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Not so-dark: High resolution H I imaging of J0139+4328 and identification of an optical counterpart

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ABSTRACT

Dark galaxies – systems rich in neutral hydrogen (HI) gas but with no stars – are a common prediction of numerous theoretical models and cosmological simulations. However, the unequivocal identification of such sources in current HI surveys has proven challenging. In this work, we present interferometric follow-up observations with the Karl G. Jansky Very Large Array (VLA) of a former dark galaxy candidate J0139+4328, originally detected with the single-dish FAST telescope. The improved spatial resolution of the VLA data allows us to identify a faint optical counterpart and characterize the galaxy. Located at a distance of about 31 Mpc, J0139+4328 has a stellar mass of $3 \times 10^6 M_{\odot}$ and a relatively high gas richness of $M_{\text{HI}}/M_{*} = 18$. Despite its high ratio, the galaxy is consistent, within the scatter, with the stellar-to-HI mass relation of HI-selected samples in the literature and with the baryonic Tully–Fisher relation (BTFR), although its kinematic measurement is subject to large uncertainties. This case highlights the potential of modern high-sensitivity HI surveys for detecting low surface brightness and gas-rich galaxies, but underscores the need for careful interpretation of low-resolution HI data, with potentially large centroid errors, and for sufficiently deep optical imaging to ensure robust identification.

Key words. galaxies: dwarf – galaxies: formation – galaxies: fundamental parameters – galaxies: ISM – galaxies: stellar content

1. Introduction

The search for dark galaxies, that is, galaxies containing gas but no stars, has been of interest for many decades, dating back to early surveys of the neutral hydrogen (HI) content of galaxies (e.g., [Pisano et al. 2002](#); [Doyle et al. 2005](#); [Davies et al. 2006](#)). Theoretically, these systems are thought to be dark matter halos in which neutral gas could not reach the critical density required to trigger star formation. This is attributed to a combination of poor cooling efficiency in low-metallicity environments and suppression of gas collapse due to cosmic ultraviolet (UV) background radiation (e.g., [Kepner et al. 1997](#)). Some analytic models (e.g., [Davies et al. 2006](#); [Benítez-Llambay & Frenk 2020](#)) and cosmological simulations (e.g., [Benítez-Llambay et al. 2017](#); [Lee et al. 2024](#)) predict the existence of such systems and provide further insight into their evolution. In particular, [Lee et al. \(2024\)](#) found that dark galaxies in the Illustris TNG50 simulation ([Vogelsberger et al. 2014](#); [Nelson et al. 2015](#)) tend to reside in dark matter halos with slightly higher spin parameters, have larger gas and dark matter sizes compared to luminous systems of similar mass, and are found in lower-density regions where reduced numbers of interactions allow them to persist without merging with luminous systems. Finding such systems in obser-

vations serves as a direct test of current theoretical models compared to predictions from simulations.

While the existence of dark galaxies is theoretically anticipated, there are very few convincing dark galaxy candidates to date. Numerous HI-rich but optically faint or undetected sources have been identified in untargeted HI surveys such as the Arecibo Legacy Fast ALFA (ALFALFA, [Giovannelli et al. 2005](#)) and the FAST all sky HI survey (FASHI, [Zhang et al. 2024](#)). However, most previous candidates were later revealed to be more likely remnants of tidal interactions or gas stripping (e.g., VIRGOHI21, [Bekki et al. 2005](#); [Taylor et al. 2017](#); SECCO 1, [Bellazzini et al. 2015a,b, 2018](#)). Some of the most promising current candidates include HI 1225+01 ([Giovannelli & Haynes 1989](#); [Chengalur et al. 1995](#); [Matsuoka et al. 2012](#)), AGESVC1 282 ([Taylor et al. 2012, 2013](#); [Bílek et al. 2020](#)), Cloud 9 ([Zhou et al. 2023](#); [Benítez-Llambay & Navarro 2023](#); [Benítez-Llambay et al. 2024](#)), and AC G185.0–11.5 ([Liu et al. 2025](#)). However, all of the above candidates are close to another galaxy or belong to a cluster, complicating their interpretations. The lack of truly isolated candidates questions the existence of these systems and/or their detectability with current HI surveys. Recent untargeted HI surveys such as the APERTURE TIE In Focus (Apertif, [van Cappellen et al. 2022](#); [Adams et al. 2022](#)) and the Widefield ASKAP L-band Legacy All-sky Blind survey (WALLABY, [Koribalski et al. 2020](#)) are continuing to expand the search for such systems, with some candidates already

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proposed (with WALLABY, O’Beirne et al. 2025), which, however, have yet to be confirmed.

Despite the scarcity of unequivocal dark galaxies, deep HI surveys have uncovered a variety of systems that are extremely gas rich and underluminous. These include low-mass galaxies with unusually high gas richness (e.g., Coma P, Janowiecki et al. 2015; AGC 229101, Leisman et al. 2021), and ultra-diffuse galaxies (UDGs), galaxies of low surface brightnesses ($\mu > 24$ mag arcsec⁻²) but large physical sizes ($R_e > 1.5$ kpc; e.g., Leisman et al. 2017; Mancera Piña et al. 2019; Šiljeg et al. 2024). Together, these objects define a broader regime of interest in galaxy evolution: probing the physical conditions under which dark matter halos can accumulate baryons and the circumstances in which those baryons fail to form stars efficiently. The study of such galaxies offers a promising path to refine models of gas accretion, cooling, and feedback in low-mass systems (e.g., Somerville & Davé 2015).

In this work, we report follow-up observations of a galaxy that was initially discovered with FAST and reported by Xu et al. (2023, hereafter X23) as a dark galaxy candidate. Using interferometric HI observations from the Karl G. Jansky Very Large Array (VLA), we were able to refine the position of the HI emission and identify the faint optical counterpart in the Pan-STARRS1 (PS1) data (Chambers et al. 2016). The system remains underluminous, with a log (M_{HI}/L_V) of 0.94, comparable to UDGs (Leisman et al. 2017; Mancera Piña et al. 2019; Janowiecki et al. 2019; Leisman et al. 2021), but at lower masses. Compared with larger samples of HI-selected galaxies, we find that J0139+4328 is not an outlier from the $M_{\text{HI}}-M_*$ scaling relation, although it lies in a relatively unexplored regime of dwarf galaxies with low stellar mass and high gas richness. This case highlights the challenges of associating optical counterparts with HI detections from low-resolution single-dish surveys, particularly at larger distances where the optical counterpart is expected to be faint in existing large-scale optical surveys.

