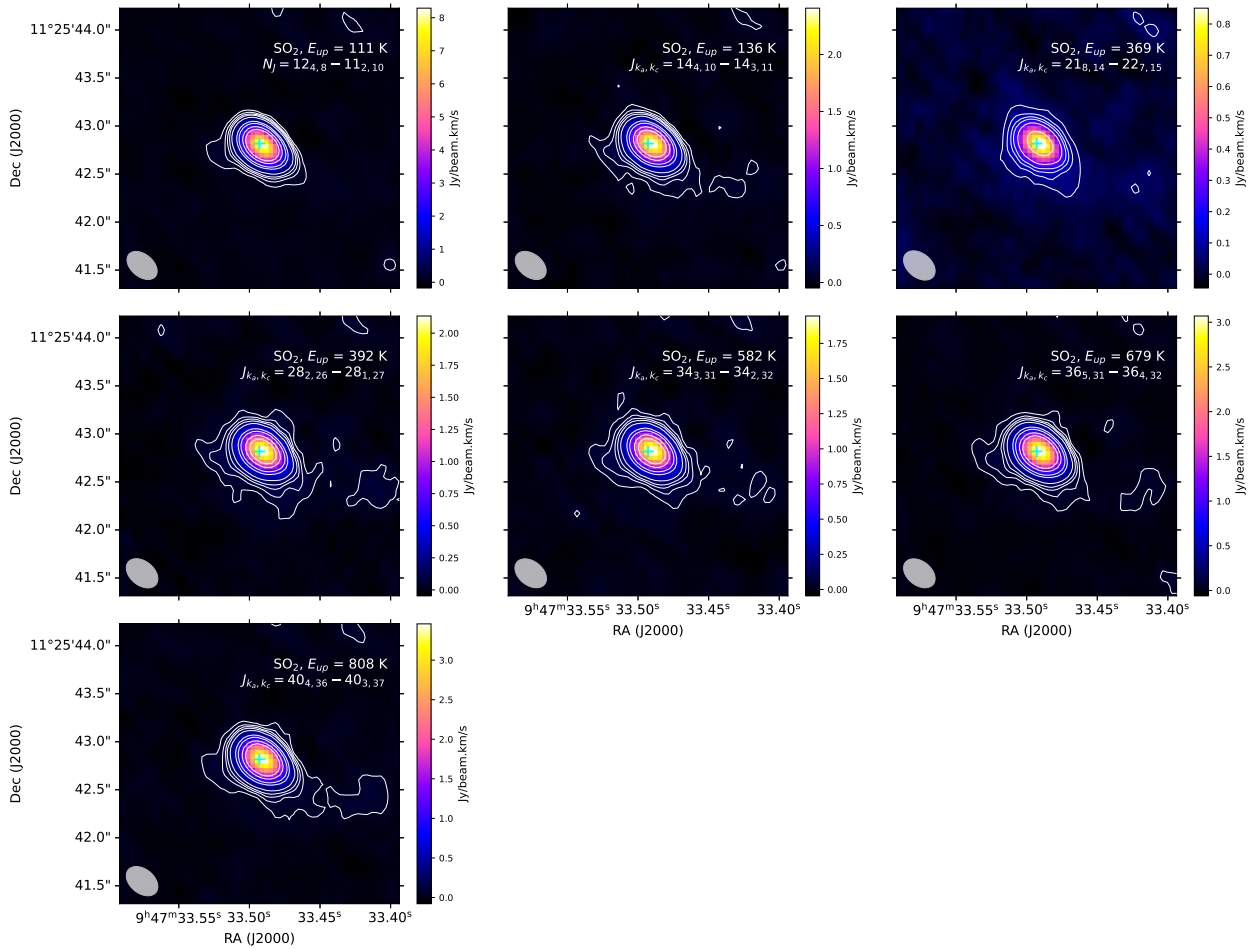


Fig. C.7: Moment 0 maps of SO₂ transitions towards R Leo using ALMA 12m array observations. Contours are drawn with steps of 3σ , 4σ , 5σ , 6σ , 9σ , 12σ , 18σ , 36σ , 48σ , 72σ , and 96σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