In Sect. 2, we describe the VLA observations and ground-based imaging data. In Sect. 3, we derive and report the global HI and optical properties of the galaxy, and compare them with previous HI detections of the galaxy with FAST. We report on the environment of J0139+4328 and place it in the context of other HI detected galaxies in Sect. 4. Finally, in Sect. 5, we provide our conclusions and considerations for future studies of dark galaxies.

2. Data

2.1. Resolved HI observations

Our target J0139+4328 was observed with the Karl G. Jansky VLA in D-configuration using exploratory time under the program 23A-424. The observations were made in four blocks of three hours each on December 20, 21, 22, and 27, 2023, with 8.8 hours in total on J0139+4328. In addition, observations using the C-array configuration were made under project 25A-251 on July 29, 30 and August 2, 10 of 2025 with a total on-target observing time of 11.7 h. In all observations, the flux and band-pass calibrator 3C48 was observed for 15 minutes at the start of each observing block. J0139+4328 was observed as a set of 15-minute scans interleaved with 5-minute scans on the secondary calibrator J0136+4751. The spectral setup consisted of an 8 MHz bandwidth with 1024 channels giving a channel width of 7.81 kHz corresponding to 1.65 km s⁻¹ (radio convention). The data were flagged for minor radio frequency interference and standard cross-calibrations using the primary and the phase

calibrators were applied to the target data. No self-calibration was performed. Data from both D and C-array configurations were combined and imaged with a robust weighting of -0.5 and a channel width of 1.65 km s⁻¹. Hanning smoothing was applied to the data cube, resulting in a velocity resolution of 3.3 km s⁻¹. The continuum was subtracted by fitting a straight line through the line-free channels. The noise level of the spectral line cube is a 0.64 mJy beam⁻¹ and the beam size is 21''8 × 15''5 (position angle 48°1). This gives a 3- σ column density sensitivity of 3.8×10^{19} cm⁻² for a line width of 10 km s⁻¹. The line cube was initially cleaned without a mask down to five times the noise level. The result of this was smoothed to twice the spatial resolution and with a 5-channel wide Hanning function in velocity. A mask was created from this smoothed cube using a clip level of twice the noise level of this lower-resolution cube. The original line cube was then cleaned using this mask down to a level of half the noise level.

After determining the systemic velocity of the source (Sect. 3.2), we transferred the cube to the rest frame of the source using

$$V_{\text{source frame}} = \frac{v_{\text{sys}} - v_{\text{obs}}}{v_{\text{sys}}} \cdot c \quad (1)$$

where v_{sys} is the systemic frequency of the source, v_{obs} the observed frequency of the spectral line emission, and c the speed of light. The channel width and spectral resolution in this rest frame become 1.66 km s⁻¹ and 3.33 km s⁻¹, respectively. All HI properties of J0139+4328 and all plots showing the HI emission are given in the rest frame of the source.

2.2. Pan-STARRS1

We used PS1 data (Chambers et al. 2016) to characterize the stellar counterpart, as this is the only optical photometric survey with publicly available data in this region. PS1 is a broadband photometric survey conducted using the 1.8-meter telescope located at the Haleakalā Observatory in Hawaii. We use the PS1 cutout service to obtain 3.3' × 3.3' images in the g -, r -, and i -bands centered at the HI centroid position (see Sect. 3.2). The median seeings in g -, r -, and i -bands are estimated by Gaussian fits to foreground stars and are equal to 1.27'', 1.11'', and 1.03'', respectively.

3. Properties of J0139+4328

In this section, we derive the HI and optical properties of J0139+4328. All derived quantities can be found in Table 1, while the methods are reported below.

3.1. Distance

We derive the velocity-based distance to J0139+4328 using the calculators¹ provided by the Extragalactic Distance Database (Kourkchi et al. 2020). We use the numerical action method (NAM) model from Shaya et al. (2017), which is more precise for small distances (the model extends to ~ 38 Mpc). J0139+4328 has a heliocentric optical recessional velocity of 2484 km s⁻¹, corresponding to a velocity of 2651 km s⁻¹ in the Galactic standard of the rest frame (Eq. 4 from Kourkchi et al. 2020). The NAM model gives a distance of (31.03 ± 5.4) Mpc, where the error is estimated using the prescription from Haubner et al. (2025).

¹ <https://edd.ifa.hawaii.edu/>

Table 1. Properties of J0139+4328.

Property	Value
HI	
Centroid	01:39:29.7 +43:28:30
$V_{\text{sys}}^{\text{hel,op}}$	$(2483.5 \pm 0.5) \text{ km s}^{-1}$
D	$(31.03 \pm 5.6) \text{ Mpc}$
f_c	$150 \text{ pc arcsec}^{-1}$
S_{HI}	$(259 \pm 26) \text{ mJy km s}^{-1}$
M_{HI}	$(5.9 \pm 2.2) \times 10^7 M_{\odot}$
W_{50}	$(28.5 \pm 1.2) \text{ km s}^{-1}$
W_{20}	$(43.4 \pm 1.9) \text{ km s}^{-1}$
PA_{kin}	$\sim 150^\circ$
Optical	
Centroid	01:39:29.5 +43:28:30
i	$(56 \pm 7)^\circ$
PA_{op}	$(128 \pm 2)^\circ$
m_g	$20.55^{+0.35}_{-0.26}$
m_r	$20.32^{+0.21}_{-0.18}$
m_i	$20.31^{+0.38}_{-0.27}$
M_*	$3.2^{+6.4}_{-2.0} \times 10^6 M_{\odot}$

Notes. Systemic velocity ($V_{\text{sys}}^{\text{hel,op}}$) is given in heliocentric rest frame using optical convention. All HI properties are reported in the rest frame of the source. f_c is the conversion factor from arcsec to pc. PA_{kin} is defined as the angle of the receding side measured counterclockwise from the north.

3.2. HI properties

We construct the moment zero map by collapsing all channels with visually identified emission (in the heliocentric rest frame and optical convention this corresponds to $2459.1\text{--}2509.4 \text{ km s}^{-1}$; channel maps are shown in Appendix A). Figure 1 shows the moment zero map overlaid on an optical image, with the peak emission clearly situated on top of an optical counterpart. We convert to column densities (N_{HI}) using

$$\frac{N_{\text{HI}}}{\text{cm}^{-2}} = 1.1 \times 10^{24} \frac{S_{\text{HI}}}{\text{Jy km s}^{-1}} \frac{\text{arcsec}^2}{\theta_{\text{min}} \times \theta_{\text{maj}}}, \quad (2)$$

where S_{HI} is the total HI flux and θ_{min} and θ_{maj} are the full widths at half maximum (FWHM) of the minor and major axes of the beam, respectively. The HI centroid is determined by fitting a 2D Gaussian to the innermost $24'' \times 24''$ region of the galaxy. We obtain the position of 01h39m29.7s +43d28m30s, offset by $31''$ from the X23 position obtained with FAST. The moment zero map is then smoothed to $30''$ resolution and clipped at $3\text{-}\sigma$ to define a mask for the original resolution HI cube, then applied to each channel. This procedure includes faint emission into the mask and has been shown to robustly recover the total HI flux of a marginally resolved HI source (Oosterloo et al. 2013; Adams & Oosterloo 2018). Figure 2 shows the global spectrum within this mask. We fit a Gaussian function to derive the HI properties of the source (Table 1). The systemic velocity in the heliocentric rest frame and optical convention corresponds to 2483.5 km s^{-1} . All other HI properties are derived in the rest frame of the source, and correspond to: velocity width of the global profile at 50% of the peak value (W_{50}) of

28.5 km s^{-1} and total flux of $259 \text{ mJy km s}^{-12}$. At a distance of 31.03 Mpc (Sect. 3.1) the derived HI mass of J0139+4328 is $5.9 \times 10^7 M_{\odot}$. Its peak column density, at a resolution of $16'' \times 22''$ is $3.5 \times 10^{20} \text{ atoms cm}^{-2}$.

Figure 3 shows the moment one map, created over the same channel range as the moment zero map, with a mask applied based on a 5σ ($1.6 \times 10^{20} \text{ cm}^{-2}$) threshold in the moment zero map, to retain only high-significance emission. As apparent from the white isovelocity contours, there is a velocity gradient across the source indicating the presence of a rotating disk. This gradient also shows the potential signature of a warp, evident from the change in position angle (measured counterclockwise from the north) from 150° in the center to 190° in the outskirts. We extract multiple position-velocity (PV) slices at position angles between $100^\circ\text{--}280^\circ$. In Fig. 4, we show a PV slice at a position angle of 150° , which shows a clear gradient in velocity. We see no sign of flat rotation in any of the PV slices. Furthermore, as J0139+4328 is marginally resolved, the W_{50} is perfectly consistent with the extent of the high S/N ($\sim 4\sigma$) emission in all extracted PV slices. Hence, we estimate the rotation velocity of J0139+4328 using W_{50} .

We correct the W_{50} using standard prescriptions for instrumental (Verheijen & Sancisi 2001) and thermal broadening (Tully & Fouque 1985; Verheijen & Sancisi 2001) that have been specifically calibrated for W_{50} measurements on spatially resolved galaxies. We assume the contribution from thermal velocity dispersion of $\sigma_v^t = (8 \pm 2) \text{ km s}^{-1}$, a common value found in dwarf galaxies (de Blok & Walter 2006; Stilp et al. 2013). The projected rotational velocity after the above corrections is found to be $(10.2 \pm 2.2) \text{ km s}^{-1}$.

As we cannot determine HI geometry of J0139+4328 due to insufficient spatial resolution, we employ optical geometry for the inclination correction of the rotational velocity. However, we note that the optical position angle does not seem to be fully aligned³ with the kinematic position angle (with $\sim 22\text{--}62^\circ$ difference for $\text{PA}_{\text{kin}} 150\text{--}190^\circ$), making the following estimate potentially uncertain. In addition, a warp may introduce a systematic uncertainty in the rotational velocity estimate, which could be biased toward lower values if the outer disk (assuming it dominates the W_{50} measurement) is warped to a lower inclination than the inner disk. However, the presence and magnitude of this effect are difficult to estimate given the limited spatial resolution of the data. Taking the optically determined inclination of 56° (Sect. 3.4), we obtain a rotational velocity of $V_{\text{rot}} = (12.3 \pm 2.8) \text{ km s}^{-1}$.

3.3. Comparison of HI properties between VLA and FAST

As mentioned in Sect. 1, the galaxy was initially detected with FAST and reported in X23 as a dark galaxy. Additionally, there is another FAST detection of the source within the FASHI catalog (hereafter FASHI detection, Zhang et al. 2024). In this section, we compare the obtained HI properties of the VLA detection with the two detections from FAST. As X23 reported their HI measurements in the heliocentric rest frame and the radio

² This value is obtained by integrating the fitted Gaussian function and is perfectly consistent with the value obtained by integrating the data between -25.7 and 34.1 km s^{-1} where the spectrum is above zero.

³ Beam smearing in HI observations might be shifting the observed PA_{kin} to higher values in the inner parts due to the presence of the warp. Hence, it is possible that HI observations with improved spatial resolution would find PA_{kin} in the central part to be consistent with the optical one.