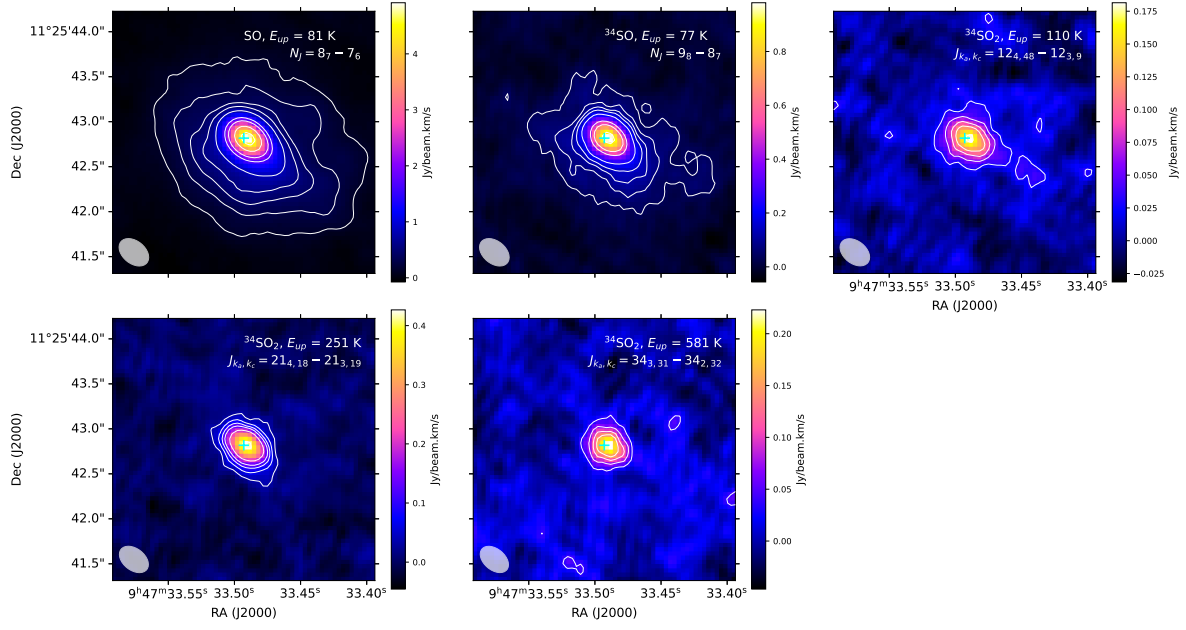


Fig. C.8: Moment 0 maps of SO, ^{34}SO , and $^{34}\text{SO}_2$ transitions towards R Leo using ALMA 12m array observations. Contours are drawn with steps of 3σ , 6σ , 9σ , 12σ , 18σ , 36σ , 48σ , 72σ , and 96σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.

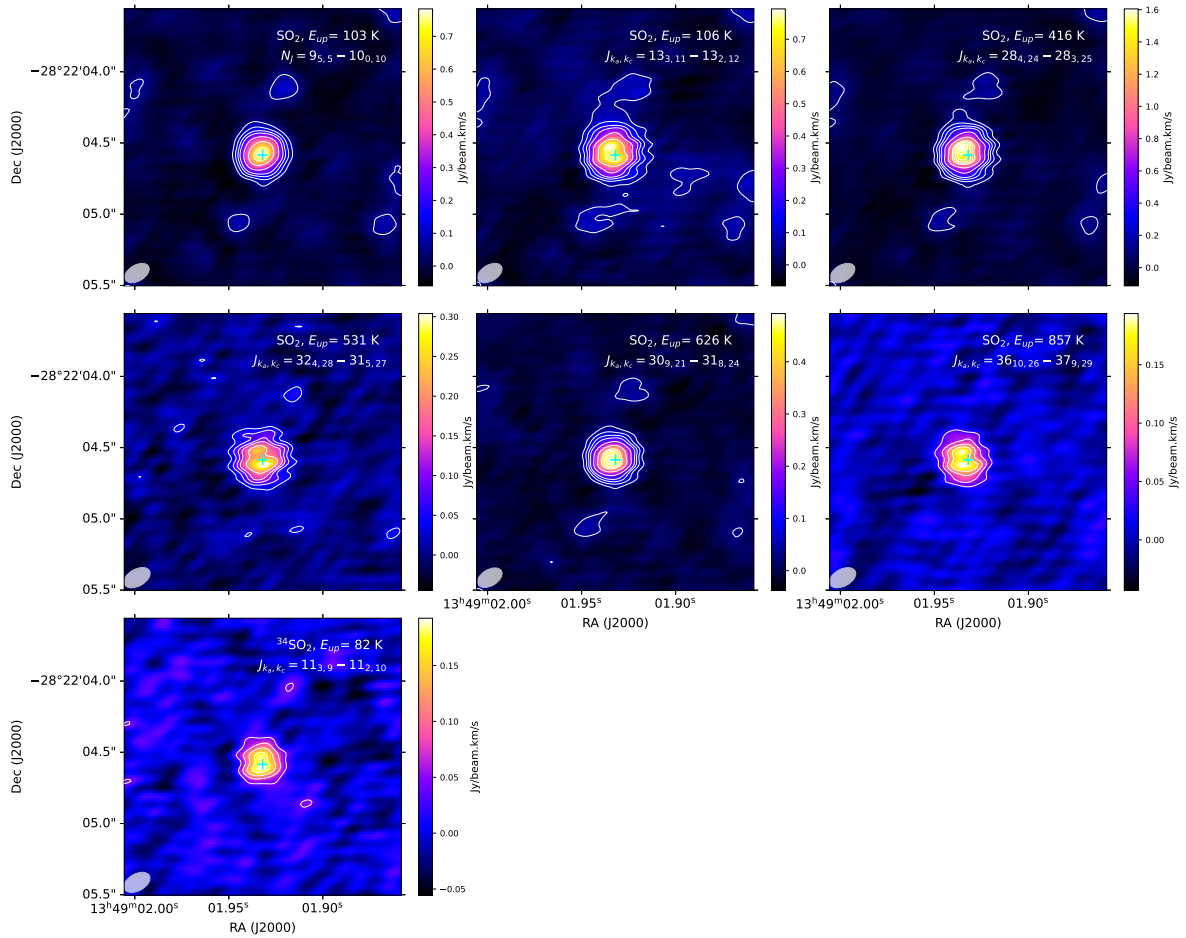


Fig. C.9: Moment 0 maps of SO_2 transitions towards W Hya using ALMA 12m array observations. Contours are drawn with steps of 3σ , 4σ , 5σ , 6σ , 9σ , 12σ , 18σ , and 36σ . The beam size is shown at the lower left corner of each subplot as a filled grey ellipse. The cyan cross symbol indicates the position of the continuum peak.