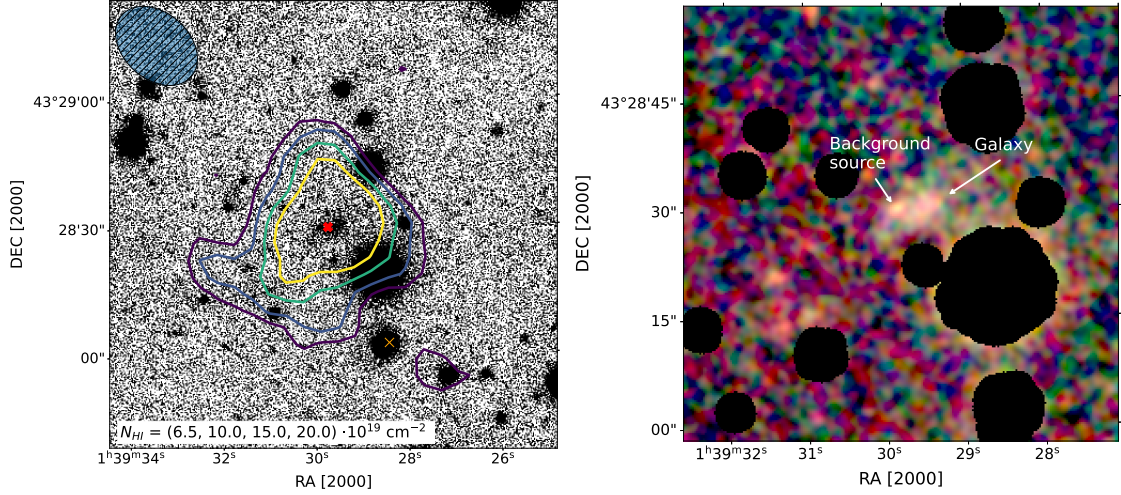


Fig. 1. Identification of the optical counterpart. *Left:* HI contours on optical image obtained by stacking PS1 g , r , and i -band images. The HI contours are $[6.5, 10, 15, 20] \times 10^{19}$ atoms cm^{-2} , where the lowest contour is $\sim 2\text{-}\sigma$. The thick red cross marks the VLA centroid, while the thin orange cross marks the FAST centroid from X23. *Right:* Color image of the galaxy from g (blue), r (green), and i -bands (red), smoothed by a Gaussian kernel of 2 pixels. Masked areas are in black.

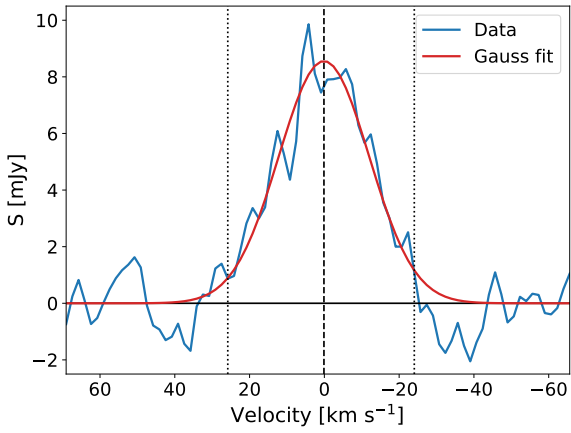


Fig. 2. VLA global HI profile of J0139+4328 in the rest frame of the source. Dotted lines indicate the extent used for creating the moment zero map and the dashed line denotes the determined systemic velocity.

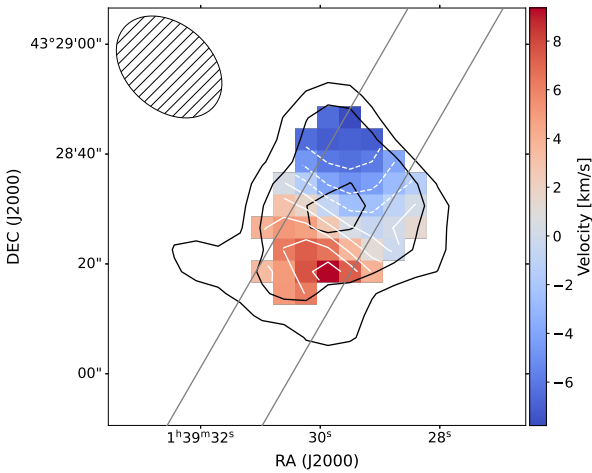


Fig. 3. Moment one map with total intensity HI contours at $[3, 5, 10]\text{-}\sigma$ overlaid in black. White contours are iso-velocity contours separated by 2 km s^{-1} . Full gray lines denote the area from which the position-velocity slice in Fig. 4 was extracted.

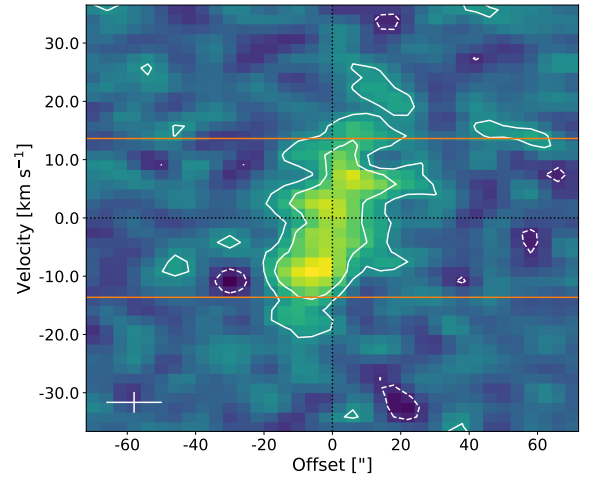


Fig. 4. Position-velocity slice at the $\text{PA} = 150^\circ$. Data contours are shown in white at $[2, 4]\text{-}\sigma$ with full (dashed) lines denoting positive (negative) emission. Horizontal orange lines denote the W_{50} spread. In the lower left corner, we denote the beam size (horizontal line) and velocity resolution (vertical line).

velocity convention, we converted our quantities to the same convention and rest frame for a consistent comparison.

The HI centroid of the X23 (FASHI) detection is $31''$ ($13''$) in the southern direction offset from the VLA one. The relatively large offset between the X23 centroid and the VLA one (which is positioned on top of the optical counterpart) is attributed to the centroiding error of the FAST telescope⁴ and was likely the reason for the misinterpretation of the galaxy as a dark HI cloud. The systemic velocities determined in X23 (2464.4 km s^{-1}) and reported in FASHI (2462.8 km s^{-1}) are well consistent with the VLA estimate of 2463.1 km s^{-1} . The total flux of the X23 detection is $424 \text{ mJy km s}^{-1}$; 1.6 times higher than the VLA flux of $259 \text{ mJy km s}^{-1}$. On the other hand, the flux from the FASHI

⁴ As shown for the FASHI catalog in Figure 9 of Zhang et al. 2024, the centroiding error of the FAST telescope peaks at $\sim 20''$ with a tail towards higher values, so the offset of $31''$ found in this work is not a significant outlier.

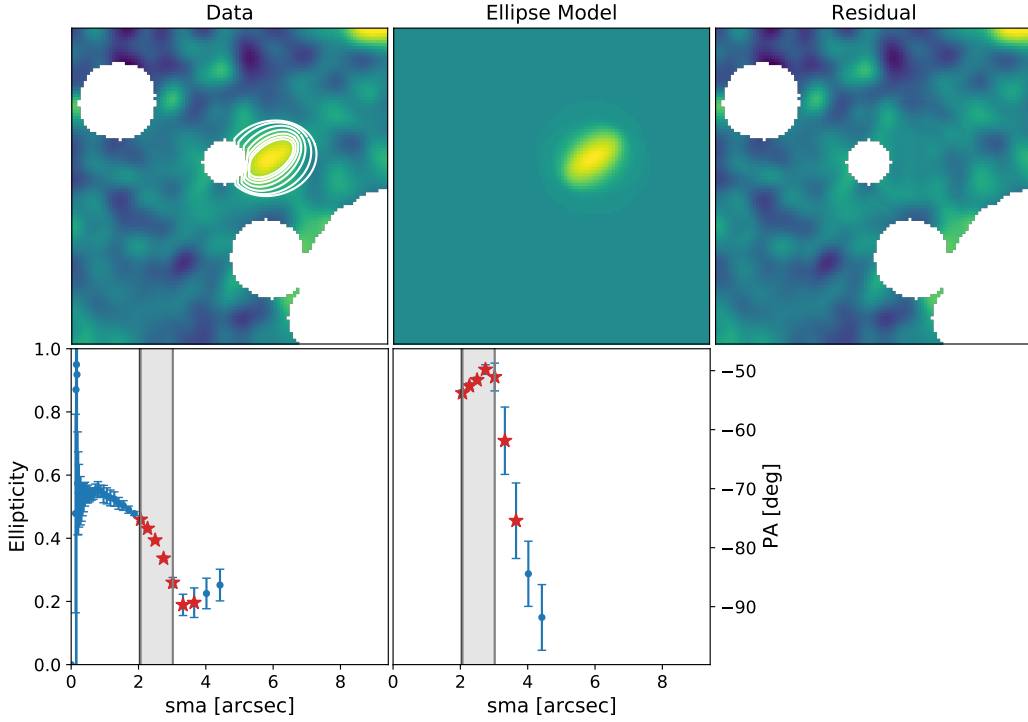


Fig. 5. Isophotal fitting. *Upper panels:* The left panel displays the smoothed image with overlaid ellipses whose corresponding parameters are indicated by red stars in the lower panels. Masked regions appear in white. The central panel presents the model reconstructed from all fitted ellipses whose parameters are shown as blue circles in the bottom panels, while the right panel shows the residual obtained by subtracting the model from the data image. *Lower panels:* Ellipticity and position angle as functions of the semi-major axis length from the second iteration of the fitting procedure. Blue points represent all fitted ellipses from this run, and red stars correspond to those displayed in the upper left panel. The vertical black line indicates the radius equivalent to the PSF’s FWHM after smoothing, and the shaded gray area marks the region used to determine the overall geometry (by taking the median within this range).

catalog amounts to only $256 \text{ mJy km s}^{-1}$ (well within 1σ of the VLA flux). Looking at the global profile, the X23 profile peaks at a higher value of $\sim 12 \text{ mJy}$, while the VLA and the FASHI detections both peak at $\sim 9 \text{ mJy}$. Furthermore, when comparing the shapes of the X23 and the VLA profiles, the VLA spectrum falls off more quickly toward lower velocities, corresponding to the approaching part of the galaxy. Consequently, the W_{50} is higher for the X23 detection (38.9 km s^{-1}) compared to the VLA (28.2 km s^{-1}), while the W_{50} of the FASHI detection is between those two with 35.1 km s^{-1} . With the same peak value, but a larger W_{50} , the FASHI detection likely has a more similar shape to the X23 detection, but is scaled toward lower fluxes, in line with the VLA flux.

While there are dissimilarities between all HI measurements of J0139+4328, we cannot conclusively determine whether their origin⁵ is due to differences in spatial resolution, column density sensitivity, or overall calibration of the HI data. However, it is clear that the source is at most marginally resolved in all of these detections at their respective column density sensitivities. Deeper HI data with at least the same resolution as the VLA and column density sensitivities of FAST ($\sim 10^{18} \text{ cm}^{-2}$) are needed to truly tackle the observed differences between different detections, as well as to robustly characterize the source in terms of its kinematic properties and the morphology of the HI disk.

⁵ Differences in flux measurements are unlikely to result from the VLA’s interferometric short-spacing problem, since the maximum angular extent of the source in the X23 detection ($\sim 6'$) is well below the spatial scales at which the VLA begins to lose sensitivity (10–12', set by the 35.5 m minimum baseline in the C configuration).

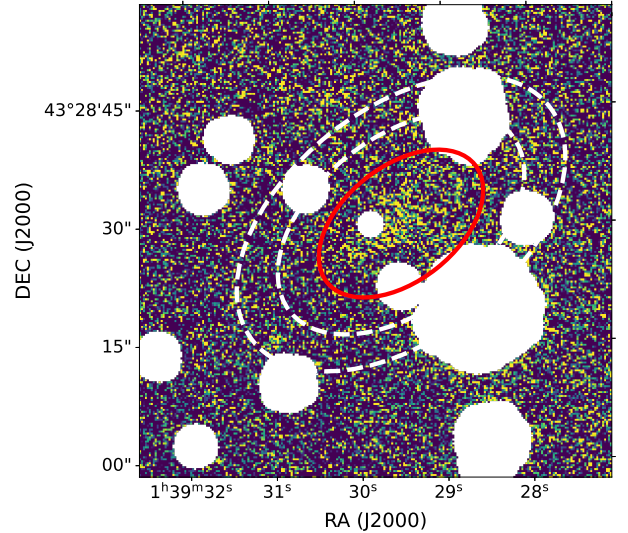


Fig. 6. Aperture photometry on the masked image in the *r*-band. The red ellipse corresponds to the $12''$ (along the semi-major axis) aperture in which the magnitude was measured, while dashed white ellipses denote the annulus within which the background level was estimated at semi-major axes of $18''$ – $24''$.

3.4. Optical properties

In this section, we describe our photometric measurements and the derivation of the stellar mass estimate for J0139+4328. We note that there is a red background source coincident with

the galaxy, as seen in the right panel of Fig. 1, although offset sufficiently from the center to still enable reliable photometry. We masked this source together with foreground stars before estimating the galaxy’s geometry using isophotal fitting. We employed the isophotal fitting procedure described in Šiljeg et al. (2024). In short, we fit the *i*-band image using the Astropy affiliated package `photutils` (Bradley et al. 2022) twice, the first time constraining the position of the center and the second time constraining the position angle and ellipticity. We smoothed the *i*-band image with a 3 pixel smoothing kernel before fitting to increase S/N. We took the median value of each parameter within the region outside the FWHM of the smoothed point spread function (PSF) and the surface brightness limit of 27 mag/arcsec² (below this limit, the geometry is highly influenced by noise in the image). The optical geometry obtained is given in Table 1, and the fitted model is shown in Fig. 5.

Due to the faintness and the small size of the source, we prefer to use fixed aperture photometry instead of a fitted model. Aperture photometry was performed on full spatial resolution images using the optical geometry obtained and a semi-major axis extent of 12″. As seen in Fig. 6, the 12″ aperture incorporates all the visible flux from the galaxy. For confirmation, we have experimented with smaller (6″ and 9″) and larger extents (15″ and 18″), confirming that the obtained magnitudes tend to be fainter in the former case (indicating that some galaxy emission has likely been excluded), and compatible with the chosen extent in the latter (no additional emission was included). For example, a 9″ aperture resulted in systematically ~0.14 fainter magnitudes, while a 15″ aperture yielded magnitudes consistent within 1.5σ. An elliptical annulus of the same shape with semi-major axes from 18″–24″ was used for a local sky background subtraction. The errors in magnitude are estimated by placing the same aperture in 12 positions without any sources in the image and repeating the same procedure. The standard deviation of these values was used as the error estimate for the magnitude obtained. We measured photometry using masked images and report the results in Table 1. We note that the flux within the 12″ aperture of an unmasked image (keeping the same background level measured from the masked image) gave fully compatible results within the 1σ error.

Correcting for Galactic extinction from Schlafly & Finkbeiner (2011) and using the mass-to-light ratio – (*g* – *r*) color relation from Herrmann et al. (2016), we find a stellar mass of $3.2 \times 10^6 M_{\odot}$ ⁶, reported in Table 1. Furthermore, if we transfer to the V-band magnitude using the filter transformations of Tonry et al. (2012), we obtain the V-band luminosity of $L_V = -12.28$. This gives a $\log(M_{\text{HI}}/L_V)$ of 0.94, a value closely comparable to the UDG values (~0.2–1), and significantly larger than the median of the ALFALFA-SDSS sample (~–0.3; Durbala et al. 2020; Leisman et al. 2021).

4. Discussion

4.1. Environment of J0139+4328

We searched for the closest known galaxy to J0139+4328 within a systemic velocity difference of 1000 km s^{–1} using the NASA/IPAC Extragalactic Database⁷ (NED). The closest neighbor to J0139+4328 is NGC 620, with an angular separation of

⁶ During the reviewing process of this paper, Mitrašinović et al. (2026) published another detection of the optical counterpart of J0139+4328 identified using deep optical imaging. The stellar mass estimate obtained in their work is fully consistent with that reported here.

⁷ <https://ned.ipac.caltech.edu/>

1.24° and a nearly identical recession velocity of 2507 km s^{–1} (separated by only 23 km s^{–1}, Theureau et al. 1998). At a distance of 31.03 Mpc, that corresponds to a projected separation of 671 kpc. NGC 620 is included by Karachentsev et al. (2011) in their isolated nearby galaxy sample (“local orphan galaxies” or LOGs). For a distance of 31.03 Mpc, the stellar mass of NGC 620 is $2.7 \times 10^9 M_{\odot}$ (Leroy et al. 2019). Based on abundance matching, NGC 620 has a halo mass of $1.9 \times 10^{11} M_{\odot}$ (Behroozi et al. 2010). This halo has a virial radius of around 155 kpc (assuming the over-density parameter of $\Delta = 300$), much smaller than the projected separation of 671 kpc from J0139+4328. Thus, J0139+4328 is very unlikely to be a satellite galaxy or a tidal dwarf, making it highly probable that its properties (e.g., its large gas richness) are intrinsic to the galaxy.

4.2. Scaling relations

Although initially proposed as a dark galaxy, we have now confirmed the presence of a faint stellar counterpart. Given the low stellar content, we wish to understand if this is a typical dwarf galaxy, simply selected via HI emission and lying near the detection threshold of current optical surveys for its stellar mass and distance, or if it represents an unusual dwarf galaxy, indicative of new populations we might find with advancing HI surveys.

We first look at the $M_{\text{HI}}-M_*$ relation, shown in Fig. 7. We plot the ALFALFA-SDSS catalog (Durbala et al. 2020) for context. We also plot a subsample of the Survey of HI in Extremely Low-mass Dwarfs (SHIELD, McQuinn et al. 2021), a sample selected from the ALFALFA survey with HI mass $M_{\text{HI}} \lesssim 10^{7.2} M_{\odot}$. We include two samples of UDGs as potential analogs of optically faint and gas rich galaxies at higher stellar masses: edge-on HI-bearing UDGs (He et al. 2019) and UDGs from Gault et al. (2021), both selected from the ALFALFA survey. We also add Leo P (Giovanelli et al. 2013; McQuinn et al. 2024) and Leo T (Irwin et al. 2007; Simon & Geha 2007; Ryan-Weber et al. 2008) as local HI-rich faint dwarf galaxies for context. Finally, we add a low surface brightness dwarf galaxy in the Dorado group (LSBG-Dorado-A) which we refer to as LSB-D (Maccagni et al. 2024) and Pavo (Jones et al. 2023, 2025), one of the lowest-baryonic mass-HI-rich galaxies known outside the Local Group so far.

J0139+4328 has a stellar mass similar to the SHIELD sample and the LSB-D, but is significantly gas richer, falling into a regime that has not been well studied before at high resolution. This regime is naturally excluded from SHIELD due to the upper HI mass limit in their selection, and is also excluded from UDG selections due to the small intrinsic sizes of such low stellar mass objects. With a gas richness (M_{HI}/M_*) of 18, J0139+4328 is significantly more gas rich than our local faint dwarfs Leo P and Leo T, and is instead in the regime of UDG galaxy samples, but is an order of magnitude smaller in mass. It would therefore be interesting to explore whether J0139+4328 could be a low-mass counterpart to HI-rich UDGs, that is, having lower surface brightness and larger effective radii compared to its similar mass dwarfs, using deep optical imaging. Finally, while J0139+4328 is gas rich for its stellar mass, it is not a significant outlier from the relation traced by the ALFALFA-SDSS catalog, which is naturally biased toward gas-rich systems at these masses.

Next, we look at the BTFR. The BTFR connects the baryonic mass of the galaxy to the circular velocity of its dark matter halo, and is one of the tightest scaling relations for late-type galaxies. Interestingly, the HI-rich UDGs have been shown to be systematically offset from this relation, having lower circular velocities compared to galaxies of similar baryonic

content (Mancera Piña et al. 2019, 2020). We therefore investigate whether J0139+4328 follows this trend and shares dynamical properties with HI-rich UDGs, or if it aligns more closely with the scaling relations of typical dwarf galaxies.

The BTFR is shown in Fig. 8. We plot various samples from the literature for context. We take 123 galaxies from the SPARC sample (Lelli et al. 2016) with reliable rotation curves (quality flags $Q = 1$ or 2) and inclinations greater than 25° , excluding three galaxies that are part of the LITTLE THINGS subsample. The LITTLE THINGS subsample with 17 galaxies comes from Iorio et al. (2017), which has a more detailed analysis (based on the 3D modeling of the whole data cube). In addition, we incorporated a subsample of 16 galaxies from the SHIELD survey with published velocity measurements (McNichols et al. 2016; McQuinn et al. 2022). We also include two sets of UDGs: a sample of six systems from Mancera Piña et al. (2019, 2020), including the updated velocity for AGC 114905 from Mancera Piña et al. (2024), and a sample of 11 edge-on UDGs from He et al. (2019).

All of the above samples (except the edge-on UDGs from He et al. 2019) are based on spatially well resolved HI measurements (making them more reliable for tracing the true V_{rot}). Furthermore, all resolved HI rotation curves (except the SHIELD subsample of McNichols et al. 2016) were corrected for the asymmetric drift, thereby obtaining the circular speed of the system V_{circ} . For a more consistent comparison with the marginally resolved J0139+4328 for which such detailed analysis is impossible, we additionally overplot the same SHIELD galaxies again, this time applying the same procedure for V_{rot} estimation we used for J0139+4328 (see Sect. 3.2; from W_{50} reported in the ALFALFA catalog and with the same inclinations used for their resolved V_{rot} estimates). We note that in five galaxies our assumed thermal broadening of 8 km s^{-1} (together with instrumental broadening) could explain the whole ALFALFA W_{50} measurement, leaving no space for rotation. Hence, we cannot estimate the rotation velocity using our method in these cases and we leave them out of this analysis.

J0139+4328 has a lower rotational velocity compared to spatially resolved measurements of galaxies with similar baryonic mass. However, after applying the same procedure for V_{rot} estimation to the SHIELD sample, we see that many galaxies move toward lower rotational velocities, in the regime similar to J0139+4328. In addition, we are not tracing the flat part of the rotation curve, and our inclination correction is potentially uncertain due to the slight misalignment between the kinematic and optical position angles. Hence, although J0139+4328 exhibits potentially similar properties to UDGs, we cannot discern its exact position in the BTFR based on our data.

5. Conclusions and future considerations

In this work, we have presented the follow-up VLA HI observations of J0139+4328, a former dark galaxy candidate originally detected with the single-dish FAST telescope (X23). The VLA centroid is $31''$ offset from the FAST detection and is situated on top of a faint optical counterpart. Although J0139+4328 is optically faint, it does not seem to be a significant outlier from the stellar-to-HI mass relation traced by the ALFALFA-SDSS sample nor the BTFR when the scatter is taken into account (although the kinematic measurement is too uncertain to support strong conclusions). It does, however, populate a parameter space of low stellar mass with high gas richness, which has not been well studied before and could potentially exhibit similar properties to UDGs.

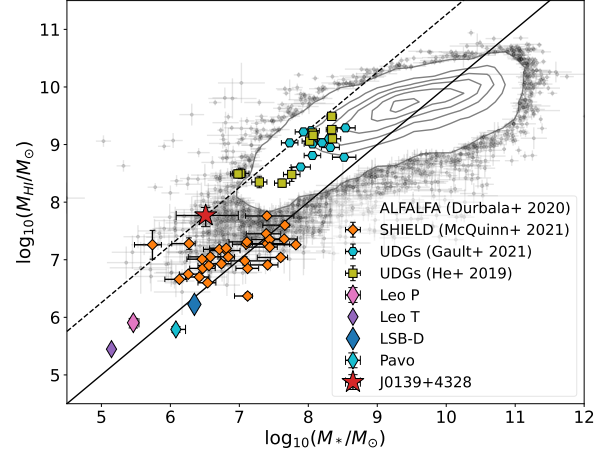


Fig. 7. M_{HI} vs. M_* relation. The full black line denotes the 1:1 relation (gas richness of 1), and the dashed line denotes the 1:18 relation (gas richness of 18).

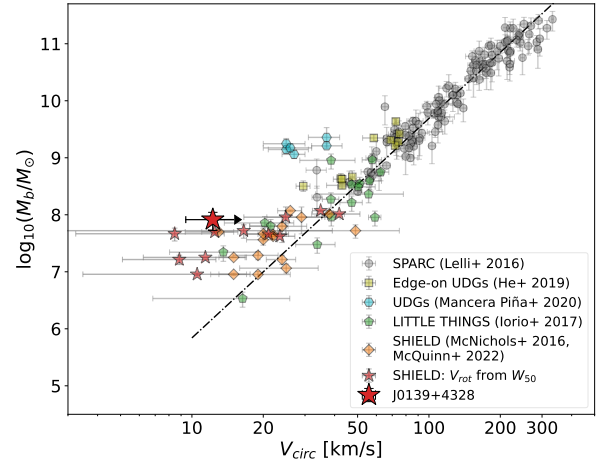


Fig. 8. Position of J0139+4328 on the BTFR. The dash-dotted black line is the best-fit model to the SPARC sample from Lelli et al. (2016).

This case highlights the importance of a cautious interpretation of current state-of-the-art single-dish HI data. In the example of J0139+4328, a galaxy of $3 \times 10^6 M_\odot$ in stellar mass detected at a distance of 31 Mpc, it is clear that the FAST HI surveys are outstripping optical surveys in sensitivity to low-mass, low surface brightness, gas-rich galaxies. The claim of J0139+4328 as a dark galaxy highlights the importance of care in interpreting the shallow optical data when comparing to the deep but low-resolution FAST data. This case demonstrates that single-dish HI data require a meticulous search for the optical counterpart within the centroiding error of the telescope and even beyond, as also seen for some of the detections from ALFALFA (Cannon et al. 2015). Hence, in cases such as the one of J0139+4328, interferometric follow-ups might be crucial for a robust identification of the optical counterpart.

The identification of an optical counterpart to the FAST HI detection highlights that there are very few truly “dark” systems for the sensitivities of surveys to date. Most “dark” HI detections are tidal debris and even those often have optical counterparts with deep enough optical data (e.g., the ALFALFA Virgo 7 cloud complex, Jones et al. 2024; or the WALLABY J103508–283427, O’Beirne et al. 2024). As ongoing and upcoming future HI surveys push to fainter sensitivities, a “dark” space may await

discovery but requires optical data of a matching sensitivity for robust claims. The “almost” dark galaxies present an exciting discovery in and of themselves, especially as new HI observations probe lower column densities, allowing us to explore whether gas disks can persist in more diffuse and tenuous states.

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Appendix A: Channel maps

In Fig. A.1 we show the channel maps of the VLA detection containing emission from the galaxy.

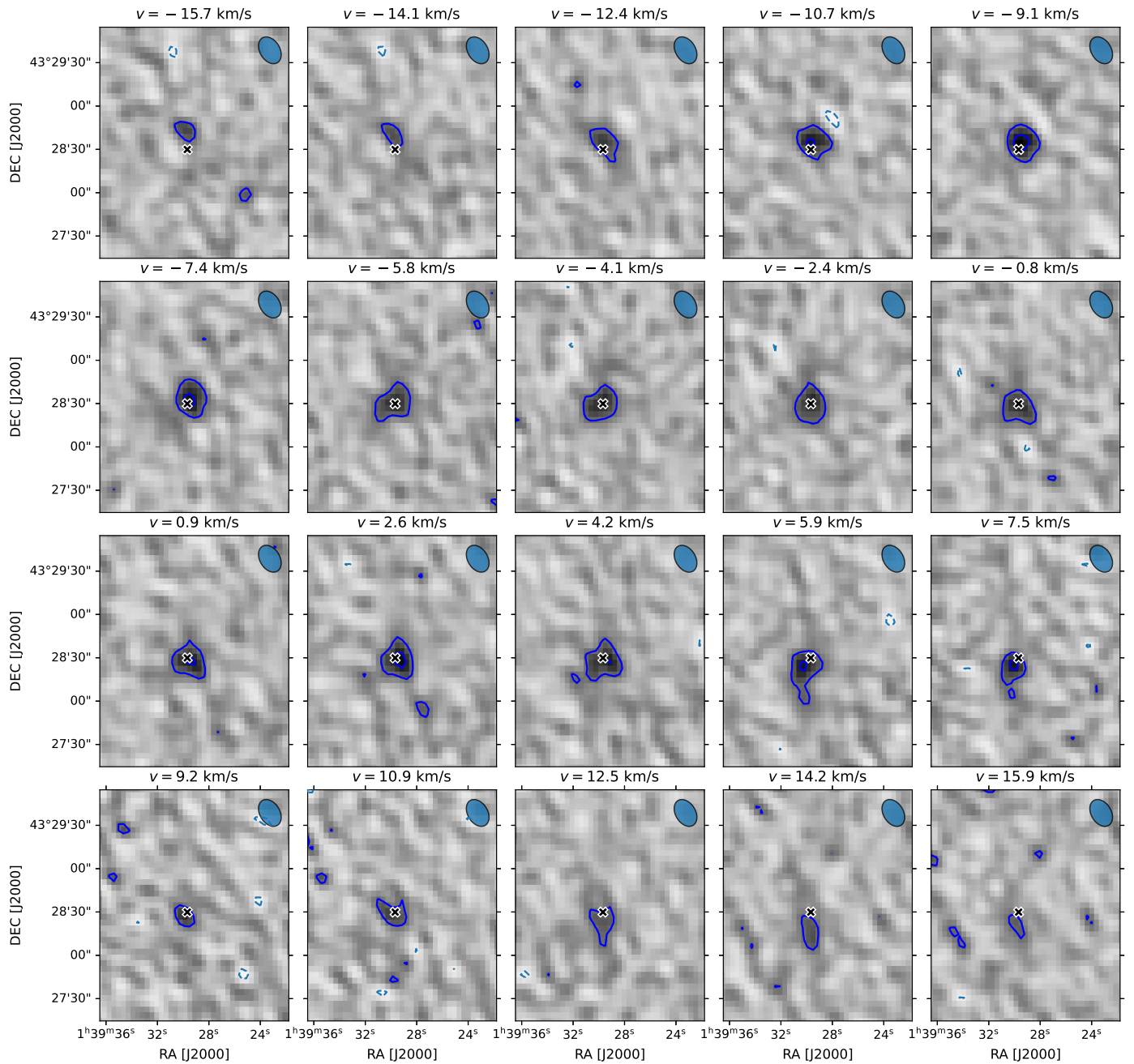


Fig. A.1. Channel maps of the VLA detection in the rest frame of the source. Blue contours represent the data, with dark blue lines denoting positive emission and light dashed blue lines the negative. Contours are plotted starting from 3 times the noise in the cube ($0.65 \text{ mJy beam}^{-1}$), and are spaced by a factor of 2 in intensity. The black X indicates the H I centroid of the galaxy.